

Wilson Flow and Applications

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- 3 Wilson Flow and Renormalisation
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Motivation

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Many applications: renormalisation, separation of topological sectors...

Flow equation

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- A continuous map $\mathcal{F}: V \rightarrow U$ can be thought of as a diff. eq. evolving U through an extra dimension t .
- One can obtain this differential equation by considering infinitesimal transformations:

$$U \rightarrow U + \epsilon Z(U)U + O(\epsilon^2),$$

whose continuous composition is given by

$$\dot{U}_t = Z_t(U_t)U_t,$$

and Z lives in the algebra.

Flow equation II

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- Since $\ln \det \mathcal{F}_{t,*}(V)$ is equal to the LHS of the above eq.:

$$S(\mathcal{F}_t(V)) - \ln \det \mathcal{F}_{t,*}(V) = (1-t)S(\mathcal{F}_t(V)) - C_t,$$

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- and we see how for $t = 1$ we obtain a trivialising map.

Solutions for the flow generator I

A better equation to find Z_t is:

$$\sum_{x,\mu} [\partial_{x,\mu}^a Z^a(U_s) - t \partial_{x,\mu}^a S(U) Z^a(U_s)] = S(U) + \dot{C}_t.$$

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Why the choice of the Wilson action? \implies simplicity

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Thus, to first order the adequate flow generator is just the gradient of the Wilson action!

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- Locality guaranteed where series converges. But the latter is not clear.
- At least to leading order, the use of the Wilson flow decreases the action for increasing t :

$$\frac{d}{dt} S_w(U_t) = -\frac{3}{16} \sum_{x,\mu} \partial_{x,\mu}^a S_w(U_t) \partial_{x,\mu}^a S_w(U_t) \leq 0.$$

Wilson Flow and Renormalisation

Wilson flow and perturbation theory

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and inserting that in the above eq. we obtain a series of equations. To leading order, we have that:

$$B_{\mu,1}(t, x) = \int d^4y \frac{e^{-(x-y)^2/4t}}{(4\pi t)^2} A_\mu(y).$$

Explicitly, we see that to this order the gauge potential is averaged over a volume whose mean-square radius is $\sqrt{8t}$.

$\langle E \rangle$ in perturbation theory

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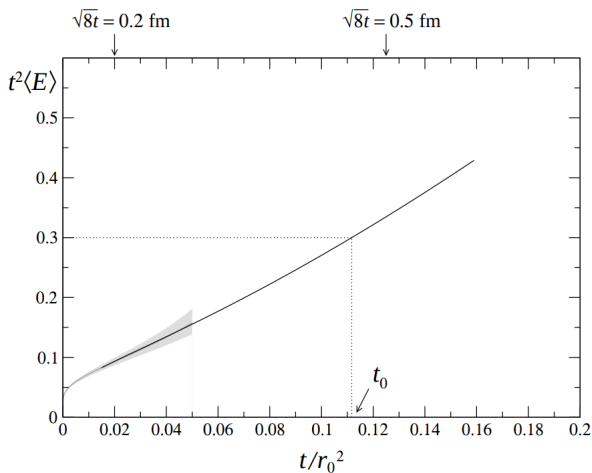
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It does from the next order and onward. Using the relation between the bare and renormalised couplings in the $\overline{\text{MS}}$ scheme and going to an energy scale $q = (8t)^{-1/2}$:

$$\langle E \rangle = \frac{3}{4\pi t^2} \alpha(q) [1 + k_1 \alpha(q) + O(\alpha^2)].$$

Flow dependence of $t^2\langle E \rangle$ 

© M. Lusher 2010

What about fermions?

We can also flow the fermionic fields $\not\in$ [M. Lusher 2013](#). The simplest extension is to keep the evolution of the gauge fields untouched and:

$$\partial_t \chi_t = D_\mu D_\mu \chi_t, \quad \chi_t|_{t=0} = \psi.$$

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Also, similarly to the gauge field case, now the energy scale of the renormalised operators can be understood as driven by the change in t .

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Recall that for $t > 0$, the Wilson action tends to decrease!

Wilson Flow and Topology II

We can define a way to measure the smoothness of a certain gauge field as:

$$h = \max \{s_p\}, \quad s_p = \text{Re tr}[1 - V(p)].$$

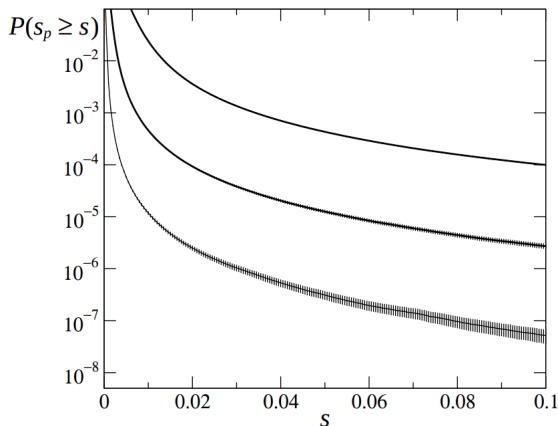
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Large values for h not favoured \implies large values for the plaquettes are suppressed.

Suppression of the plaquettes



Top to bottom $a = 0.1, 0.07, 0.05$ fm *M. Lusher 2010*

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Long story short, fields satisfying:

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Hence, Wilson flow dynamically drives the gauge fields into smooth and well-defined topological configurations.

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We have talked about:

- How the flow equation can be obtained.
- That Wilson flow drives fields to configurations decreasing the Wilson action.
- That the gauge and fermionic fields renormalise at positive flow time.
- An intuition for how disconnected topological sectors appear at $t > 0$.

Thank you for your attention!