

# On branching random walks on periodic lattices

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# A random walk

Suppose there is one particle at a point  $v \in \mathbb{Z}^d$  at time  $t = 0$ . This particle will either move to a point  $u \neq v$ ,  $u \in \mathbb{Z}^d$  or will remain at  $v$  over a short period of time  $\delta t$ .

a probability of the transition  $v \rightarrow u$

$$p(v, u, \delta t) = a(v, u)\delta t + o(\delta t),$$

a probability of the transition  $v \rightarrow v$

$$p(v, v, \delta t) = 1 + a(v, v)\delta t + o(\delta t).$$

# The transition intensity 1

The value  $a(v, u)$  is called the transition intensity between  $v$  and  $u$ .

$$(i) \ a(v, u) \geq 0, \quad v \neq u;$$

$$(ii) \ a(v, v) < 0;$$

$$(iii) \ \sum_{u \in \mathbb{Z}^d} a(v, u) = 0;$$

# The transition intensity 2

Let  $g_1, \dots, g_d$  be a family of linearly independent (not necessarily orthogonal) vectors with integer coordinates. By a lattice we mean a set

$$\Gamma = \left\{ g \in \mathbb{Z}^d : g = \sum_{j=1}^d n_j g_j, n_j \in \mathbb{Z}, j = 1, \dots, d \right\}.$$

(iv)  $a(v, u) = a(u, v) = a(v + g, u + g), \quad \forall g \in \Gamma;$

(v)  $\sum_{u \in \mathbb{Z}^d} \|u\|^2 |a(v, u)| < \infty, \quad v \in \mathbb{Z}^d;$

(vi) the graph  $G = (\mathbb{Z}^d, \mathcal{E})$  with the vertex set  $\mathbb{Z}^d$  and edge set

$$\mathcal{E} = \{(v, u) : a(v, u) > 0, v, u \in \mathbb{Z}^d\}$$

is connected.

# A branching source

Suppose there is one particle at a point with branching source and it can't move anywhere from there. We assume that a particle can generate several descendants over a short period of time  $\delta t$ .

a probability of generating  $k \neq 1$  descendants

$$p_k = b_k \delta t + o(\delta t).$$

a probability of generating  $k = 1$  descendant

$$p_1 = 1 + b_1 \delta t + o(\delta t).$$

# The branching intensity

The value  $b_k$  is called a branching intensity into  $k$  descendants.

$$(1) b_k \geq 0, \quad k \neq 1;$$

$$(2) b_1 \leq 0;$$

$$(3) \sum_{k=0}^{+\infty} b_k = 0;$$

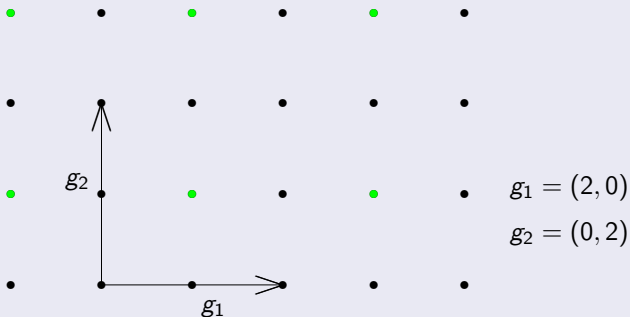
$$(4) \beta = \sum_{k=1}^{+\infty} kb_k < \infty;$$

# Locations of the branching sources

Let branching intensities depend on the point  $\nu$  and value  $\beta(\nu)$  be periodic with respect to the lattice  $\Gamma$ .

$$(5) \beta(\nu + g) = \beta(\nu), \quad g \in \Gamma.$$

## Example



# A random walk with a periodic set of branching sources

We assume that each new particle evolves according to the same law independently of other particles.

Each particle located at a point  $v \in \mathbb{Z}^d$  at time  $t$  can either move to a point  $u \neq v$  or remain at the source and produce  $k \neq 1$  descendants located at the point  $v$  (for  $k = 0$  we assume that the number of descendants is 0; that is, the particle dies) or remain unchanged (that is, no changes occur) over a short period of time  $[t; t + \delta t)$ .

a probability of transition  $v \rightarrow u$

$$p(v, u, \delta t) = a(v, u)\delta t + o(\delta t),$$

a probability of generating  $k \neq 1$  offsprings in  $v$

$$p_k(v, \delta t) = b_k(v)\delta t + o(\delta t).$$

a probability of remaining unchanged

$$p(v, \delta t) = 1 + a(v, v)\delta t + b_1(v)\delta t + o(\delta t).$$



# Finite number of branching sources

Yarovaya E., et. all [1998-2007]

One source. Asymptotic behavior of all moments. Limit theorems.

Vatutin V. Topchii V. [2005]

Limit theorem for catalytic BRW on  $\mathbb{Z}$  with one source of branching in super critical case.

Rytova A. Yarovaya E. [2018]

Extinction probability for sub and super critical cases.

Khristolyubov I. Yarovaya E. [2018-2019]

$N$  sources. Asymptotic behavior of all moments in subcritical and supercritical cases. Limit theorems.

# Mean value of particles

By  $M(v, u, t)$  we denote the mean number of particles at a point  $u$  at time  $t$ , provided that at the initial time  $t = 0$  there was one particle at a point  $v$ . The function  $M(v, u, t)$  satisfies the following Cauchy problem:

## The Cauchy problem

$$\begin{cases} M'_t(v, u, t) = \mathcal{A}M(v, u, t), \\ M(v, u, 0) = \delta_u(v). \end{cases}$$

## The operator $\mathcal{A}$

$$\begin{aligned} \mathcal{A} &= \mathcal{A}_0 + Q, \\ (\mathcal{A}_0 f)(v) &= \sum_{w \in \mathbb{Z}^d} a(v, w) f(w), \\ (Qf)(v) &= \beta(v) f(v). \end{aligned}$$

# Properties of $\mathcal{A}$

The operator  $\mathcal{A} : \ell^2(\mathbf{Z}^d) \rightarrow \ell^2(\mathbf{Z}^d)$  satisfies the following properties:

- $\mathcal{A}$  is self-adjoint;
- $\mathcal{A}$  is bounded;
- $\mathcal{A}_0$  is non-positive;
- $\mathcal{A}$  is periodic with respect to the lattice  $\Gamma$ .

## Connection with a discrete Laplacian

If condition (v) is replaced with stronger condition that for any  $v \in \mathbb{R}^d$  there are only finite number of transition probabilities  $a(v, u)$  are not zero, then operator  $-\mathcal{A}_0$  is a discrete combinatorial Laplacian on the graph  $G$  defined in (vi). In this case operator  $-\mathcal{A}$  is a discrete Schrödinger with periodic potential.

# A discrete Laplacian

Sunada T., Sy P. [1992]

Local spectral properties near the left edge of the spectrum.

Higuchi Y., Shirai T. [2001]

Local spectral properties near the left edge for the magnetic Schrödinger.

Higuchi Y., Shirai T. [2004]

Direct integral decomposition. Spectral properties.

Higuchi Y., Nomura Y. [2008]

Direct integral decomposition. Conditions for absolute continuity of the spectrum (absence of flat bands).

Korotyaev E., et. all [2010-2018]

Direct integral decomposition with explicit form of fiber operators.  
Estimates for spectrum of Laplacian with different types of perturbations.

# Asymptotic behaviour of $M(v, u, t)$

## Theorem

The function  $M(v, u, t)$  has the following asymptotic behaviour as  $t \rightarrow \infty$

$$M(v, u, t) = e^{\lambda_1(0)t} t^{-\frac{d}{2}} c_0(v, u) \left(1 + O(t^{-1})\right)$$

where

$$c_0(v, u) = \frac{|C|}{(2\pi)^{\frac{d}{2}}} \frac{\psi_1(v', 0) \psi_1(u', 0)}{\sqrt{\left| \det \left\{ \frac{\partial^2 \lambda_1(\theta)}{\partial \theta^2} \Big|_{\theta=0} \right\} \right|}}$$

where  $v = v' + \gamma_v$ ,  $u = u' + \gamma_u$ ,  $v', u' \in \Omega$ ,  $\gamma_v, \gamma_u \in \Gamma$ .

If  $a(v, u)$  decays fast enough, then

$$M(v, u, t) \stackrel{as}{=} e^{\lambda_1(0)t} t^{-\frac{d}{2}} \sum_{k=0}^{\infty} c_k(v, u) t^{-k}.$$

# Notations

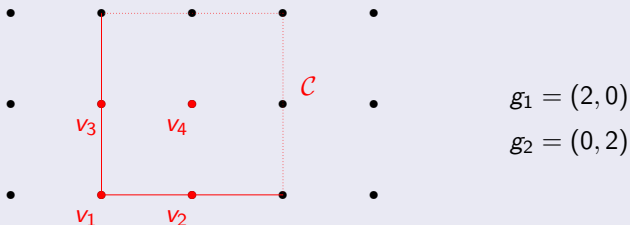
## Fundamental cell

$$\mathcal{C} = \{x \in \mathbf{R}^d : x = \sum_{j=1}^d x_j g_j, 0 \leq x_j < 1, j = 1, \dots, d\}.$$

## Fundamental vertex set

$$\Omega = \mathcal{C} \cap \mathbb{Z}^d, \quad \Omega = \{v_1, \dots, v_p\}.$$

## Example



# An auxiliary matrix family

We introduce a new variable  $\theta$  with range such that  $2\pi\theta \in \mathcal{C}$ .

$$A(\theta) = \begin{pmatrix} \tilde{a}_{11}(\theta) + \beta_1 & \tilde{a}_{12}(\theta) & \cdots & \tilde{a}_{1p}(\theta) \\ \tilde{a}_{21}(\theta) & \tilde{a}_{22}(\theta) + \beta_2 & \cdots & \tilde{a}_{2p}(\theta) \\ \vdots & \vdots & \ddots & \vdots \\ \tilde{a}_{p1}(\theta) & \tilde{a}_{p2}(\theta) & \cdots & \tilde{a}_{pp}(\theta) + \beta_p \end{pmatrix},$$

where the functions  $\tilde{a}_{jk}(\theta)$  and constants  $\beta_j$  is defined by

$$\tilde{a}_{jk}(\theta) = \sum_{g \in \Gamma} e^{-i\langle g, \theta \rangle} a(v_j + g, v_k), \quad \beta_j = \beta(v_j).$$

## Connection between $\mathcal{A}$ and $A(\theta)$

Let the eigenvalues of the matrix family  $A(\theta)$  be ordered in non-increasing order for every parameter  $\theta$ :  $\lambda_1(\theta) \geq \dots \geq \lambda_p(\theta)$ .

$$\sigma(\mathcal{A}) = \bigcup_{j=1}^p \bigcup_{2\pi\theta \in \mathcal{C}} \lambda_j(\theta),$$

# Properties of $\lambda_1(\theta)$

## Theorem

For  $\lambda_1(\theta)$  the following statements hold:

- a) the right edge of the spectrum of  $\mathcal{A}$  coincides with  $\lambda_1(0)$ , i.e.

$$\max \sigma(\mathcal{A}) = \lambda_1(0);$$

- b) the distance between the right edge of the spectrum of  $\mathcal{A}$  and the right edge of the second spectral band is positive, i.e.

$$\lambda_1(0) - \sup_{\theta \in \tilde{\mathcal{C}}} \lambda_2(\theta) > 0;$$

- c) the determinant of the Hessian matrix of  $\lambda_1(\theta)$  does not vanish at  $\theta = 0$ , i.e.

$$\det \left\{ \frac{\partial^2 \lambda_1(\theta)}{\partial \theta^2} \Big|_{\theta=0} \right\} \neq 0;$$

- d)  $\lambda_1(0)$  is not an eigenvalue of  $\mathcal{A}$ .



# Properties of $\lambda_1(0)$

## Theorem

$\lambda_1(0)$  is positive if either  $\sum_{j=1}^p \beta_j > 0$  or for some  $m = 1, \dots, p$   
 $\beta_m \geq \|A_0(0)\|$ .

## Theorem

If we add the coupling constant  $\mu$  then

- for  $\mu \rightarrow 0$

$$\lambda_1(0) = \frac{\mu}{p} \sum_{j=1}^p \beta_j + O(\mu^2).$$

- for  $\mu \rightarrow \infty$

$$\lambda_1(0) = \mu \max_{j=1, \dots, p} \beta_j + O(1).$$

# The second moment (supercritical case)

Let  $M_2(v, u, t)$  be the second moment for our BRW. Assume that additionally to (4) and (5) following conditions are satisfied:

$$(4') \beta_2(v) = \sum_{k=2}^{+\infty} k(k-1)b_k(v) < \infty;$$

$$(5') \beta_2(v+g) = \beta_2(v), \quad g \in \Gamma.$$

The Cauchy problem for the second moment

$$\begin{cases} M_{2t}'(v, u, t) &= (\mathcal{A}M_2)(v, u, t) + \beta_2(v)M^2(v, u, t), \\ M_2(v, u, 0) &= \delta_u(v), \end{cases}$$

# Asymptotic behaviour of $M_2(v, u, t)$

## Theorem

Suppose that  $\lambda_1(0) > 0$ . Then the function  $M_2(v, u, t)$  has the following asymptotic behavior as  $t \rightarrow \infty$

$$M_2(v, u, t) = \frac{e^{2\lambda_1(0)t} t^{-d} d_0(v, u)}{\lambda_1(0)} (1 + O(t^{-1})),$$

where the  $d_0(v, u)$  is defined by

$$d_0(v, u) = \frac{|C|^2 \psi_1(v', 0) |\psi_1(u', 0)|^2}{2^{d/2} (2\pi)^d \left| \det \left\{ \frac{\partial^2 \lambda_1(\theta)}{\partial \theta^2} \Big|_{\theta=0} \right\} \right|} \sum_{w \in \Omega} \psi_1(w, 0) \beta_2(w) |\psi_1(w, 0)|^2$$

- Higuchi Y., Nomura Y. Spectral structure of the Laplacian on a covering graph //European Journal of Combinatorics. – 2009. – V. 30.2. – P. 570-585.
- Higuchi Y., Shirai T. Some spectral and geometric properties for infinite graphs //Contemporary Mathematics. – 2004. – V. 347. – P. 29-56.
- Khristolyubov I., Yarovaya E. A Limit Theorem for Supercritical Branching Random Walks with Branching Sources of Varying Intensity //arXiv preprint arXiv:1904.01468. – 2019.
- Korotyaev E., Saburova N. Schrödinger operators on periodic discrete graphs //Journal of Mathematical Analysis and Applications. – 2014. – V. 420.1. – P. 576-611.
- Korotyaev E., Saburova N. Schrödinger operators with guided potentials on periodic graphs //Proceedings of the American Mathematical Society. – 2017. – V. 145.11. – P. 4869-4883.

- Higuchi Y., Shirai T. Weak Bloch property for discrete magnetic Schrödinger operators //Nagoya Mathematical Journal. – 2001. – V. 161. – P. 127-154.
- Rytova A., Yarovaya E. Survival Analysis of Particle Populations in Branching Random Walks //arXiv preprint arXiv:1812.09909. – 2018.
- Sy P. W., Sunada T. Discrete Schrödinger operators on a graph //Nagoya Mathematical Journal. – 1992. – V. 125. – P. 141-150.
- Vatutin V. A., Topchii V. A. Limit theorem for critical catalytic branching random walks //Theory of Probability and Its Applications. – 2005. – V.3. – P. 498-518.
- Yarovaya E. B. A limit theorem for critical branching of a random walk on with a single source //Russian Mathematical Surveys. – 2005. – V. 60.1. – P. 173.

Thank you for your attention!