

From Tensor Integral to IBP

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in collaboration with
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Motivation

NNLO

Tensor Integral

Tarasov's method

Projection method

Applications

Heavy Quark Form Factors

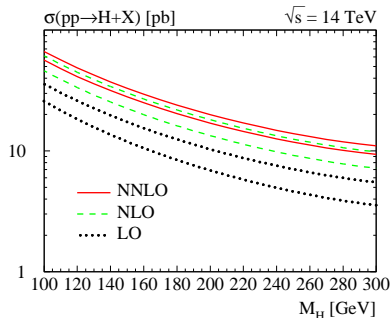
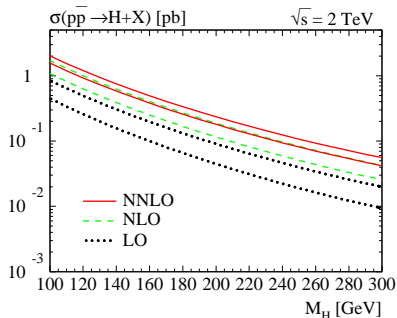
Single Top Quark Production

Conclusions

Multi loop calculation

- *More precision in calculated results*

Ex.: Total cross section for Higgs production in gluon fusion



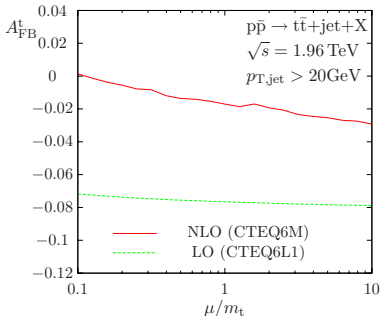
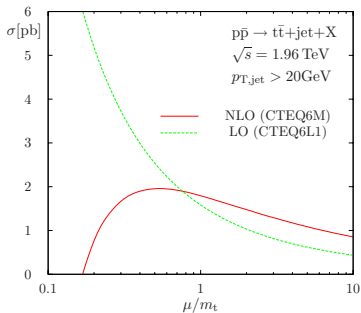
[R. Harlander, W. Kilgore Nov. '02]

- *Perturbative convergence $LO \rightarrow NLO(\approx 70\%)$ and $NLO \rightarrow NNLO(\approx 30\%)$*

Multi loop calculation

- *More precision in calculated results*
- *New effects*

Ex.: forward-backward charge asymmetry of the top quark



[S. Dittmaier, P. Uwer, S. Weinzierl Apr. '08]

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Backgrounds for New Physics Searches

When do we need NNLO?

- *When NLO corrections are large*

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*For instance **single top quark** production :*

Single top quark production

<i>Process</i>	\sqrt{S}	$\sigma_{LO}(pb)$	$\sigma_{NLO}(pb)$
<i>t-channel</i>	2.0 TeV $p\bar{p}$	1.068	1.062
	14.0 TeV pp	152.7	155.9

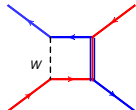
[B.Harris, E. Laenen, L.Phaf, Z. Sullivan, S. Weinzierl '02]

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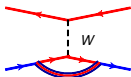
No colour exchange at NLO:



$$\times \left(\begin{array}{c} \text{quark} \\ \text{---} \\ \text{W} \\ \text{---} \\ \text{quark} \end{array} \right)^*$$

$$\propto \text{tr}[T_a] \text{tr}[T_a] = 0$$

Only vertex corrections contribute:

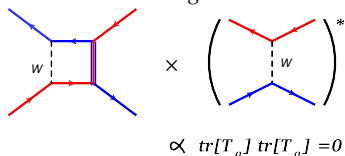


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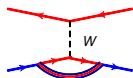
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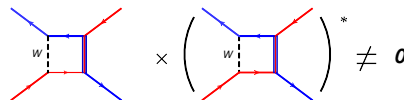


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Only vertex corrections contribute:



Colour exchange at NNLO:



$$\neq 0$$



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- *The true uncertainty due to missing higher order correction may be greater because new colour exchange diagrams first contribute at NNLO.*
- *A significant effect on kinematical distributions*
- *Important for studies of the V–A structure*
- *Source of polarised top quarks*
- *Access to the b quark pdfs*
- ...

In the NNLO-corrections occur tensor integrals:

$$\mathcal{I}(d, a_1, \dots, a_n)_{[1, k_1^\mu, k_2^\nu, \dots]} = \int d^d k_1 \int d^d k_2 \frac{\prod_{ij} k_1^{\mu_i} k_2^{\nu_j}}{p_1^{a_1} \dots p_n^{a_n}}$$

Possibilities to reduce tensor integrals to scalar integrals:

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Possibilities to reduce tensor integrals to scalar integrals:

- *By Schwinger parametrization*

[O. V. Tarasov, Phys. Rev. '96, Nucl. Phys. '81]

- *By projection method*

[T. Binoth, E.W.N. Glover, P. Marquard and J.J. van der Bij '02; E.W.N. Glover '04]

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Tensor reduction \Rightarrow various scalar integrals with the same structure of the integrand with different powers of propagators

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Reduction techniques:

- *Laporta: efficient algorithm to solve linear system of IBP-Identities*

AIR *[Anastasiou, Lazopoulos '04]*

FIRE *[Smirnov '08]*

Crusher *[Marquard, Seidel (to be published)]*

REDUZE 1&2 *[Studerus '09; Manteuffel, Studerus '12]*

Tensor reduction leads to a very large number of scalar integrals which are shifted in dimension and have other powers of propagators

$$\mathcal{J}(d, a_1, \dots, a_n) [k_1^\mu k_2^\nu, \dots] \rightarrow g^{\mu\nu} \sum_i \mathcal{J}(d + x_i, a_1^i, \dots, a_n^i) [1]$$

Example for two loop corrections to Axial Vector Form Factors

$$\begin{aligned} \mathcal{J}(d, 1, 1, 1, 1, 1, 1) [1, k_1^{\mu_1} k_1^{\mu_2} k_2^{\nu_1} k_2^{\nu_2}] &\rightarrow \mathcal{J}(2 + d, 2, 1, 1, 1, 1, 2) + \\ &\dots + \mathcal{J}(4 + d, 1, 1, 1, 2, 3, 1) + \dots + \mathcal{J}(8 + d, 3, 3, 3, 2, 1, 2) \end{aligned}$$

Shift in the dimension

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An arbitrary scalar Feynman integral:

$$\mathcal{I}^{(d)}(\{s_i\}, \{m_s^2\}) \propto \prod_{j=1}^N c_j \int_0^\infty \cdots \int_0^\infty \frac{d\alpha_j \alpha_j^{a_j-1}}{[D(\alpha)]^{\frac{d}{2}}} e^{i \left[\frac{Q(\{s_i\}, \alpha)}{D(\alpha)} - \sum_{l=1}^N \alpha_l (m_l^2 - i\epsilon) \right]}$$

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$D\left(\frac{\partial}{\partial m_j^2}\right)$ (polynomial differential operator) obtained from $D(\alpha)$ by substituting $\alpha_i \rightarrow \partial_j \equiv \partial/\partial m_j^2$. The application of $D(\partial_i)$ to the scalar integral:

$$\mathcal{I}^{(d-2)}(\{s_i\}, \{m_s^2\}) \propto D(\partial_j) \mathcal{I}^{(d)}(\{s_i\}, \{m_s^2\}),$$

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$$\mathcal{I}^{(d-2)}(\{s_i\}, \{m_s^2\}) \propto D(\partial_j) \mathcal{I}^{(d)}(\{s_i\}, \{m_s^2\}),$$

apply this to master integrals

$$\mathcal{I}_{master}(d-2, a_1, \dots, a_n) = \sum_i c_i \mathcal{I}(d, a_1^i, \dots, a_n^i),$$

all scalar integrals in rhs. of that equation have to be replaced by master integrals.

i.e.

$$\mathcal{I}_{master}(d-2, a_1, \dots, a_n) = \sum_j D_j \mathcal{I}_{master}(d, a_1^j, \dots, a_n^j),$$

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$$\mathcal{I}_{master}(d-2, a_1, \dots, a_n) = \sum_j D_j \mathcal{I}_{master}(d, a_1^j, \dots, a_n^j),$$

we have for all master integrals:

$$\begin{pmatrix} \mathcal{I}_1^{d-2} \\ \vdots \\ \mathcal{I}_l^{d-2} \end{pmatrix}_{master} = D_{ll} \cdot \begin{pmatrix} \mathcal{I}_1^d \\ \vdots \\ \mathcal{I}_l^d \end{pmatrix}_{master}$$

where l is the number of master integrals.

By this method we get scalar products between loop momenta and external momenta and no shift in the dimension of integrals

$$\mathcal{I}(d, a_1, \dots, a_n)_{[1, k_1^\mu, k_2^\mu, \dots]} \rightarrow g^{\mu\nu} \sum_{ij} \mathcal{I}(d, a_1, \dots, a_n)_{[1]} k_i p_j$$

Example for two loop corrections to Axial Vector Form Factors

$$\begin{aligned} \mathcal{I}(d, 1, 1, 1, 1, 1, 1)_{[1, k_1^{\mu 1} k_1^{\mu 2} k_2^{\nu 1} k_2^{\nu 2}]} &\rightarrow \mathcal{I}(d, -2, 1, 1, 1, 1, 1) + \dots \\ &+ \mathcal{I}(d, -1, 0, 1, 1, 1, 1) + \dots + \mathcal{I}(d, 0, 1, 1, 1, -2, -2) \end{aligned}$$

The general tensor structure for the amplitude \mathcal{A} :

$$\mathcal{A} = \sum_{i=1}^n B_i(t, u, s) \mathcal{S}_i,$$

where t, u and s are the Mandelstam variables and \mathcal{S}_i are the Dirac structures.

Projectors for the tensor coefficients:

$$\mathcal{S}_j^\dagger \mathcal{A} = \sum_{i=1}^n B_i(t, u, s) \underbrace{(\mathcal{S}_j^\dagger \mathcal{S}_i)}_{\mathcal{M}_{ji}} \Rightarrow B_i(t, u, s) = \sum_j \mathcal{M}_{ij}^{-1} (\mathcal{S}_j^\dagger \mathcal{A})$$

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essential for this method : to be able to calculate the inverse matrix \mathcal{M}_{ij}^{-1}

Tarasov's method

- *Positive powers for propagators (the sum of the powers of all propagators is large)*
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Projection method

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Tarasov's method

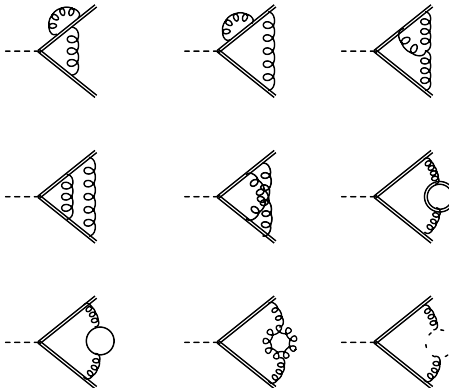
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Two loop corrections to Heavy Quark Form Factors

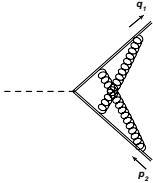
We implemented both methods to calculate the two loop corrections to Heavy Quark Vector and axial Vector Form Factors:

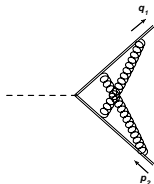


[W. Bernreuther, R. Bonciani, T. Gehrmann, R. Heinesch, T. Leineweber, P. Mastrolia, E. Remiddi '04]

[J. Gluza, A. Mitov, S. Moch, T. Riemann '09]

Two loop corrections to Heavy Quark Form Factors





There are 6 Dirac structures

(Heavy Quark Vector and axial Vector

Form Factors):

$$\mathcal{S}_1 = \bar{u}(q_1)(1 + \gamma_5)u(p_2)p_2^\mu$$

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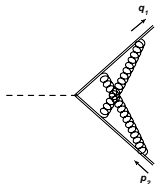
$$\mathcal{S}_3 = \bar{u}(q_1)(1 + \gamma_5)u(p_2)q_1^\mu$$

$$\mathcal{S}_4 = \bar{u}(q_1)(1 - \gamma_5)u(p_2)q_1^\mu$$

$$\mathcal{S}_5 = \bar{u}(q_1)(1 + \gamma_5)\gamma_\mu u(p_2)$$

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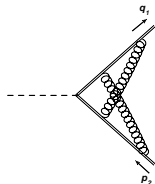
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number of integrals	564	671
max sum of powers of propagators	6	14
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reduction time	7500 s	433260 s

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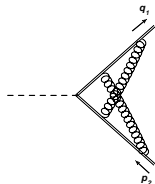
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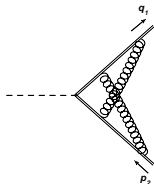
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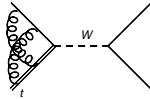
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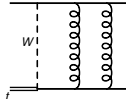
but ...

Two loop corrections to single Top Quark Production

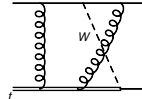
There are three topological families:



Vertex corrections



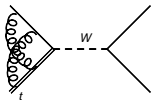
Planar double boxes



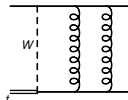
Non Planar double boxes

Two loop corrections to single Top Quark Production

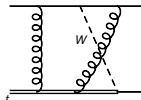
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Non Planar double boxes

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$$S_3 = \bar{u}(q_1) \gamma_6 \gamma_{\mu_1} u(p_2) \bar{u}(q_2) \gamma_6 \gamma_{\mu_1} u(p_1)$$

$$S_4 = \bar{u}(q_1) \gamma_7 \gamma_{\mu_1} \gamma_{p_1} u(p_2) \bar{u}(q_2) \gamma_6 \gamma_{\mu_1} u(p_1)$$

$$S_5 = \bar{u}(q_1) \gamma_7 \gamma_{\mu_1} \gamma_{\mu_2} u(p_2) \bar{u}(q_2) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{q_1} u(p_1)$$

$$S_6 = \bar{u}(q_1) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{p_1} u(p_2) \bar{u}(q_2) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{q_1} u(p_1)$$

$$S_7 = \bar{u}(q_1) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} u(p_2) \bar{u}(q_2) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} u(p_1)$$

$$S_8 = \bar{u}(q_1) \gamma_7 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} \gamma_{p_1} u(p_2) \bar{u}(q_2) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} u(p_1)$$

$$S_9 = \bar{u}(q_1) \gamma_7 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} \gamma_{\mu_4} u(p_2) \bar{u}(q_2) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} \gamma_{\mu_4} \gamma_{q_1} u(p_1)$$

$$S_{10} = \bar{u}(q_1) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} \gamma_{\mu_4} \gamma_{p_1} u(p_2) \bar{u}(q_2) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} \gamma_{\mu_4} \gamma_{q_1} u(p_1)$$

$$S_{11} = \bar{u}(q_1) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} \gamma_{\mu_4} \gamma_{\mu_5} u(p_2) \bar{u}(q_2) \gamma_6 \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} \gamma_{\mu_4} \gamma_{\mu_5} u(p_1)$$

- *Vertex corrections: both methods*

$$\begin{aligned}
 & \text{Diagrammatic representation of vertex corrections: a central vertex with two external lines, surrounded by various loop and self-energy corrections.} \\
 & = \left[N_C^2 T C_F g_s^4 \left\{ S_1 \left[\frac{1}{m_t^2} \left(\frac{-8-11t+t^2}{2\epsilon(t-1)^3} - \frac{71-275t+44t^2}{12(t-1)^3} + \mathcal{O}(\epsilon) \right) \right] \right. \right. \\
 & \quad - \left(\frac{2n_f}{3m_t^2(t-1)} + \mathcal{O}(\epsilon) \right) \left. \right] \text{Diagram} + \frac{1}{m_t^2} \left(\frac{-8-11t+t^2}{3\epsilon(t-1)^3} + \frac{-103-226t+41t^2}{18(t-1)^3} + \mathcal{O}(\epsilon) \right) \text{Diagram} \\
 & \quad + \left(\frac{2n_f}{3m_t^2(t-1)} + \mathcal{O}(\epsilon) \right) \text{Diagram} + \left(\frac{12t}{m_t^2(t-1)^3} + \mathcal{O}(\epsilon) \right) \text{Diagram} - \left(\frac{2m_t(t+1)}{3(t-1)} + \mathcal{O}(\epsilon) \right) \text{Diagram} \\
 & \quad + S_3 \left[\frac{1}{m_t^4} \left(-\frac{11+12t+13t^2}{24\epsilon(t-1)^2} + \frac{-72+169t-324t^2+303t^3+76t^4-24t^5}{72(t-1)^3t} + \mathcal{O}(\epsilon) \right) \right] \text{Diagram} \\
 & \quad + \frac{1}{m_t^2} \left(\frac{2}{3\epsilon} - \frac{46+28t}{9t} + \mathcal{O}(\epsilon) \right) \text{Diagram} + \frac{1}{m_t^2} \left(\frac{2t}{3\epsilon(t-1)} + \frac{2t(1+5t-14t^2)}{9(t-1)^3} + \mathcal{O}(\epsilon) \right) \text{Diagram} \\
 & \quad + \frac{1}{m_t^2} \left(\frac{n_f(5-19t)}{36(t-1)} - \frac{n_f(t+1)}{6\epsilon(t-1)} + \mathcal{O}(\epsilon) \right) \text{Diagram} + \frac{1}{m_t^2} \left(\frac{5t^2+12t+19}{36\epsilon(t-1)^2} + \frac{29t^2-86t-87}{72(t-1)^2} + \mathcal{O}(\epsilon) \right) \text{Diagram} \\
 & \quad + (-n_f/9 + n_f/(3\epsilon) + \mathcal{O}(\epsilon)) \text{Diagram} + \left(\frac{n_f t}{3\epsilon(t-1)} - \frac{n_f(3+t)}{9(t-1)} + \mathcal{O}(\epsilon) \right) \text{Diagram} \\
 & \quad + (2/t - t/3 + \mathcal{O}(\epsilon)) \text{Diagram} + \left(\frac{(1+3t+3t^2-t^3)}{3(t-1)^2} + \mathcal{O}(\epsilon) \right) \text{Diagram} \\
 & \quad - \left(\frac{m_t^2(3t^2+19t+28)}{9} + \mathcal{O}(\epsilon) \right) \text{Diagram} - \left(\frac{m_t^2 t(3t^2+4t+1)}{9(t-1)} + \mathcal{O}(\epsilon) \right) \text{Diagram} \Big\} \\
 & \quad + N_C^2 C_F^2 \{ \dots \} + N_C^2 C_F C_A \{ \dots \} \Big] \frac{1}{m_t^2(t-m_W^2)}
 \end{aligned}$$

- *Vertex corrections: both methods* ✓

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try with common computer algebra system, e.g.

Mathematica or Maple

runtime ≈ 1 month with 64GB RAM

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- *We calculated \mathcal{M}_{ji}^{-1} and all Planar double boxes diagrams* ✓

- *Vertex corrections: both methods* ✓
- *Planar double boxes : projection method*

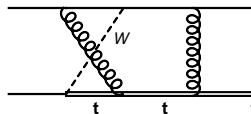
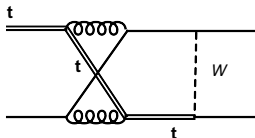
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- *We calculated \mathcal{M}_{ji}^{-1} and all Planar double boxes diagrams* ✓
- *Non Planar double boxes: a challenge !*

<i>Topology</i>	<i># Diagrams</i>	<i>reduction</i>	<i>performed checks</i>
<i>Vertex corrections</i>	29	✓	✓
<i>Planar double boxes</i>	6	✓	<i>work in progress</i>
<i>Non Planar double boxes</i>	12	<i>work in progress</i>	—

There are two most complicated topologies, which could not be reduced completely until now :



- *We have seen two possibilities to reduce tensor integrals to scalar integrals*
- *The choice of reductions method determines how difficult the next step (IBP) is*
- *As a test of our setup, we have calculated the $\mathcal{O}(\alpha_s^2)$ contributions to the Heavy Quark Vector and Axial Vector Form Factors, confirming the results of Bernreuther et al. and Gluza et al.*
- *We have also calculated the two loop vertex corrections to single Top Quark Production*