

Edge Universality for non-Hermitian Random Matrices

joint work with László Erdős and Dominik Schröder

Giorgio Cipolloni[†], IST Austria

August 21, 2019

Randomness in Physics and Mathematics: From Stochastic Processes to Network, Bielefeld

[†] Partially supported by ERC Advanced Grant No. 338804 and European Union's Horizon 2020 research and innovation programme under the Marie Skłodowska-Curie Grant Agreement No. 665385

- Eugene Wigner predicted that energy gaps of large quantum systems exhibit a universal behaviour and suggested random matrices as a mathematical model to observe this phenomenon (1955).

A brief history

- Eugene Wigner predicted that energy gaps of large quantum systems exhibit a universal behaviour and suggested random matrices as a mathematical model to observe this phenomenon (1955).
- In the last decade universality has been proven for very general **Hermitian** ensembles in the bulk, at the edge and at the cusp of the self consistent density of states.
[Bourgade, C., Erdős, Huang, Johansson, Knowles, Krüger, Landon, Lee, Schlein, Schnelli, Schröder, Sodin, Soshnikov, Sosoë, Tao, Vu, Yau, Yin,...]

A brief history

- Eugene Wigner predicted that energy gaps of large quantum systems exhibit a universal behaviour and suggested random matrices as a mathematical model to observe this phenomenon (1955).
- In the last decade universality has been proven for very general **Hermitian** ensembles in the bulk, at the edge and at the cusp of the self consistent density of states.
[Bourgade, C., Erdős, Huang, Johansson, Knowles, Krüger, Landon, Lee, Schlein, Schnelli, Schröder, Sodin, Soshnikov, Sosoë, Tao, Vu, Yau, Yin,...]
- Not much is known about **universality** for **non-Hermitian** ensembles. Main reason: lack of Dyson Brownian Motion (DBM), that is the backbone for the Hermitian case of **Erdős-Schlein-Yau three step proof strategy**.

A brief history

- Eugene Wigner predicted that energy gaps of large quantum systems exhibit a universal behaviour and suggested random matrices as a mathematical model to observe this phenomenon (1955).
- In the last decade universality has been proven for very general **Hermitian** ensembles in the bulk, at the edge and at the cusp of the self consistent density of states.
[Bourgade, C., Erdős, Huang, Johansson, Knowles, Krüger, Landon, Lee, Schlein, Schnelli, Schröder, Sodin, Soshnikov, Sosoë, Tao, Vu, Yau, Yin,...]
- Not much is known about **universality** for **non-Hermitian** ensembles. Main reason: lack of Dyson Brownian Motion (DBM), that is the backbone for the Hermitian case of **Erdős-Schlein-Yau three step proof strategy**.
- In this talk I report on **universality at the edge of non-Hermitian ensembles**.

Model Assumptions

Let X be an $N \times N$ non-Hermitian matrix with entries i.i.d. real or complex centred random variables with variance N^{-1} , i.e.

$$x_{ab} \stackrel{d}{=} N^{-1/2} \chi, \quad \forall a, b,$$

with $\mathbb{E}\chi = 0$, $\mathbb{E}|\chi|^2 = 1$, and $\mathbb{E}\Re\chi\Im\chi = 0$.

Model Assumptions

Let X be an $N \times N$ non-Hermitian matrix with entries i.i.d. real or complex centred random variables with variance N^{-1} , i.e.

$$x_{ab} \stackrel{d}{=} N^{-1/2} \chi, \quad \forall a, b,$$

with $\mathbb{E}\chi = 0$, $\mathbb{E}|\chi|^2 = 1$, and $\mathbb{E}\Re\chi \Im\chi = 0$. Moreover, we assume:

- χ has bounded moments.

Model Assumptions

Let X be an $N \times N$ non-Hermitian matrix with entries i.i.d. real or complex centred random variables with variance N^{-1} , i.e.

$$x_{ab} \stackrel{d}{=} N^{-1/2} \chi, \quad \forall a, b,$$

with $\mathbb{E}\chi = 0$, $\mathbb{E}|\chi|^2 = 1$, and $\mathbb{E}\Re\chi\Im\chi = 0$. Moreover, we assume:

- χ has bounded moments.
- Mild regularity assumption on the density of χ .

The Circular Law

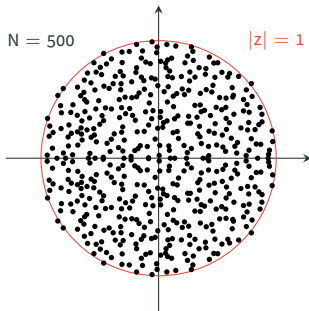


Figure 1: Real entries

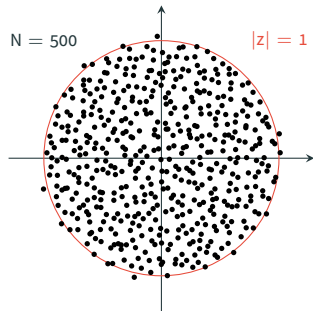


Figure 2: Complex entries with $E x_{ab}^2 = 0$

- **Circular law:** Convergence of the density of states to the uniform distribution on the unit disk.
- Eigenvalues spacing $\sim N^{-1/2}$.
- Accumulation of $\sim \sqrt{N}$ eigenvalues on the real axis for real matrices

Theorem (C.-Erdős-Schröder, '19) For any z_1, \dots, z_k with $|1 - |z_j|^2| \lesssim N^{-1/2}$, and for any test function $F: \mathbb{C}^k \rightarrow \mathbb{C}$, we have

$$\int_{\mathbb{C}^k} F(\mathbf{w}) \left[p_k^{(N)} \left(\mathbf{z} + \frac{\mathbf{w}}{\sqrt{N}} \right) - p_{\mathbf{z}}^{(\infty, \text{Gin}(\mathbb{F}))}(\mathbf{w}) \right] d\mathbf{w} = \mathcal{O} \left(N^{-c(k)} \right),$$

where $\mathbb{F} = \mathbb{R}, \mathbb{C}$.

Theorem (C.-Erdős-Schröder, '19) For any z_1, \dots, z_k with $|1 - |z_j|^2| \lesssim N^{-1/2}$, and for any test function $F: \mathbb{C}^k \rightarrow \mathbb{C}$, we have

$$\int_{\mathbb{C}^k} F(\mathbf{w}) \left[p_k^{(N)} \left(\mathbf{z} + \frac{\mathbf{w}}{\sqrt{N}} \right) - p_{\mathbf{z}}^{(\infty, \text{Gin}(\mathbb{F}))}(\mathbf{w}) \right] d\mathbf{w} = \mathcal{O} \left(N^{-c(k)} \right),$$

where $\mathbb{F} = \mathbb{R}, \mathbb{C}$.

- $p_{\mathbf{z}}^{(\infty, \text{Gin}(\mathbb{F}))}$, the correlation function of the Gaussian case (Ginibre ensemble), is explicitly known [Ginibre, '65; Metha, '67] for the complex case and [Edelman, '93; Akemann-Konziepner, '05; Forrester-Nagao, '07; Borodin-Sinclair, '08] for the real case.

Theorem (C.-Erdős-Schröder, '19) For any z_1, \dots, z_k with $|1 - |z_j|^2| \lesssim N^{-1/2}$, and for any test function $F: \mathbb{C}^k \rightarrow \mathbb{C}$, we have

$$\int_{\mathbb{C}^k} F(\mathbf{w}) \left[p_k^{(N)} \left(\mathbf{z} + \frac{\mathbf{w}}{\sqrt{N}} \right) - p_{\mathbf{z}}^{(\infty, \text{Gin}(\mathbb{F}))}(\mathbf{w}) \right] d\mathbf{w} = \mathcal{O} \left(N^{-c(k)} \right),$$

where $\mathbb{F} = \mathbb{R}, \mathbb{C}$.

- $p_{\mathbf{z}}^{(\infty, \text{Gin}(\mathbb{F}))}$, the correlation function of the Gaussian case (Ginibre ensemble), is explicitly known [**Ginibre, '65; Metha, '67**] for the complex case and [**Edelman, '93; Akemann-Konziepner, '05; Forrester-Nagao, '07; Borodin-Sinclair, '08**] for the real case.
- Previous universality result at the edge under the condition that the first four moments of the entries of X agree with the Gaussian moments [**Tao-Vu, '12**].

Sketch of the proof

- Fix $|1 - |z_0|^2| \lesssim N^{-1/2}$. We introduce the **Ornstein-Uhlenbeck flow**

$$dX_t = -\frac{1}{2}X_t dt + \frac{dB_t}{\sqrt{N}}, \quad X_0 = X, \quad X_\infty \sim \text{Gin}(N), \quad H_t^z := \begin{pmatrix} 0 & X_t - z \\ (X_t - z)^* & 0 \end{pmatrix},$$

and, denoting $G_t^z(i\eta) := (H_t^z - i\eta)^{-1}$, split **Girko's hermitization formula** as

$$\sum_{i=1}^N f(\sqrt{N}(\sigma_i - z_0)) = -\frac{1}{4\pi} \int_{\mathbb{C}} \Delta f(\sqrt{N}(z - z_0)) \left(\int_0^{\eta_0} + \int_{\eta_0}^{+\infty} \right) \Im \text{Tr} G_t^z(i\eta) d\eta dz.$$

Sketch of the proof

- Fix $|1 - |z_0|^2| \lesssim N^{-1/2}$. We introduce the **Ornstein-Uhlenbeck flow**

$$dX_t = -\frac{1}{2}X_t dt + \frac{dB_t}{\sqrt{N}}, \quad X_0 = X, \quad X_\infty \sim \text{Gin}(N), \quad H_t^z := \begin{pmatrix} 0 & X_t - z \\ (X_t - z)^* & 0 \end{pmatrix},$$
and, denoting $G_t^z(i\eta) := (H_t^z - i\eta)^{-1}$, split **Girko's hermitization formula** as
$$\sum_{i=1}^N f(\sqrt{N}(\sigma_i - z_0)) = -\frac{1}{4\pi} \int_{\mathbb{C}} \Delta f(\sqrt{N}(z - z_0)) \left(\int_0^{\eta_0} + \int_{\eta_0}^{+\infty} \right) \Im \text{Tr} G_t^z(i\eta) d\eta dz.$$

- For $\eta_0 \geq N^{-3/4-\delta}$ we can control the flow of $\int_{\eta_0}^{\infty}$ for a long time using **optimal cusp local law** for the linearized matrix H_t^z [**Alt-Erdős-Krüger, '19**]:

$$|\langle \mathbf{x}, (G^z(i\eta) - M^z(i\eta)) \mathbf{y} \rangle | \lesssim \frac{1}{N^{1/2} \eta^{1/3}} + \frac{1}{N\eta},$$

where $M^z \in \mathbb{C}^{N \times N}$ is a deterministic (block-diagonal) matrix.

Sketch of the proof

- Fix $|1 - |z_0|^2| \lesssim N^{-1/2}$. We introduce the **Ornstein-Uhlenbeck flow**

$$dX_t = -\frac{1}{2}X_t dt + \frac{dB_t}{\sqrt{N}}, \quad X_0 = X, \quad X_\infty \sim \text{Gin}(N), \quad H_t^z := \begin{pmatrix} 0 & X_t - z \\ (X_t - z)^* & 0 \end{pmatrix},$$
and, denoting $G_t^z(i\eta) := (H_t^z - i\eta)^{-1}$, split **Girko's hermitization formula** as
$$\sum_{i=1}^N f(\sqrt{N}(\sigma_i - z_0)) = -\frac{1}{4\pi} \int_{\mathbb{C}} \Delta f(\sqrt{N}(z - z_0)) \left(\int_0^{\eta_0} + \int_{\eta_0}^{+\infty} \right) \Im \text{Tr} G_t^z(i\eta) d\eta dz.$$

- For $\eta_0 \geq N^{-3/4-\delta}$ we can control the flow of $\int_{\eta_0}^{\infty}$ for a long time using **optimal cusp local law** for the linearized matrix H_t^z [**Alt-Erdős-Krüger, '19**]:

$$|\langle \mathbf{x}, (G^z(i\eta) - M^z(i\eta)) \mathbf{y} \rangle | \lesssim \frac{1}{N^{1/2} \eta^{1/3}} + \frac{1}{N\eta},$$

where $M^z \in \mathbb{C}^{N \times N}$ is a deterministic (block-diagonal) matrix.

Sketch of the proof

- Fix $|1 - |z_0|^2| \lesssim N^{-1/2}$. We introduce the **Ornstein-Uhlenbeck flow**

$$dX_t = -\frac{1}{2}X_t dt + \frac{dB_t}{\sqrt{N}}, \quad X_0 = X, \quad X_\infty \sim \text{Gin}(N), \quad H_t^z := \begin{pmatrix} 0 & X_t - z \\ (X_t - z)^* & 0 \end{pmatrix},$$
and, denoting $G_t^z(i\eta) := (H_t^z - i\eta)^{-1}$, split **Girko's hermitization formula** as
$$\sum_{i=1}^N f(\sqrt{N}(\sigma_i - z_0)) = -\frac{1}{4\pi} \int_{\mathbb{C}} \Delta f(\sqrt{N}(z - z_0)) \left(\int_0^{\eta_0} + \int_{\eta_0}^{+\infty} \right) \Im \text{Tr} G_t^z(i\eta) d\eta dz.$$

- For $\eta_0 \geq N^{-3/4-\delta}$ we can control the flow of $\int_{\eta_0}^{\infty}$ for a long time using **optimal cusp local law** for the linearized matrix H_t^z [**Alt-Erdős-Krüger, '19**]:

$$|\langle \mathbf{x}, (G^z(i\eta) - M^z(i\eta)) \mathbf{y} \rangle| \lesssim \frac{1}{N^{1/2} \eta^{1/3}} + \frac{1}{N\eta},$$

where $M^z \in \mathbb{C}^{N \times N}$ is a deterministic (block-diagonal) matrix.

- For $\eta_0 \leq N^{-1-\delta}$ we can use the classical smoothing result [**Sankar-Spielman-Teng, '05**] to exclude small eigenvalues λ_i^z of H^z :

$$\mathbb{P} \left(\min_i |\lambda_i^z| \leq N^{-1-\delta} \right) \leq N^{-c\delta}.$$

Sketch of the proof

- Fix $|1 - |z_0|^2| \lesssim N^{-1/2}$. We introduce the **Ornstein-Uhlenbeck flow**

$$dX_t = -\frac{1}{2}X_t dt + \frac{dB_t}{\sqrt{N}}, \quad X_0 = X, \quad X_\infty \sim \text{Gin}(N), \quad H_t^z := \begin{pmatrix} 0 & X_t - z \\ (X_t - z)^* & 0 \end{pmatrix},$$
and, denoting $G_t^z(i\eta) := (H_t^z - i\eta)^{-1}$, split **Girko's hermitization formula** as
$$\sum_{i=1}^N f(\sqrt{N}(\sigma_i - z_0)) = -\frac{1}{4\pi} \int_{\mathbb{C}} \Delta f(\sqrt{N}(z - z_0)) \left(\int_0^{\eta_0} + \int_{\eta_0}^{+\infty} \right) \Im \text{Tr} G_t^z(i\eta) d\eta dz.$$

- For $\eta_0 \geq N^{-3/4-\delta}$ we can control the flow of $\int_{\eta_0}^{\infty}$ for a long time using **optimal cusp local law** for the linearized matrix H_t^z [**Alt-Erdős-Krüger, '19**):

$$|\langle \mathbf{x}, (G^z(i\eta) - M^z(i\eta))\mathbf{y} \rangle| \lesssim \frac{1}{N^{1/2}\eta^{1/3}} + \frac{1}{N\eta},$$

where $M^z \in \mathbb{C}^{N \times N}$ is a deterministic (block-diagonal) matrix.

- For $\eta_0 \leq N^{-1-\delta}$ we can use the classical smoothing result [**Sankar-Spielman-Teng, '05**] to exclude small eigenvalues λ_i^z of H^z :

$$\mathbb{P} \left(\min_i |\lambda_i^z| \leq N^{-1-\delta} \right) \leq N^{-c\delta}.$$

- How to bridge the gap $\eta \in [N^{-1-\delta}, N^{-3/4-\delta}]$?

Sketch of the proof

- Fix $|1 - |z_0|^2| \lesssim N^{-1/2}$. We introduce the **Ornstein-Uhlenbeck flow**

$$dX_t = -\frac{1}{2}X_t dt + \frac{dB_t}{\sqrt{N}}, \quad X_0 = X, \quad X_\infty \sim \text{Gin}(N), \quad H_t^z := \begin{pmatrix} 0 & X_t - z \\ (X_t - z)^* & 0 \end{pmatrix},$$
and, denoting $G_t^z(i\eta) := (H_t^z - i\eta)^{-1}$, split **Girko's hermitization formula** as
$$\sum_{i=1}^N f(\sqrt{N}(\sigma_i - z_0)) = -\frac{1}{4\pi} \int_{\mathbb{C}} \Delta f(\sqrt{N}(z - z_0)) \left(\int_0^{\eta_0} + \int_{\eta_0}^{+\infty} \right) \Im \text{Tr} G_t^z(i\eta) d\eta dz.$$

- For $\eta_0 \geq N^{-3/4-\delta}$ we can control the flow of $\int_{\eta_0}^{\infty}$ for a long time using **optimal cusp local law** for the linearized matrix H_t^z [**Alt-Erdős-Krüger, '19**]:

$$|\langle \mathbf{x}, (G^z(i\eta) - M^z(i\eta)) \mathbf{y} \rangle| \lesssim \frac{1}{N^{1/2} \eta^{1/3}} + \frac{1}{N\eta},$$

where $M^z \in \mathbb{C}^{N \times N}$ is a deterministic (block-diagonal) matrix.

- For $\eta_0 \leq N^{-1-\delta}$ we can use the classical smoothing result [**Sankar-Spielman-Teng, '05**] to exclude small eigenvalues λ_i^z of H^z :

$$\mathbb{P} \left(\min_i |\lambda_i^z| \leq N^{-1-\delta} \right) \leq N^{-c\delta}.$$

- How to bridge the gap $\eta \in [N^{-1-\delta}, N^{-3/4-\delta}]$?**
- Remark:** If four moments matching is assumed [**Tao-Vu, '12**], the regime $\eta \in [N^{-1-\delta}, N^{-3/4-\delta}]$ can be covered by local law. \implies **No gap.**

Sketch of the proof

Without four moment assumption, to cover the regime $\eta \in [N^{-1-\delta}, N^{-3/4-\delta}]$, we need a completely new proof:

Theorem (C.-Erdős-Schröder, '19) Let $|1 - |z|^2| \lesssim N^{-1/2}$. For the tail of the smallest eigenvalue of H^z , we have

$$\mathbb{P} \left(\min_i |\lambda_i^z| \leq N^{-3/4-\delta} \right) \leq N^{-c\delta},$$

where λ_i^z are the eigenvalues of H^z .

Estimate of the lower tail of $\min_i |\lambda_i^z|$.

- We use a supersymmetric representation

$$\mathrm{Tr}(H - w)^{-1} = i \int_{\mathbb{C}^N} \partial_{\chi} \langle \chi, \chi \rangle e^{-i\mathrm{Tr}(H-w)\Phi\Phi^*} d\mathbf{s}, \quad \Phi := (\mathbf{s}, \chi), \quad \Phi\Phi^* = \mathbf{s}\mathbf{s}^* + \chi\chi^*$$

where χ are **Grassmannian** vectors.

Estimate of the lower tail of $\min_i |\lambda_i^z|$.

- We use a supersymmetric representation

$$\mathrm{Tr}(H - w)^{-1} = i \int_{\mathbb{C}^N} \partial_{\chi} \langle \chi, \chi \rangle e^{-i\mathrm{Tr}(H-w)\Phi\Phi^*} dS, \quad \Phi := (s, \chi), \quad \Phi\Phi^* = ss^* + \chi\chi^*$$

where χ are **Grassmannian** vectors.

- Performing the expectation,

$$\mathbb{E} e^{-i\mathrm{Tr}(X-z)(X^* - \bar{z})\Phi\Phi^*} = \mathrm{sdet} \left(1 + i \frac{\Phi^* \Phi}{N} \right)^{-N} \exp \left(-N |z|^2 \mathrm{sTr} \left(1 + \frac{i}{N} \Phi^* \Phi \right)^{-1} \frac{i}{N} \Phi^* \Phi \right),$$

$$\mathrm{sdet} \begin{pmatrix} x & \sigma \\ \tau & y \end{pmatrix} = \frac{x^2}{xy - \tau\sigma}, \quad \mathrm{sTr} \begin{pmatrix} x & \sigma \\ \tau & y \end{pmatrix} = x - y.$$

Estimate of the lower tail of $\min_i |\lambda_i^z|$.

- We use a supersymmetric representation

$$\text{Tr}(H - w)^{-1} = i \int_{\mathbb{C}^N} \partial_{\chi} \langle \chi, \chi \rangle e^{-i\text{Tr}(H-w)\Phi\Phi^*} ds, \quad \Phi := (s, \chi), \quad \Phi\Phi^* = ss^* + \chi\chi^*$$

where χ are **Grassmannian** vectors.

- Performing the expectation,

$$\mathbb{E} e^{-i\text{Tr}(X-z)(X^*-\bar{z})\Phi\Phi^*} = \text{sdet} \left(1 + i \frac{\Phi^*\Phi}{N} \right)^{-N} \exp \left(-N |z|^2 \text{sTr} \left(1 + \frac{i}{N} \Phi^*\Phi \right)^{-1} \frac{i}{N} \Phi^*\Phi \right),$$

$$\text{sdet} \begin{pmatrix} x & \sigma \\ \tau & y \end{pmatrix} = \frac{x^2}{xy - \tau\sigma}, \quad \text{sTr} \begin{pmatrix} x & \sigma \\ \tau & y \end{pmatrix} = x - y.$$

- Using the **superbosonization formula** [Bunder-Efetov-Kravtsov-Yevtushenko-Zirnbauer,'07;Littellmann-Sommers-Zirnbauer,'08]:

$$\int f(\Phi^*\Phi) = \frac{1}{2\pi i} \int_0^\infty dx \oint dy \partial_\tau \partial_\sigma \text{sdet}^N(Q) f(Q), \quad Q := \begin{pmatrix} x & \sigma \\ \tau & y \end{pmatrix},$$

to reduce the number of integrations to 2:

$$\mathbb{E} \text{Tr}((X - z)(X - z)^* - w)^{-1} = \frac{N^2}{2\pi i} \int_0^\infty dx \oint dy e^{-Nf(x) + Nf(y)} y \cdot G(x, y)$$

$$G(x, y) := \frac{1}{xy} - \frac{1}{(1+x)(1+y)} \left[1 + \frac{|z|^2}{1+x} + \frac{|z|^2}{1+y} \right], \quad f(x) := \log \frac{1+x}{x} - \frac{|z|^2}{1+x} - wx,$$

Estimate of the lower tail of $\min_i |\lambda_i^z|$.

- We use a supersymmetric representation

$$\text{Tr}(H - w)^{-1} = i \int_{\mathbb{C}^N} \partial_{\chi} \langle \chi, \chi \rangle e^{-i\text{Tr}(H-w)\Phi\Phi^*} d\mathbf{s}, \quad \Phi := (\mathbf{s}, \chi), \quad \Phi\Phi^* = \mathbf{s}\mathbf{s}^* + \chi\chi^*$$

where χ are **Grassmannian** vectors.

- Performing the expectation,

$$\mathbb{E} e^{-i\text{Tr}(X-z)(X^* - \bar{z})\Phi\Phi^*} = \text{sdet} \left(1 + i \frac{\Phi^* \Phi}{N} \right)^{-N} \exp \left(-N |z|^2 \text{sTr} \left(1 + \frac{i}{N} \Phi^* \Phi \right)^{-1} \frac{i}{N} \Phi^* \Phi \right),$$
$$\text{sdet} \begin{pmatrix} x & \sigma \\ \tau & y \end{pmatrix} = \frac{x^2}{xy - \tau\sigma}, \quad \text{sTr} \begin{pmatrix} x & \sigma \\ \tau & y \end{pmatrix} = x - y.$$

- Using the **superbosonization formula** [Bunder-Efetov-Kravtsov-Yevtushenko-Zirnbauer,'07;Littellmann-Sommers-Zirnbauer,'08]:

$$\int f(\Phi^* \Phi) = \frac{1}{2\pi i} \int_0^\infty dx \oint dy \partial_\tau \partial_\sigma \text{sdet}^N(Q) f(Q), \quad Q := \begin{pmatrix} x & \sigma \\ \tau & y \end{pmatrix},$$

to reduce the number of integrations to 2:

$$\mathbb{E} \text{Tr}((X - z)(X - z)^* - w)^{-1} = \frac{N^2}{2\pi i} \int_0^\infty dx \oint dy e^{-Nf(x) + Nf(y)} y \cdot G(x, y)$$

$$G(x, y) := \frac{1}{xy} - \frac{1}{(1+x)(1+y)} \left[1 + \frac{|z|^2}{1+x} + \frac{|z|^2}{1+y} \right], \quad f(x) := \log \frac{1+x}{x} - \frac{|z|^2}{1+x} - wx,$$

- The real case is more complicated, but also gives explicit formula as a **3-dim. integral**

Thank you very much for your attention!