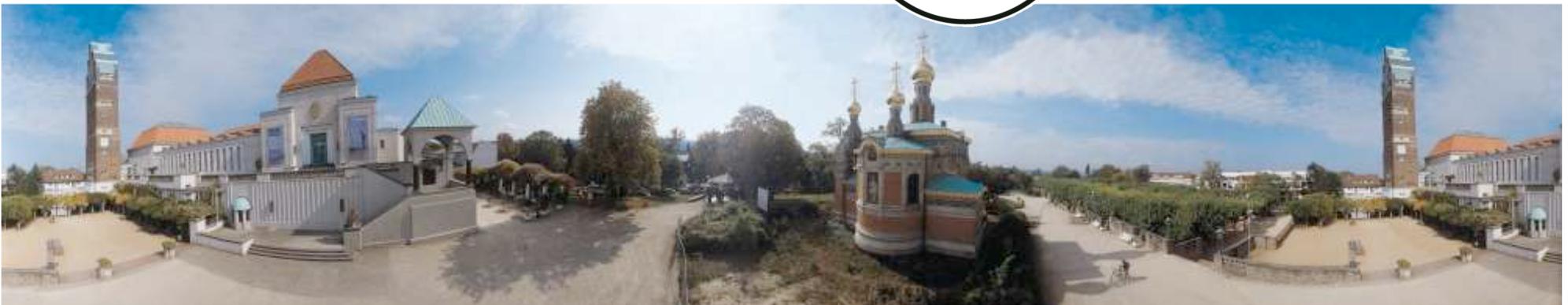


Chiral effective field theory and neutron-rich matter

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TECHNISCHE
UNIVERSITÄT
DARMSTADT



Block Course “Aspects of QCD at Finite Density”
University of Bielefeld, Sept. 22/23, 2011



DFG



*Minerva
Stiftung*



Bundesministerium
für Bildung
und Forschung



Outline

Chiral effective field theory for nuclear forces

Epelbaum, Hammer, Meissner, Rev. Mod. Phys. 81, 1773 (2009). Overview.

Epelbaum, Prog. Part. Nucl. Phys. 57, 654 (2006). Many useful details.

Bogner, Furnstahl, AS, Prog. Part. Nucl. Phys. 65, 94 (2010). Renormalization Group.

3N forces and neutron-rich nuclei

Otsuka, Suzuki, Holt, AS, Akaishi, Phys. Rev. Lett. 105, 032501 (2010).

Holt, Otsuka, AS, Suzuki, arXiv:1009.5984.

Holt, AS, arXiv:1108.2680.

3N forces and neutron matter, impact on neutron stars

Hebeler, AS, Phys. Rev. C 82, 014314 (2010).

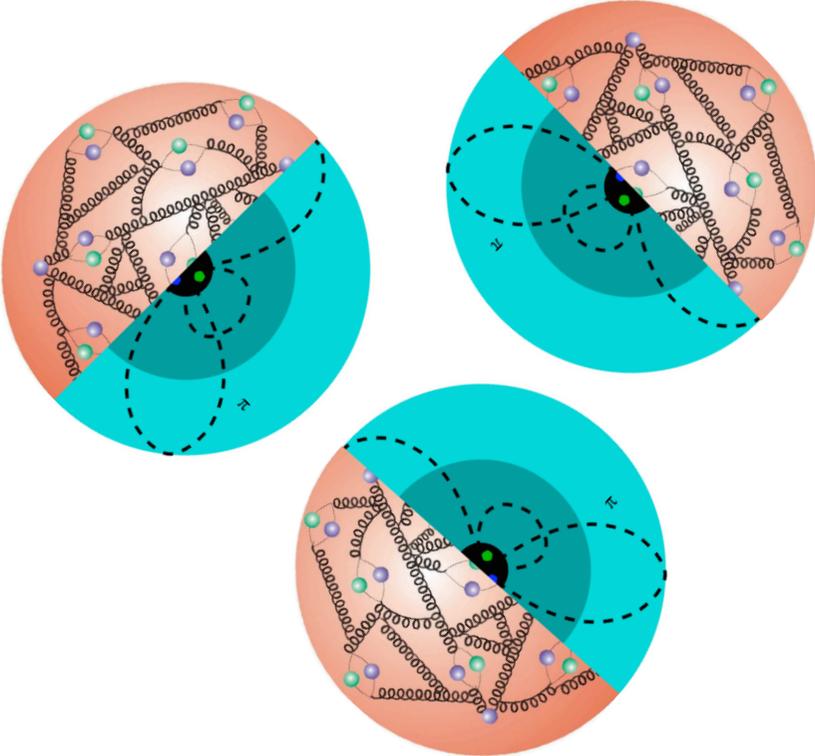
Hebeler, Lattimer, Pethick, AS, Phys. Rev. Lett. 105, 161102 (2010).

3N forces and electroweak currents

Menendez, Gazit, AS, Phys. Rev. Lett. 107, 062501 (2011).

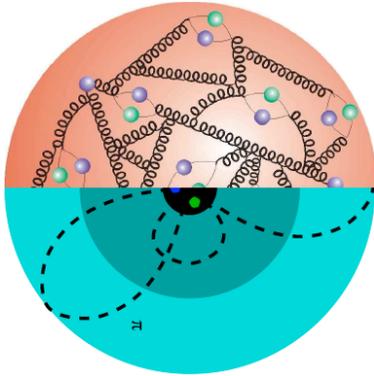
Contact: schwenk@physik.tu-darmstadt.de

The nuclear forces frontier

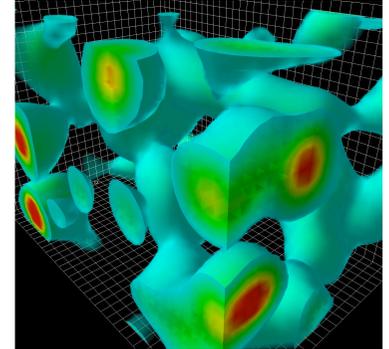


Λ / Resolution dependence

with high-energy probes:
quarks+gluons



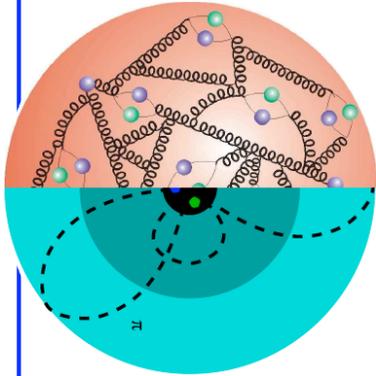
at low energies:
complex QCD vacuum



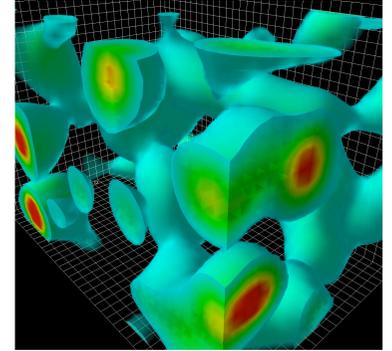
lowest energy excitations:
pions, nearly massless, $m_\pi=140$ MeV
'phonons' of QCD vacuum

Λ / Resolution dependence of nuclear forces

with high-energy probes:
quarks+gluons



at low energies:
complex QCD vacuum



lowest energy excitations:
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'phonons' of QCD vacuum

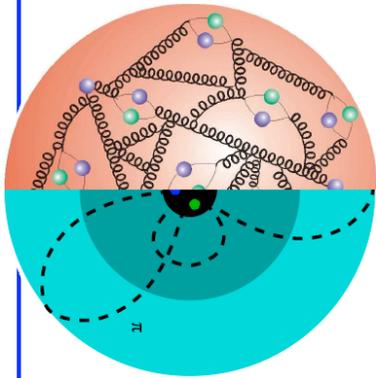
Λ_{chiral}

momenta $Q \sim \lambda^{-1} \sim m_\pi$

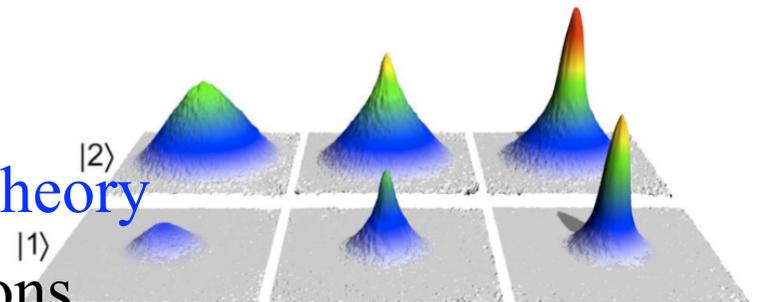
$\Lambda_{\text{pionless}}$

momenta $Q \ll m_\pi$

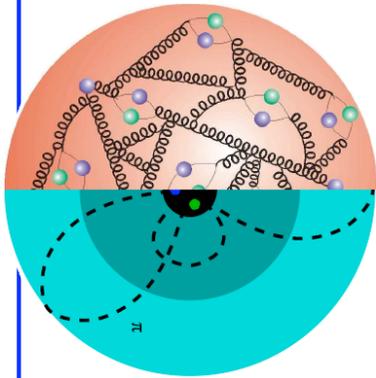
Λ / Resolution dependence of nuclear forces



$\Lambda_{\text{pionless}}$
momenta $Q \ll m_\pi$: pionless effective field theory
large scattering length physics and corrections



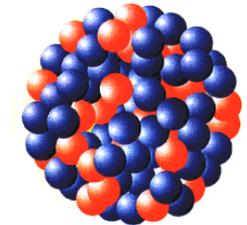
Λ / Resolution dependence of nuclear forces



Λ_{chiral}

momenta $Q \sim \lambda^{-1} \sim m_{\pi}$: chiral effective field theory

neutrons and protons interacting via pion exchanges
and shorter-range contact interactions

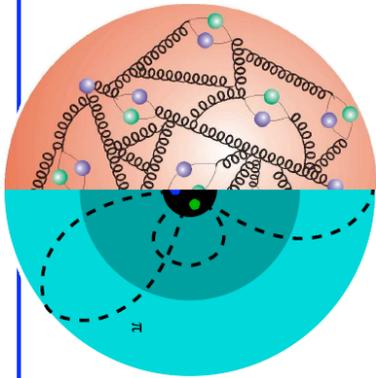


typical momenta in nuclei $\sim m_{\pi}$

$\Lambda_{\text{pionless}}$

momenta $Q \ll m_{\pi}$

Λ / Resolution dependence of nuclear forces



Effective theory for NN, 3N, many-N interactions and electroweak operators: resolution scale/ Λ -dependent

$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$

Λ_{chiral}

momenta $Q \sim \lambda^{-1} \sim m_{\pi}$: chiral effective field theory



$\Lambda_{\text{pionless}}$

momenta $Q \ll m_{\pi}$

Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV

	NN	3N	4N
LO $\mathcal{O}\left(\frac{Q^0}{\Lambda^0}\right)$			
NLO $\mathcal{O}\left(\frac{Q^2}{\Lambda^2}\right)$			
N ² LO $\mathcal{O}\left(\frac{Q^3}{\Lambda^3}\right)$			
N ³ LO $\mathcal{O}\left(\frac{Q^4}{\Lambda^4}\right)$			
	+ ...	+ ...	+ ...

limited resolution at low energies,
can expand in powers $(Q/\Lambda_b)^n$

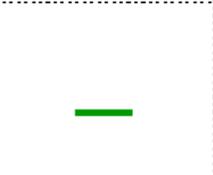
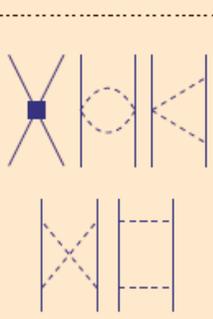
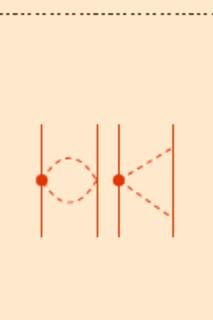
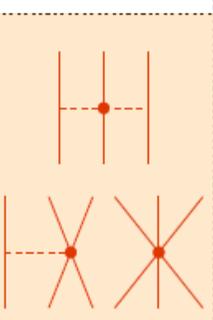
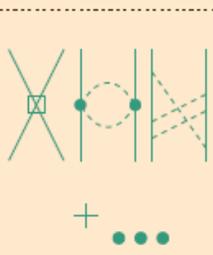
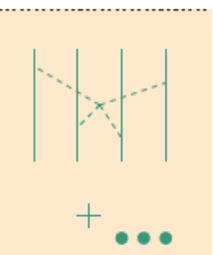
LO, $n=0$ - leading order,
NLO, $n=2$ - next-to-leading order,...

expansion parameter $\sim 1/3$

(compare to multipole expansion
for a charge distribution)

Chiral Effective Field Theory for nuclear forces

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N ² LO $\mathcal{O}\left(\frac{Q^3}{\Lambda^3}\right)$			
N ³ LO $\mathcal{O}\left(\frac{Q^4}{\Lambda^4}\right)$	 + ...	 + ...	 + ...

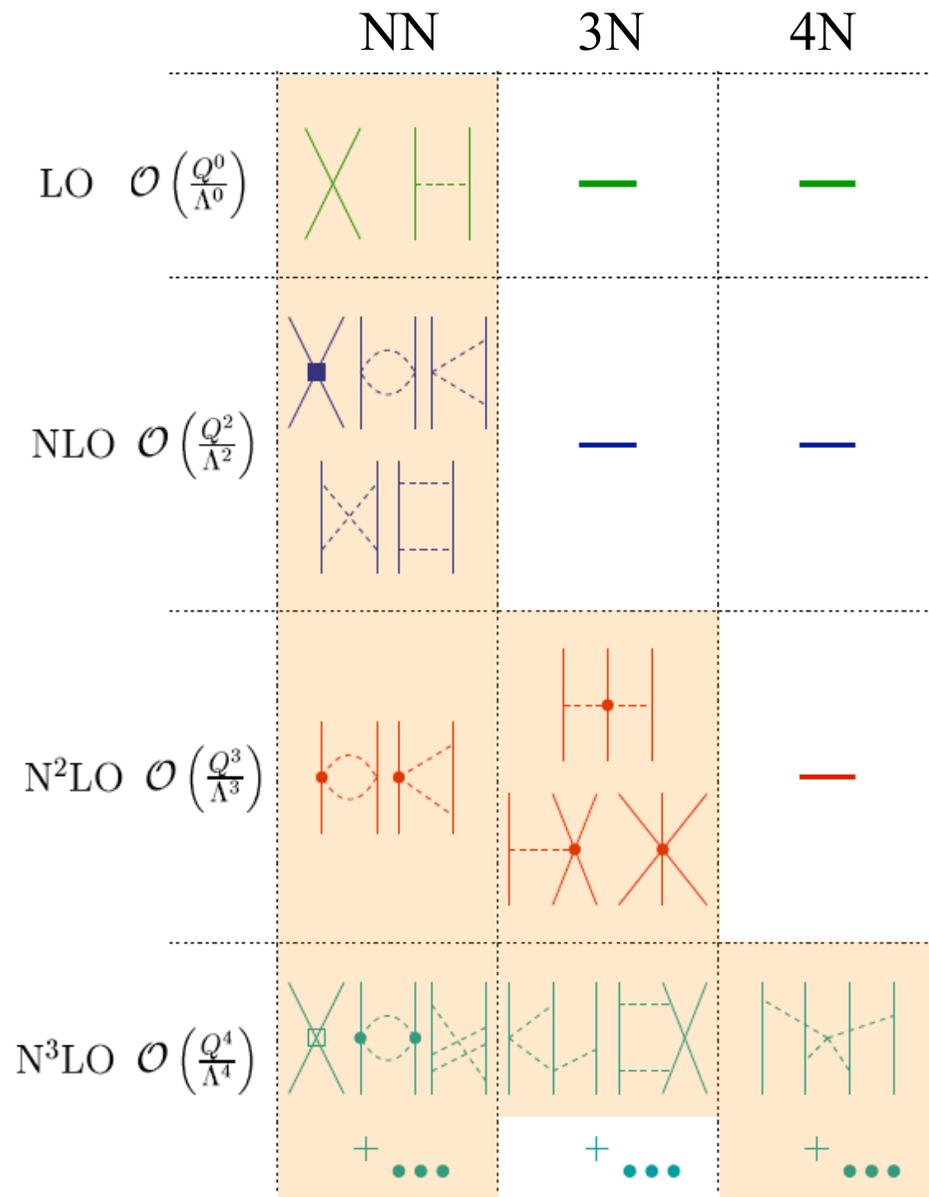
limited resolution at low energies,
can expand in powers $(Q/\Lambda_b)^n$

LO, $n=0$ - leading order,
NLO, $n=2$ - next-to-leading order,...

Question: Why is there no $n=1$ contribution?

Chiral Effective Field Theory for nuclear forces

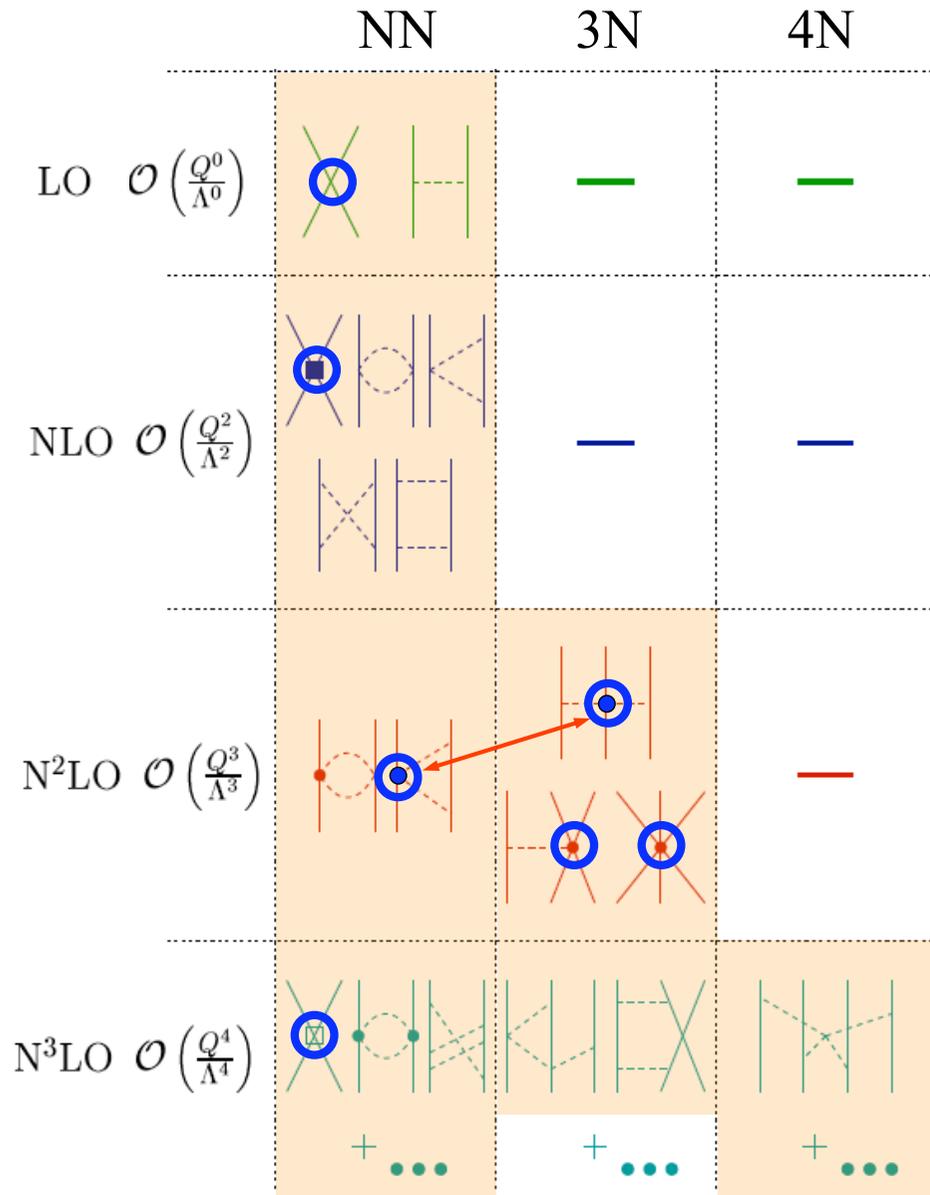
Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



include long-range pion physics

Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



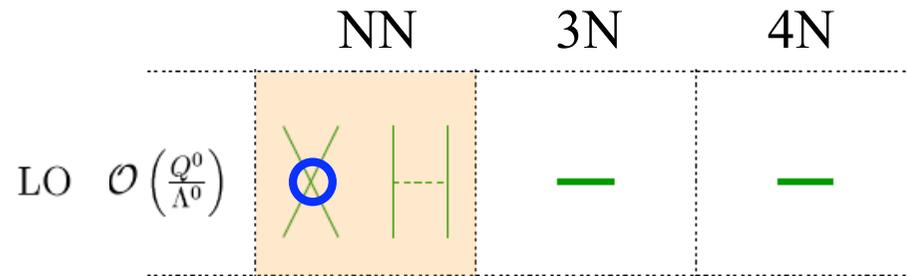
include long-range pion physics

details at short distance not resolved

capture in few **short-range couplings**,
fit to experiment once, Λ -dependent

Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



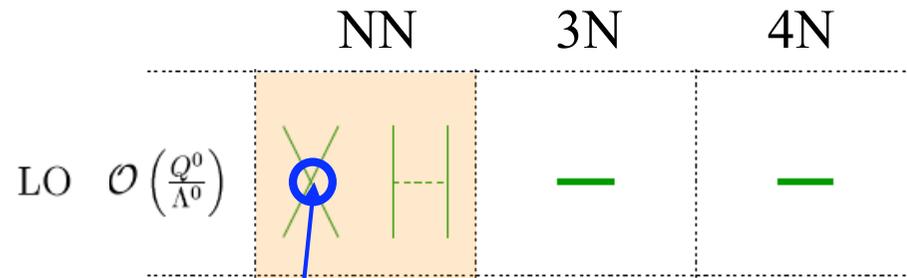
Question: What is $V(r)$ for NN interactions at LO?

Hint: One-pion exchange + contact interaction

but with resolution scale $\Lambda \sim 500$ MeV (cutoff on momenta)

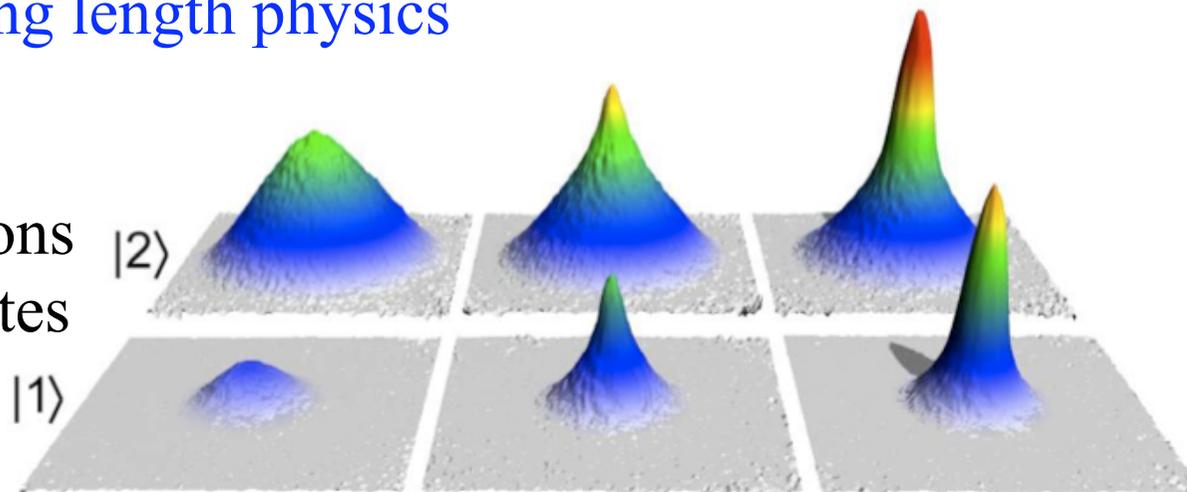
Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



large scattering length physics

${}^6\text{Li}$ fermions
2 spin states

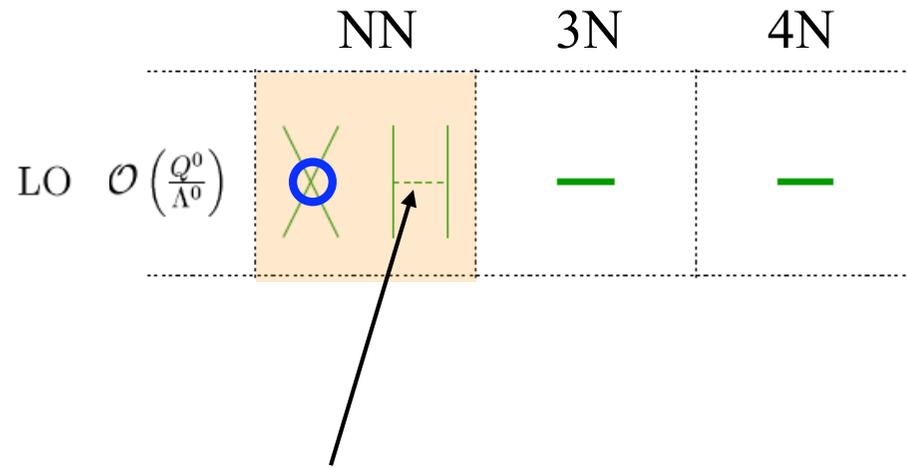


from M. Zwierlein

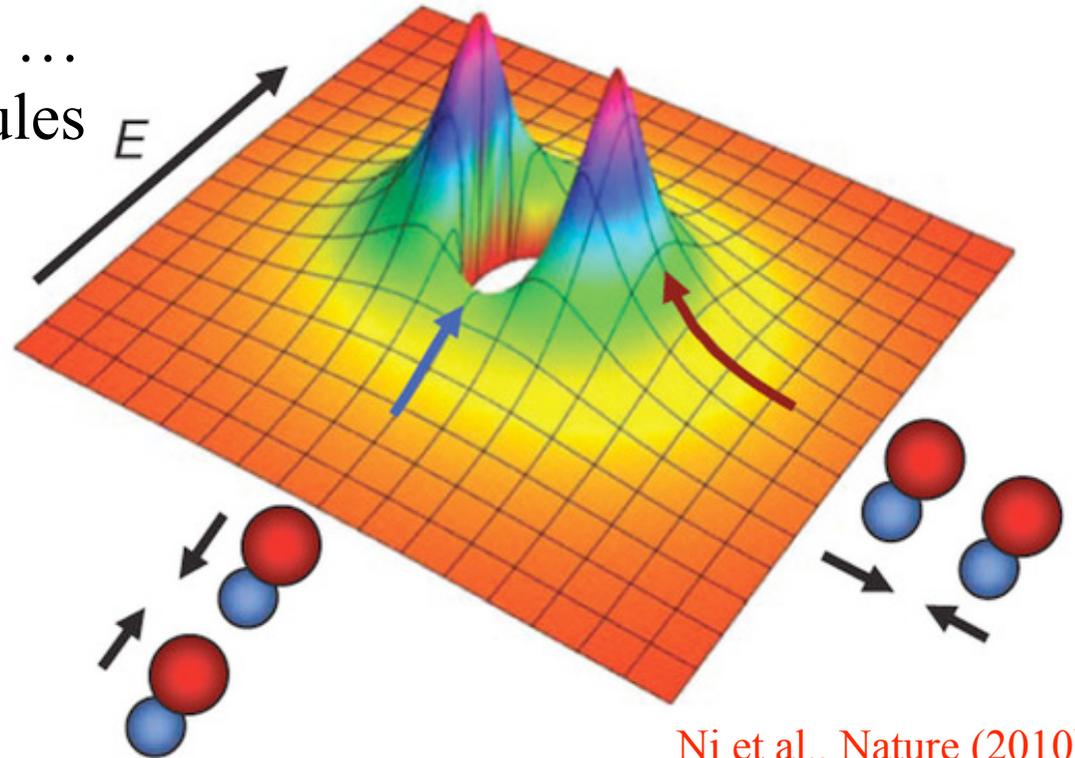
neutrons with same density and temperature
have the same properties!

Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV

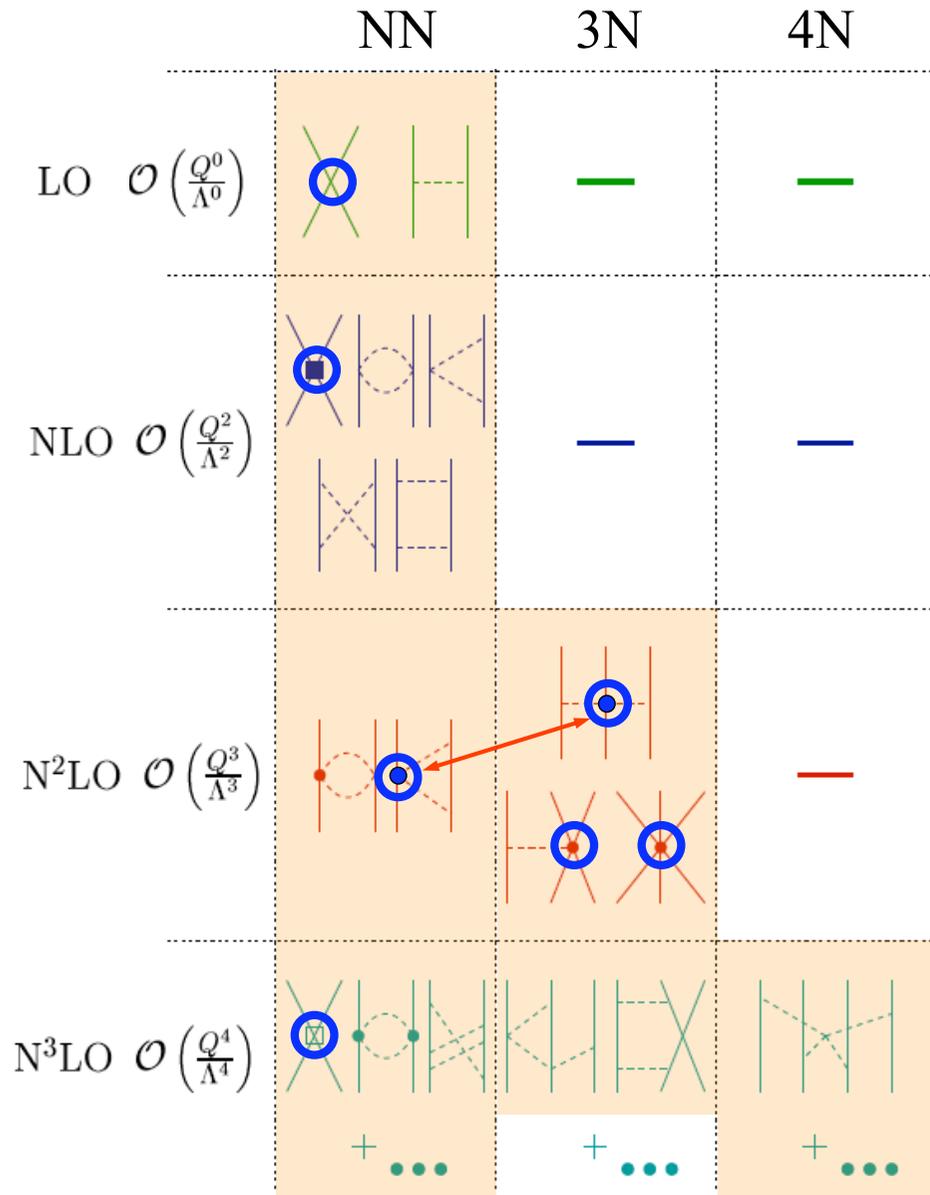


pion tensor/dipole interactions + ...
 → compare to cold polar molecules



Chiral Effective Field Theory for nuclear forces

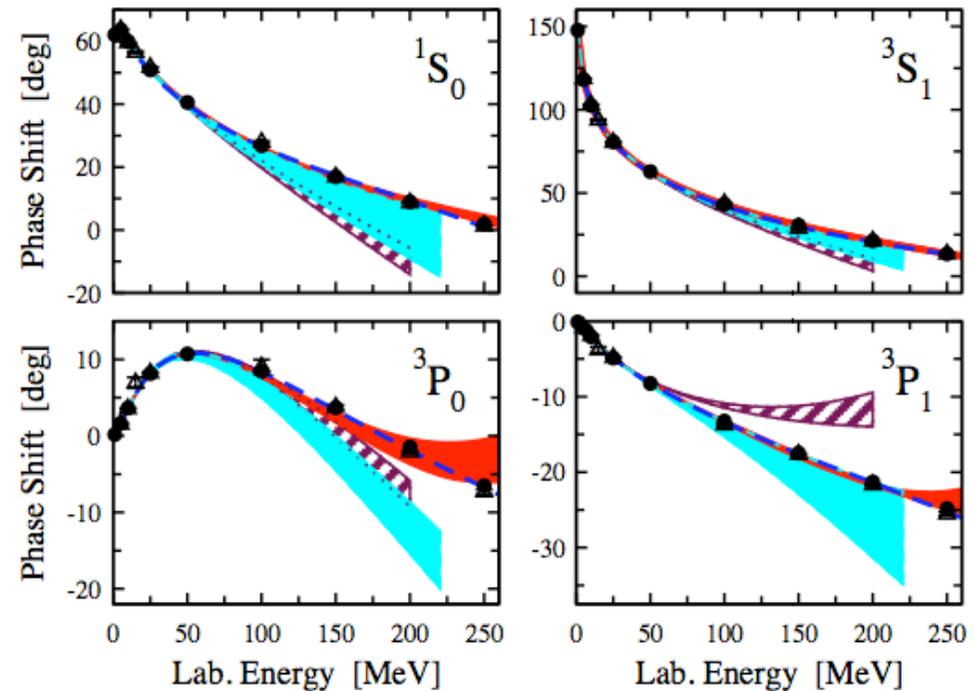
Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



systematic: can work to desired accuracy and obtain error estimates from truncation order and Λ variation

accurate reproduction of

low-energy NN scattering at N³LO

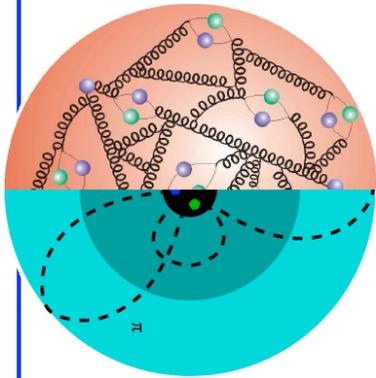


Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Kaiser, Machleidt, Meissner,...

Nuclear forces and the Renormalization Group (RG)

RG evolution to lower resolution/cutoffs Bogner, Kuo, AS, Furnstahl,...

$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$



Λ_{chiral}

Λ

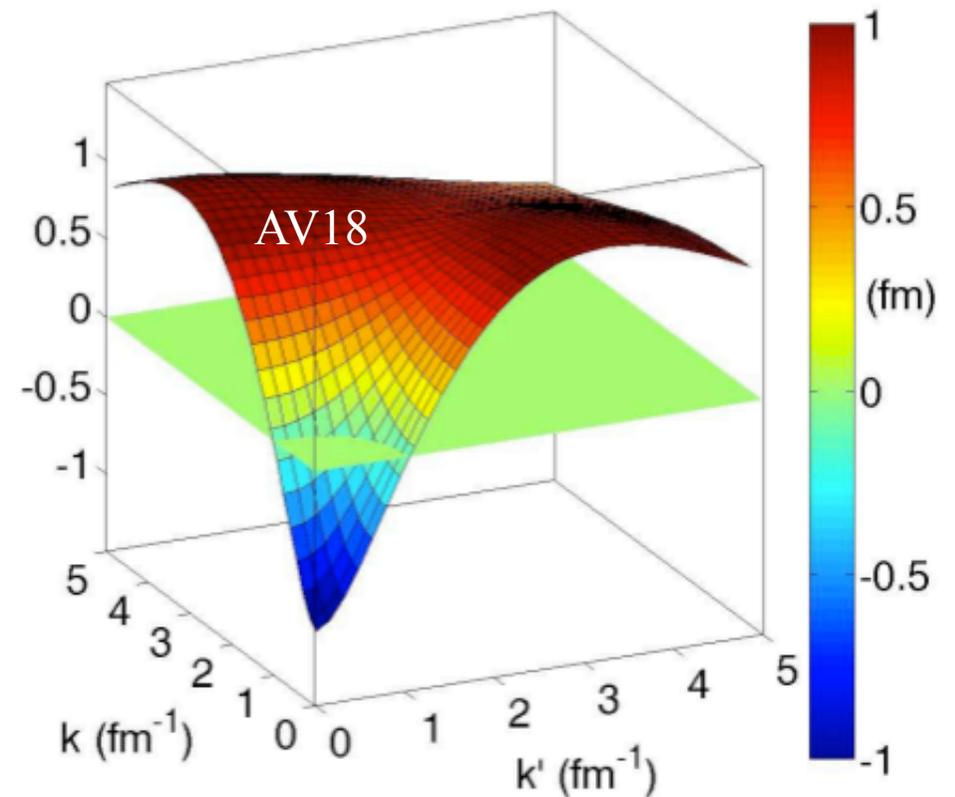
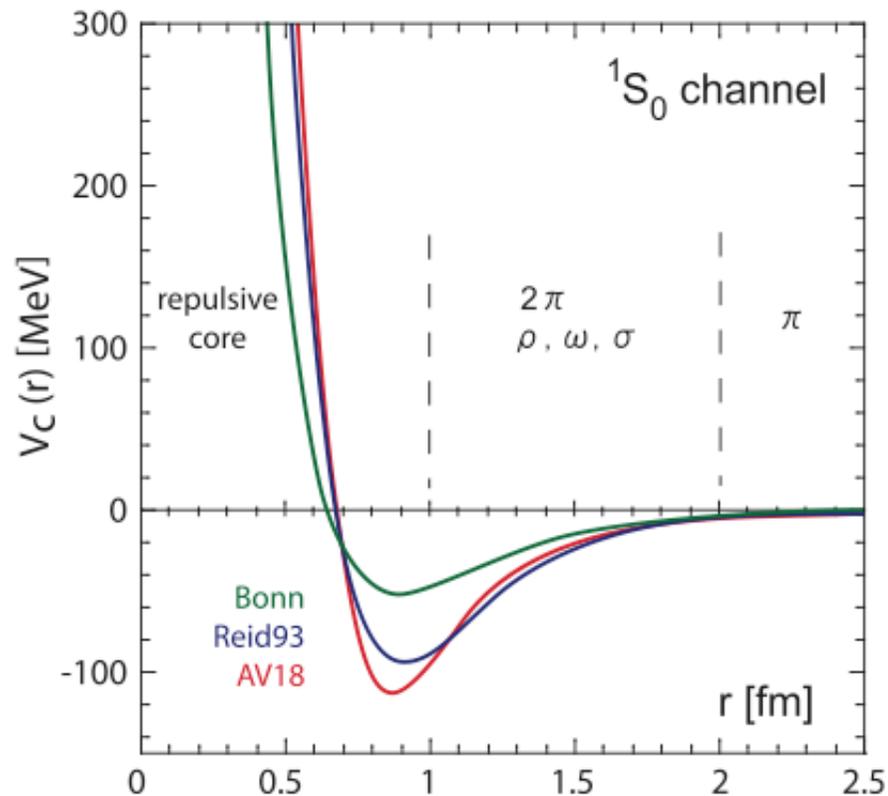


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for NN interactions (preserves NN observables)



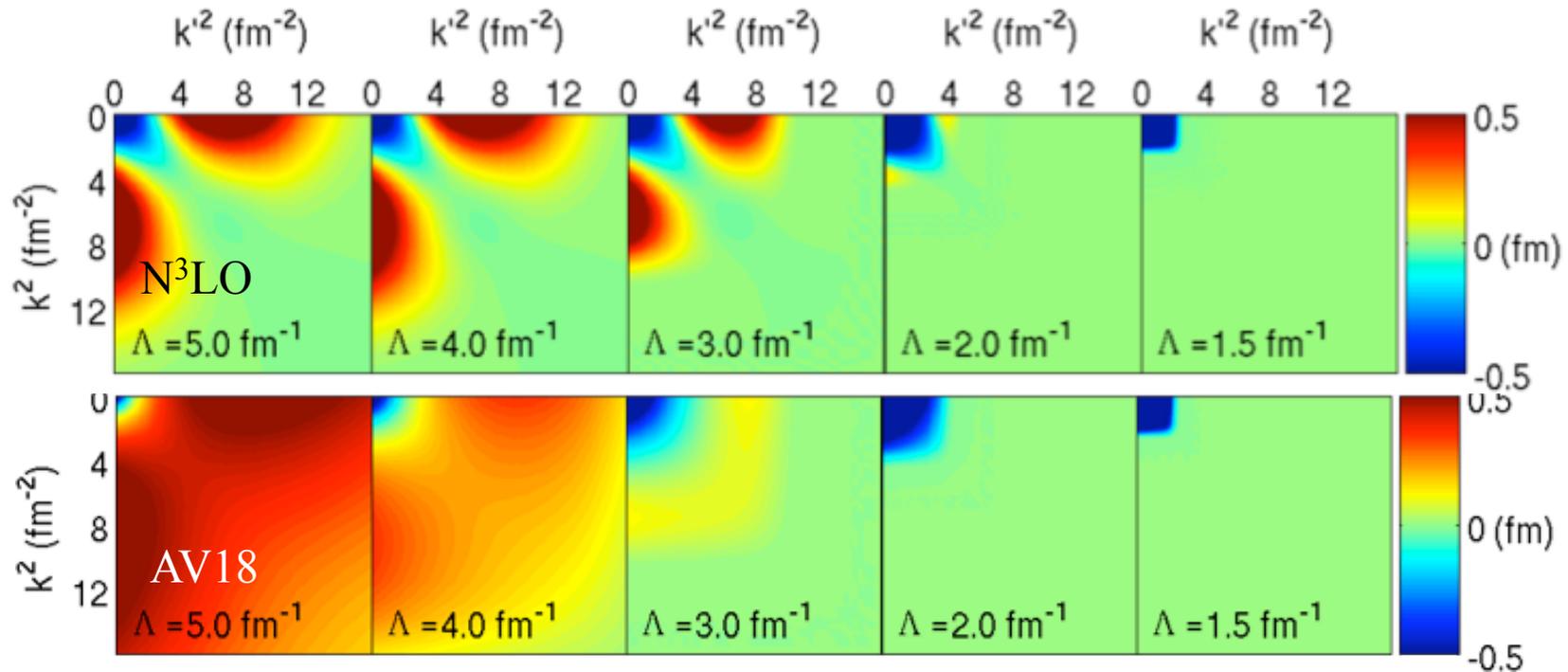
red = short-range repulsion

Nuclear forces and the Renormalization Group (RG)

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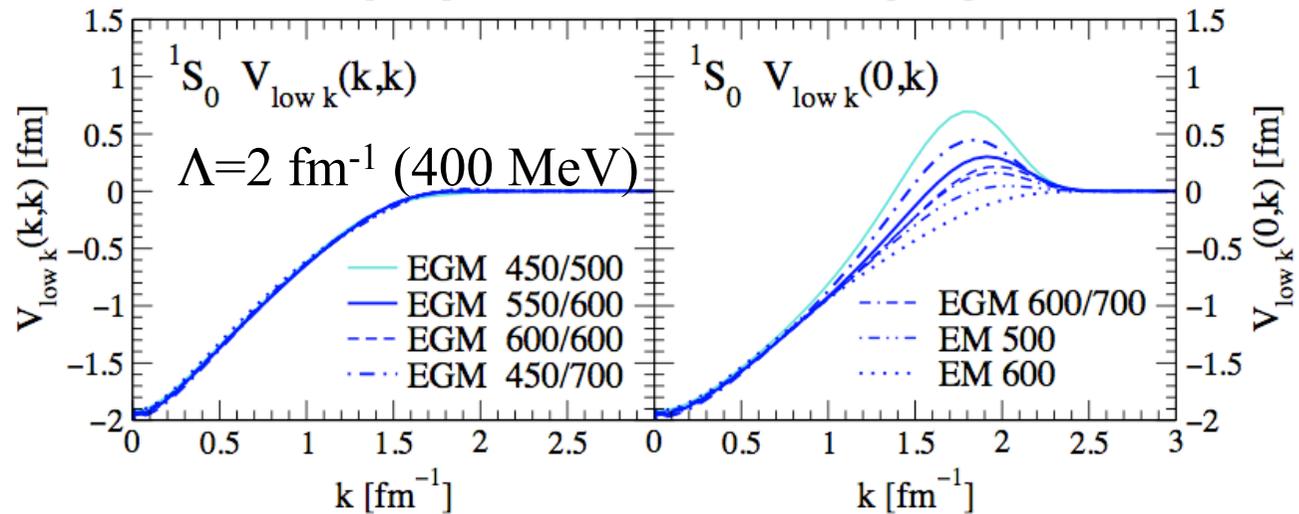
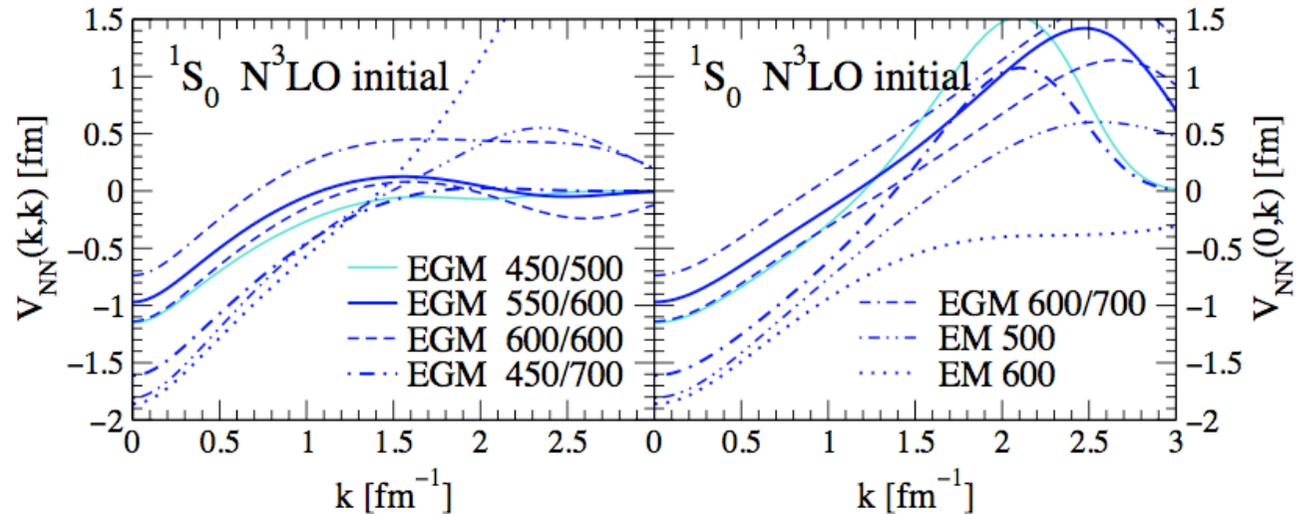
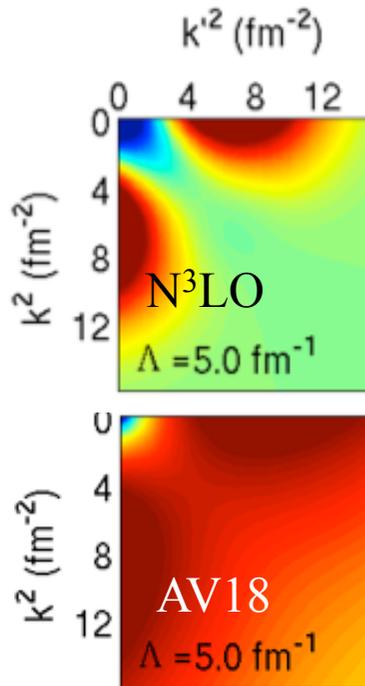
low-momentum interactions $V_{\text{low } k}(\Lambda)$

RG decouples low-momentum physics from high momenta

Nuclear forces and the Renormalization Group (RG)

RG evolution to lower resolution/cutoffs Bogner, Kuo, AS, Furnstahl,...

$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$



low-momentum i
RG decouples low

low-momentum universality from different chiral N³LO potentials

Chiral Effective Field Theory for nuclear forces

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	NN	3N	4N
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N ³ LO $\mathcal{O}\left(\frac{Q^4}{\Lambda^4}\right)$			
	+ ...	+ ...	+ ...

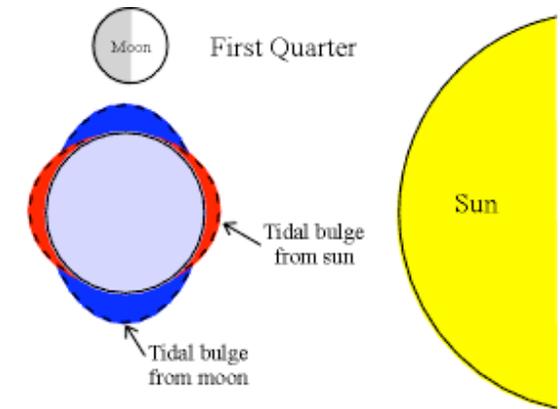
limited resolution at low energies,
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LO, $n=0$ - leading order,
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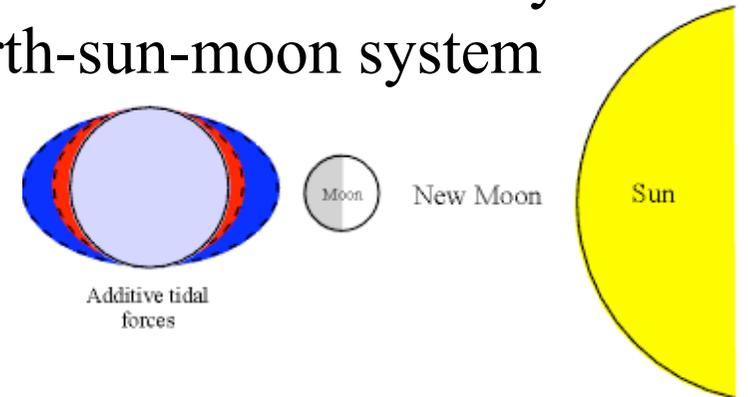
expansion parameter $\sim 1/3$

(compare to multipole expansion
for a charge distribution)

Why are there three-body forces?



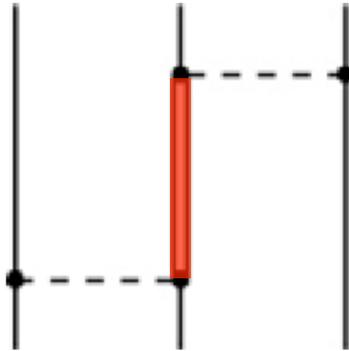
tidal effects lead to 3-body forces
in earth-sun-moon system



Why are there three-nucleon (3N) forces?

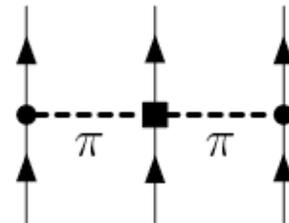
Nucleons are finite-mass composite particles,
can be excited to resonances

dominant contribution from $\Delta(1232 \text{ MeV})$



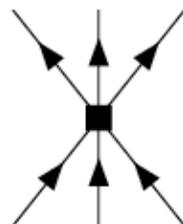
+ many shorter-range parts

in chiral EFT (Delta-less):

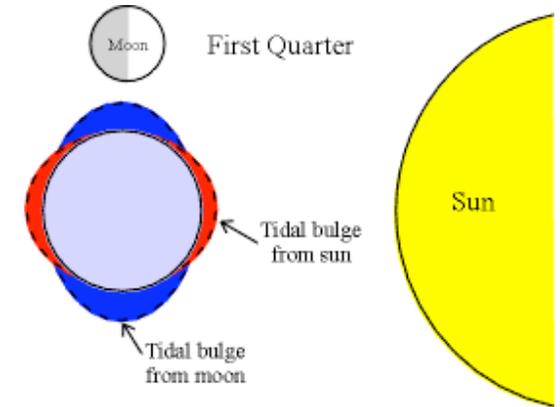


+ shorter-range parts

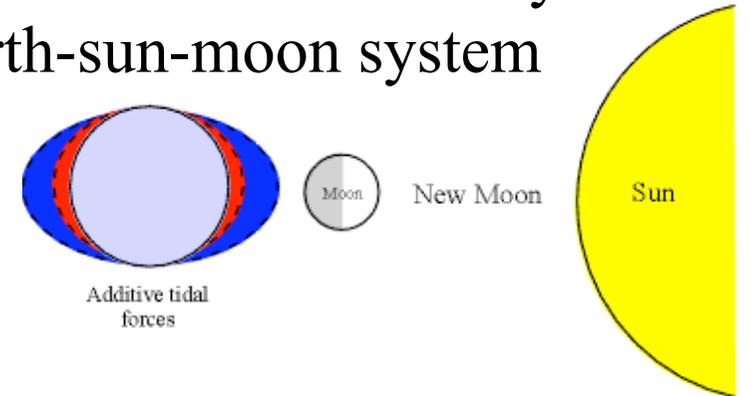
in pionless EFT:



+ higher-order parts

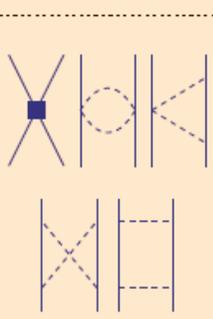
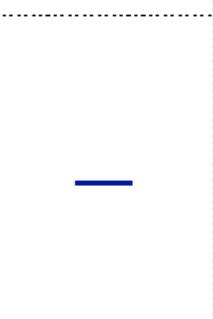
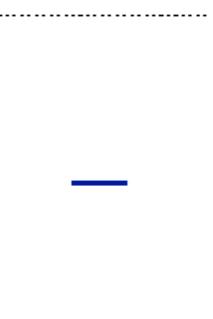
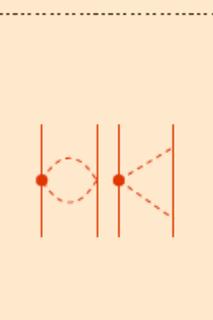
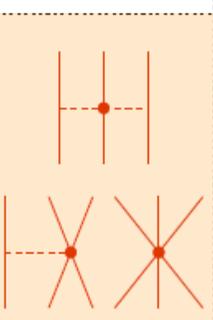
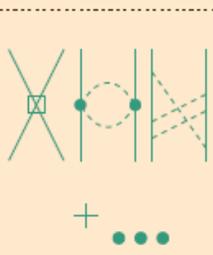
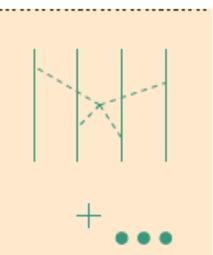


tidal effects lead to 3-body forces
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Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV

	NN	3N	4N
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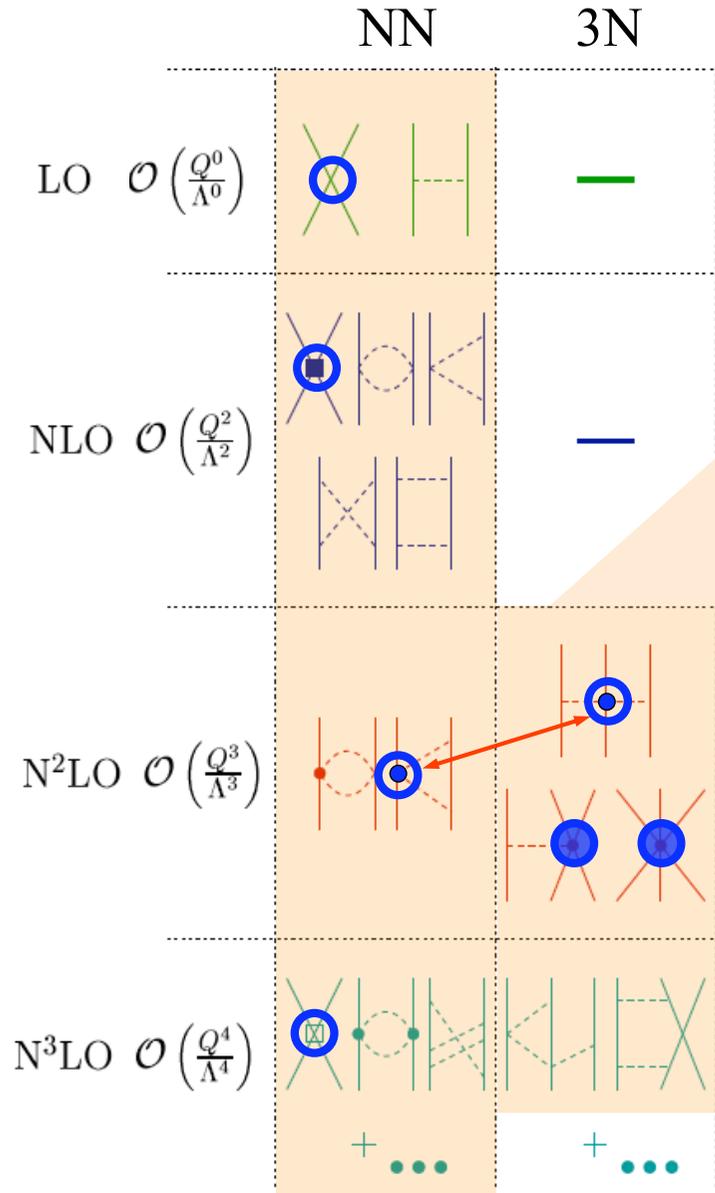
limited resolution at low energies,
can expand in powers $(Q/\Lambda_b)^n$

LO, $n=0$ - leading order,
NLO, $n=2$ - next-to-leading order,...

Question: Why do 3N forces start at $n=3$
(without explicit Deltas)?

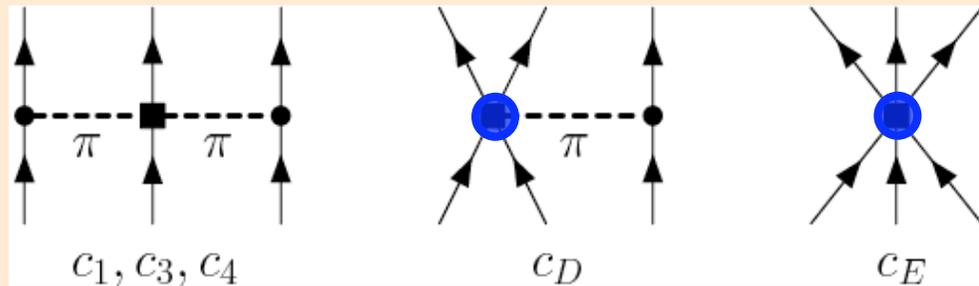
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consistent NN-3N interactions

3N,4N: only 2 new couplings to N³LO



c_i from π N and NN Meissner et al. (2007)

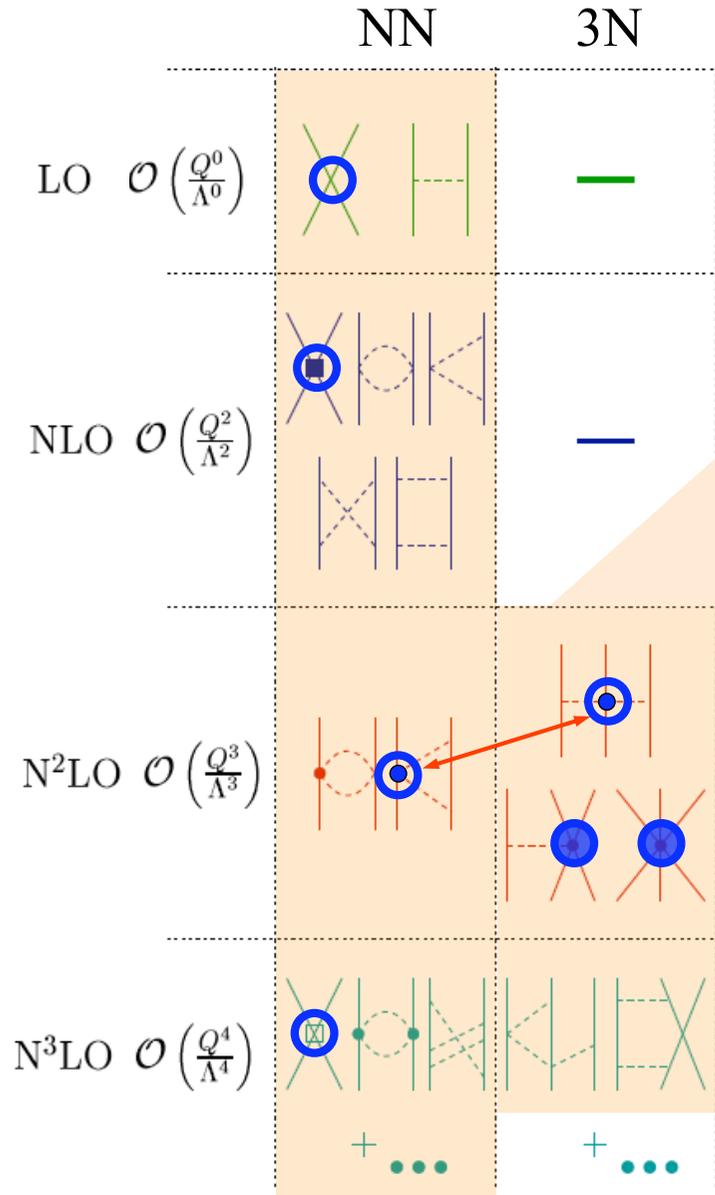
$$c_1 = -0.9_{-0.5}^{+0.2}, \quad c_3 = -4.7_{-1.0}^{+1.2}, \quad c_4 = 3.5_{-0.2}^{+0.5}$$

single- Δ : $c_1=0$, $c_3=-c_4/2=-3 \text{ GeV}^{-1}$

c_D, c_E fit to ${}^3\text{H}$ binding energy and ${}^4\text{He}$ radius (or ${}^3\text{H}$ beta decay half-life)

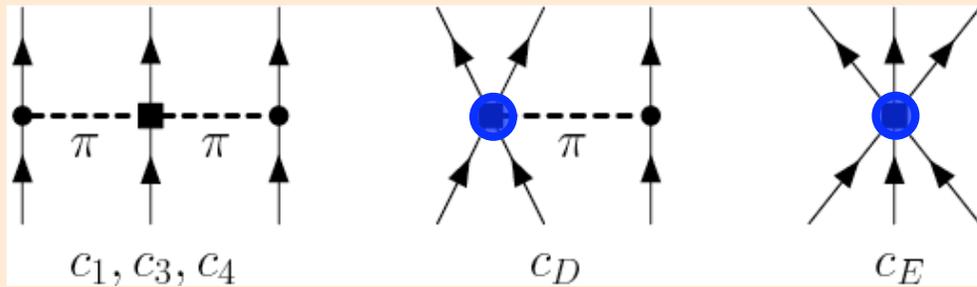
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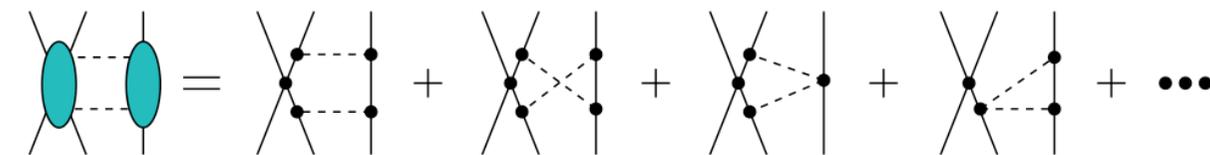
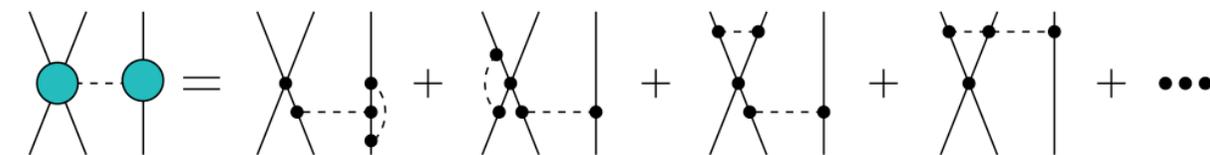
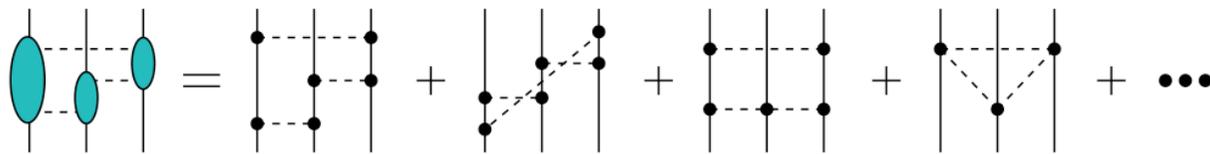
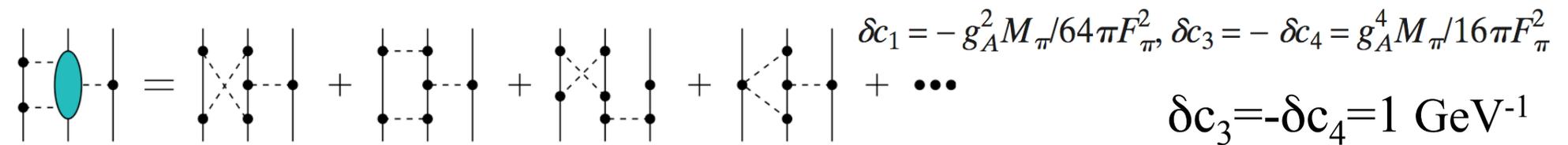
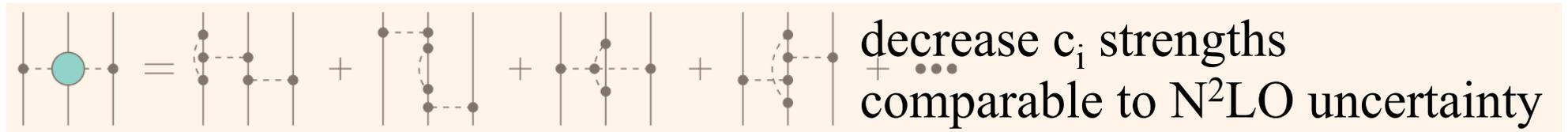
Question: Why are the next order 3N forces (n=4) parameter-free (without new 3N contact interactions)?

Subleading chiral 3N forces

parameter-free N^3LO from Epelbaum et al.; Bernard et al. (2007), Ishikawa, Robilotta (2007)

one-loop contributions:

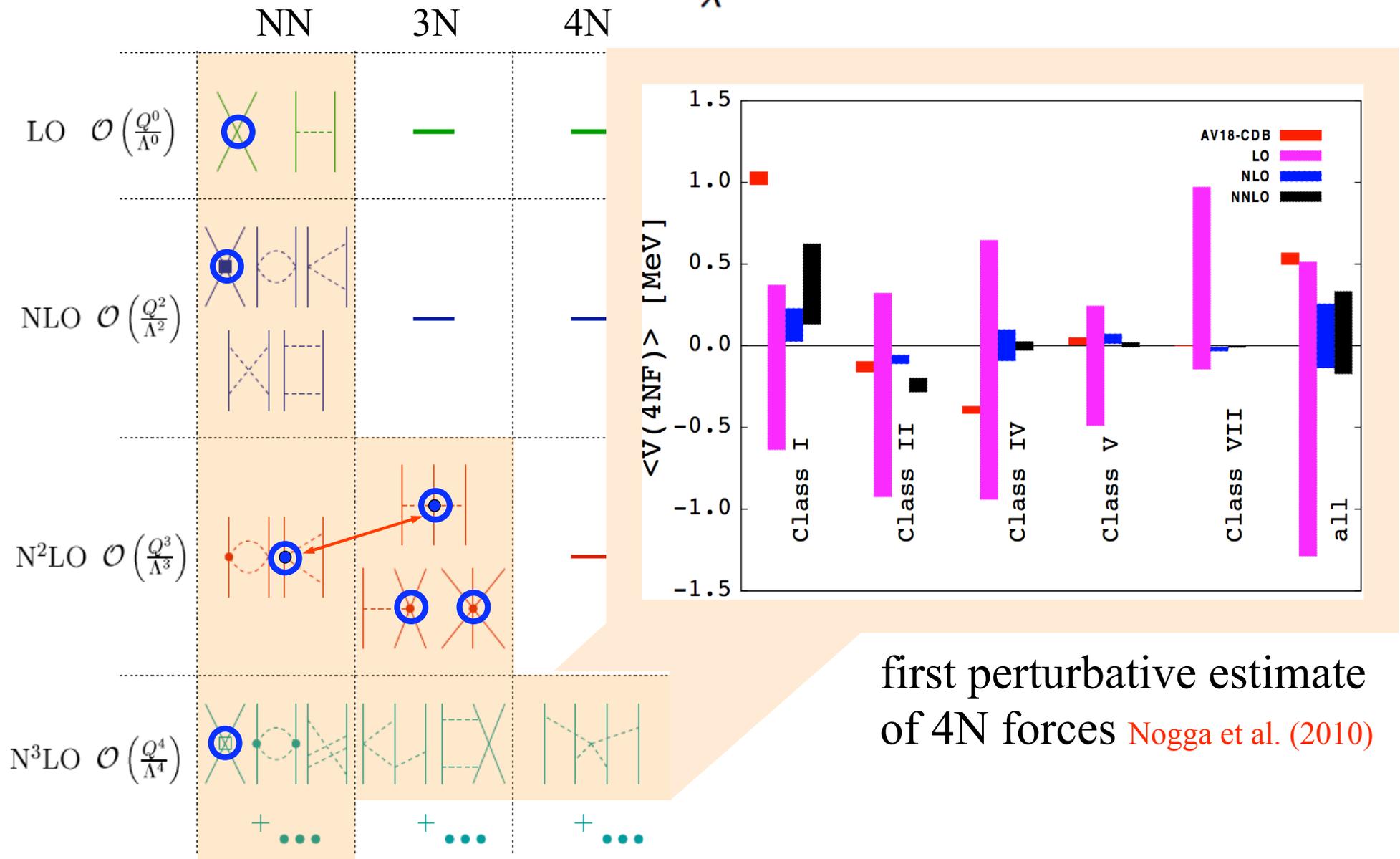
2π -exchange, 2π - 1π -exchange, rings, contact- 1π -, contact- 2π -exchange



$1/m$ corrections: spin-orbit parts, interesting for A_y puzzle

Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



first perturbative estimate
of 4N forces [Nogga et al. \(2010\)](#)

Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Kaiser, Machleidt, Meissner,...

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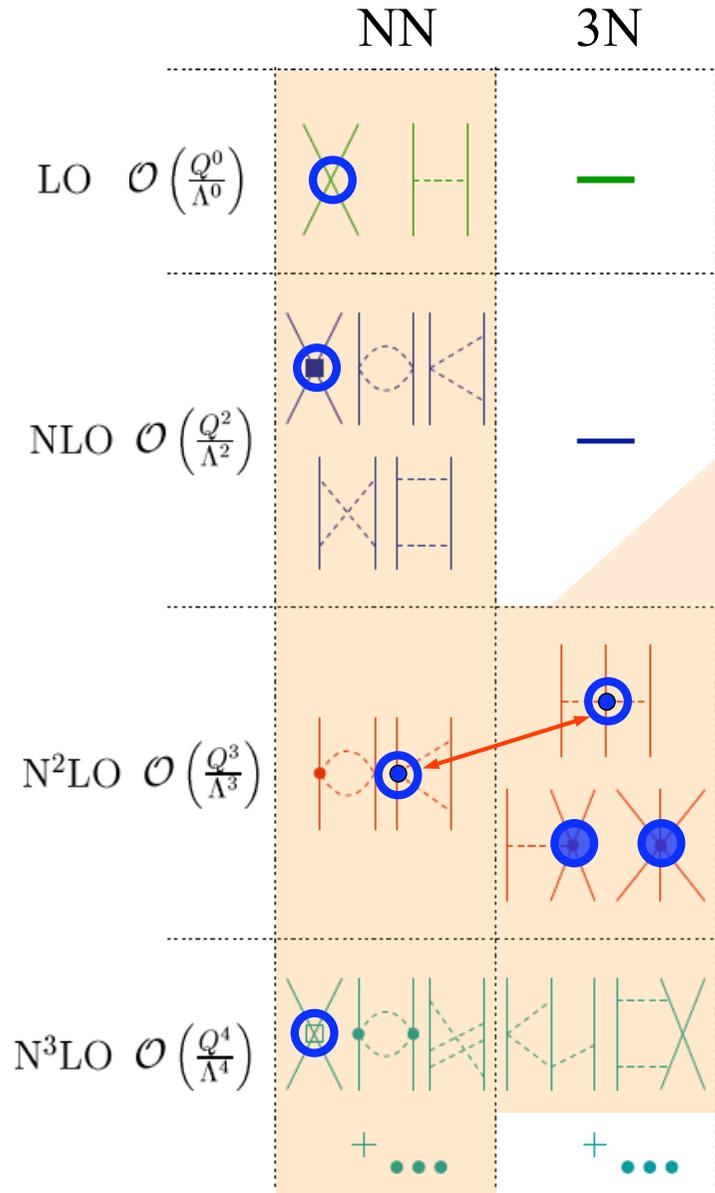
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Contact: schwenk@physik.tu-darmstadt.de

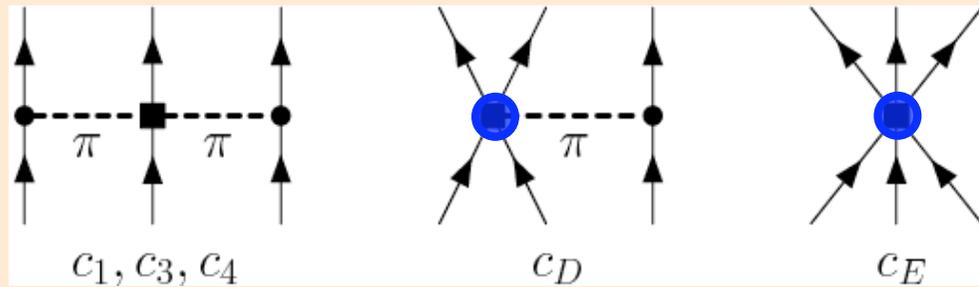
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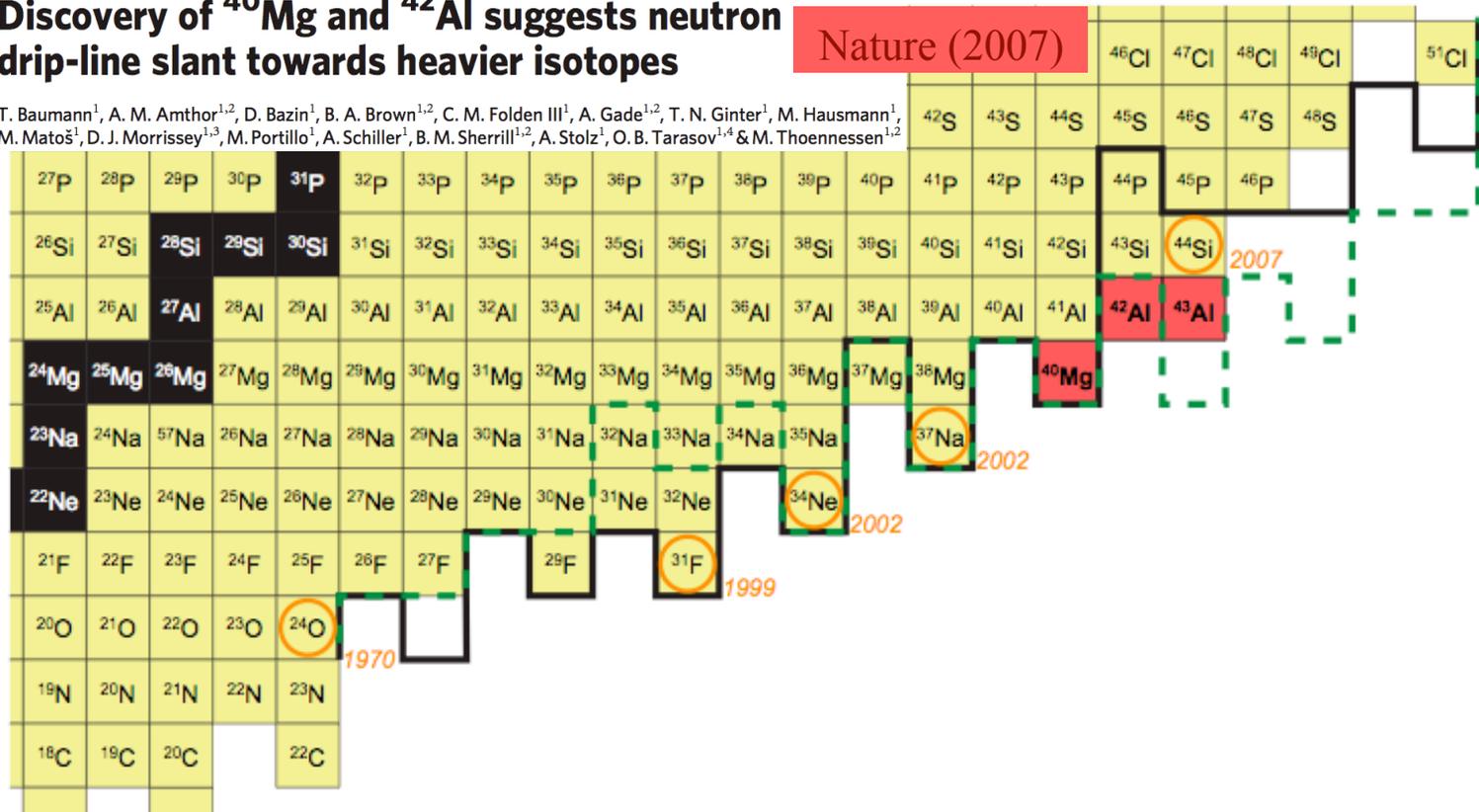
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Towards the limits of existence - the neutron drip-line

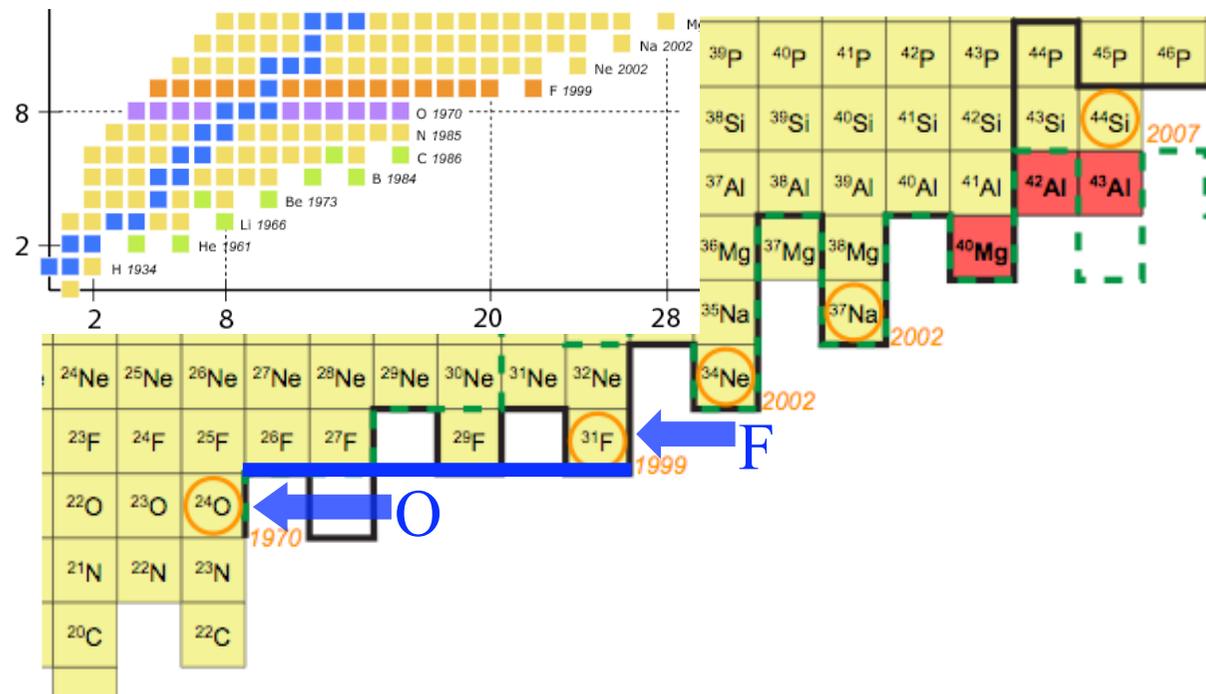
Discovery of ^{40}Mg and ^{42}Al suggests neutron drip-line slant towards heavier isotopes

Nature (2007)

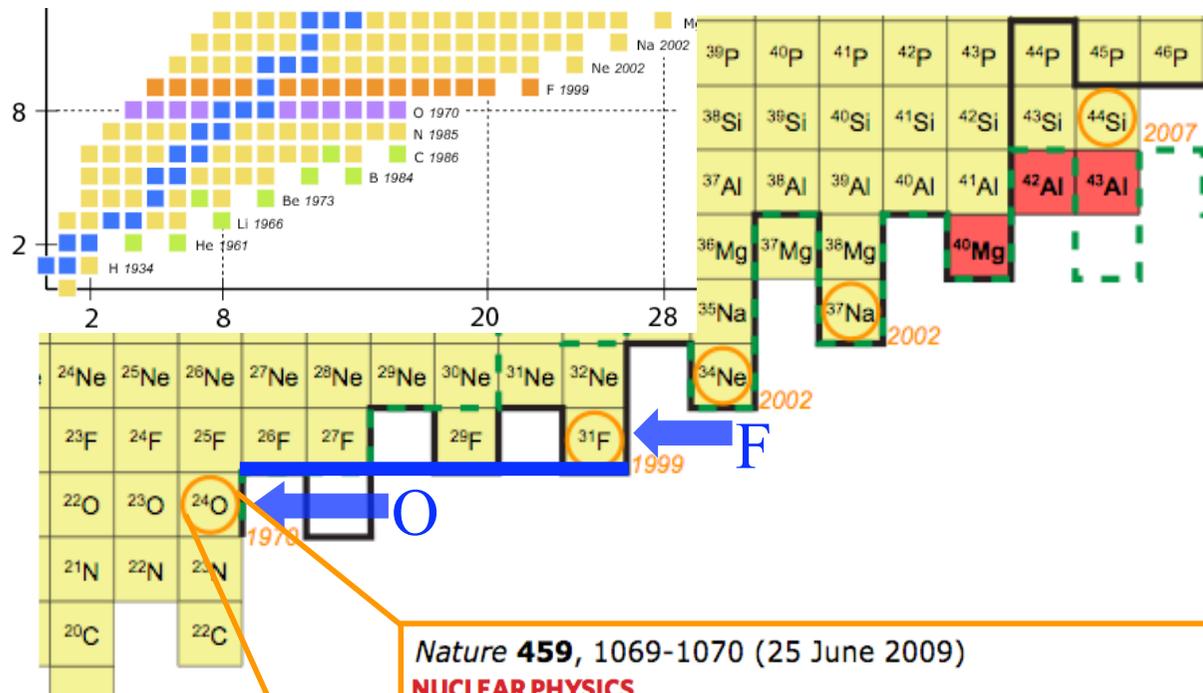
T. Baumann¹, A. M. Amthor^{1,2}, D. Bazin¹, B. A. Brown^{1,2}, C. M. Folden III¹, A. Gade^{1,2}, T. N. Ginter¹, M. Hausmann¹, M. Mátos¹, D. J. Morrissey^{1,3}, M. Portillo¹, A. Schiller¹, B. M. Sherrill^{1,2}, A. Stolz¹, O. B. Tarasov^{1,4} & M. Thoennessen^{1,2}



The oxygen anomaly



The oxygen anomaly



Nature **459**, 1069-1070 (25 June 2009)

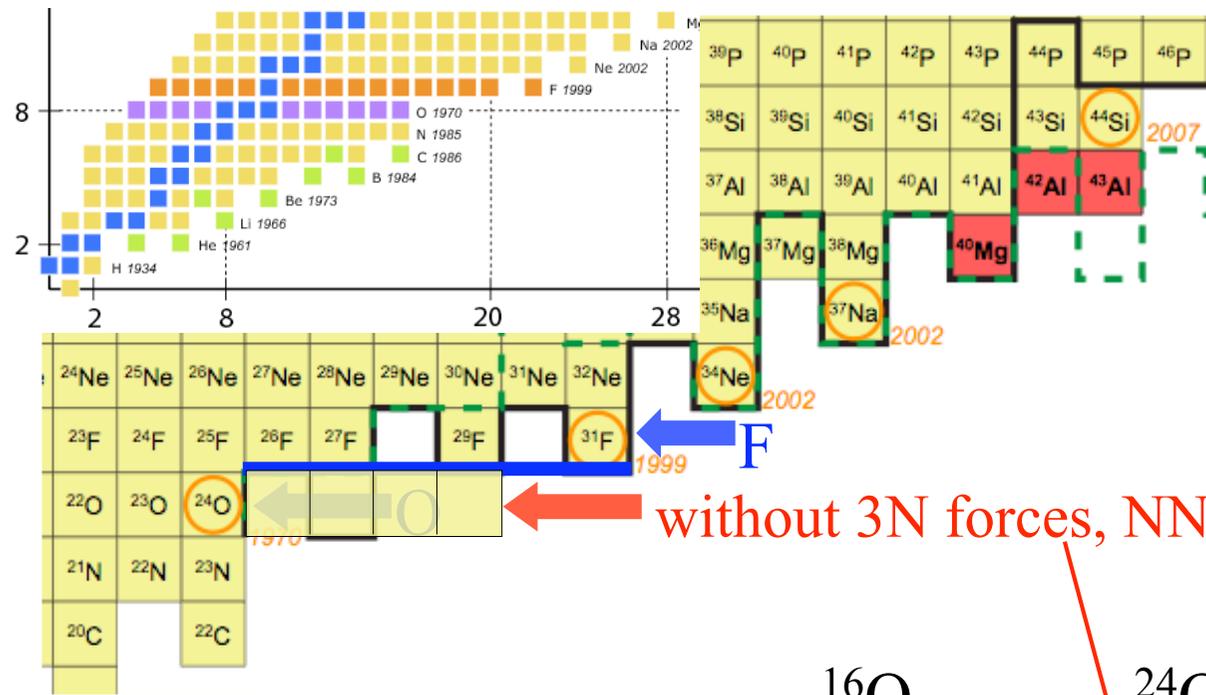
NUCLEAR PHYSICS

Unexpected doubly magic nucleus

Robert V. F. Janssens

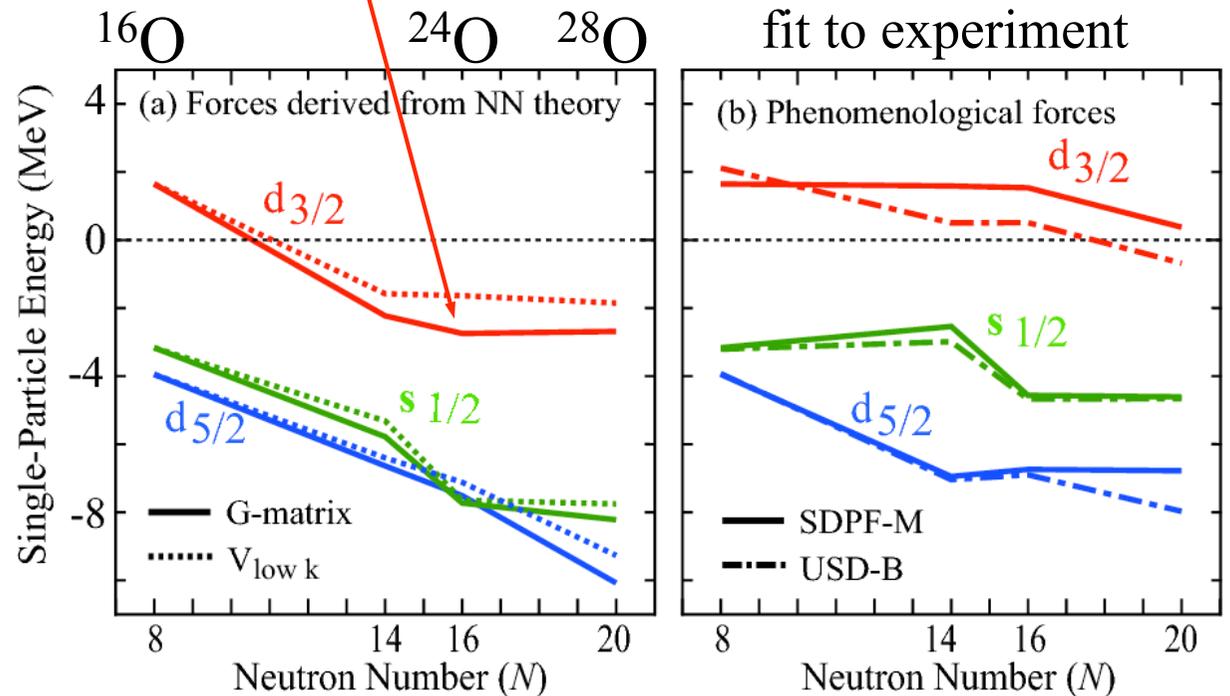
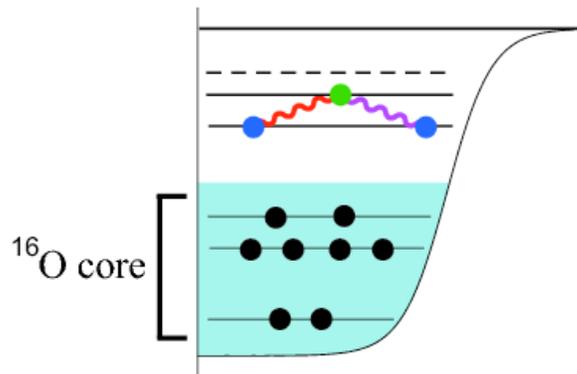
Nuclei with a 'magic' number of both protons and neutrons, dubbed doubly magic, are particularly stable. The oxygen isotope ^{24}O has been found to be one such nucleus — yet it lies just at the limit of stability.

The oxygen anomaly - not reproduced without 3N forces



without 3N forces, NN interactions too attractive

many-body theory based on two-nucleon forces:
drip-line incorrect at ^{28}O



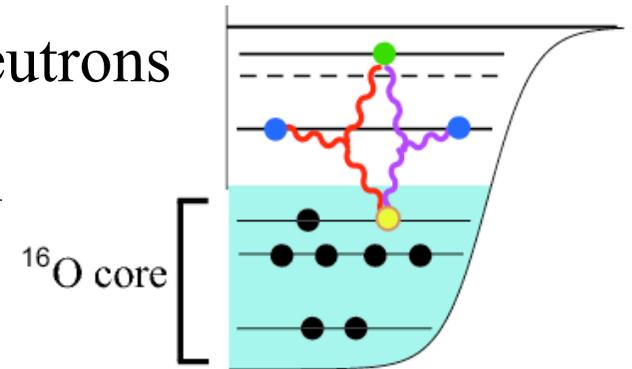
The oxygen anomaly - impact of 3N forces

include 'normal-ordered' 2-body part of 3N forces (enhanced by core A)

leads to repulsive interactions between valence neutrons

contributions from residual three valence-nucleon interactions suppressed by $E_{\text{ex}}/E_{\text{F}} \sim N_{\text{valence}}/N_{\text{core}}$

Friman, AS, arXiv:1101.4858.



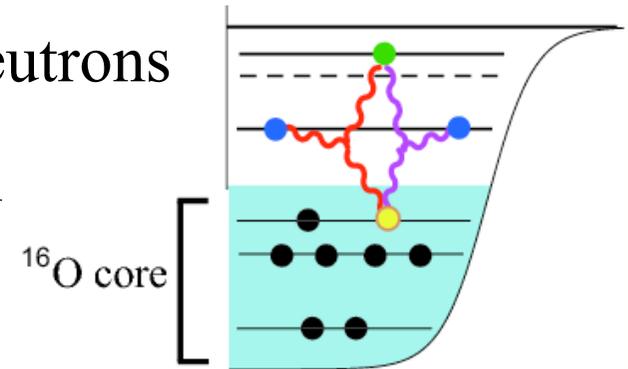
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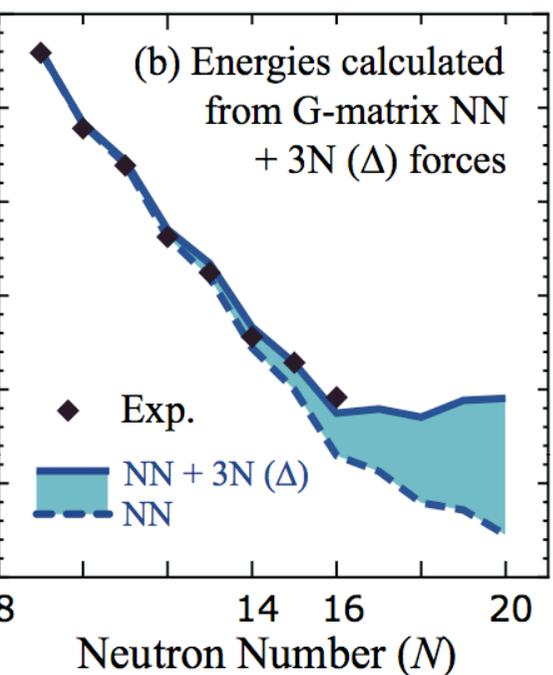
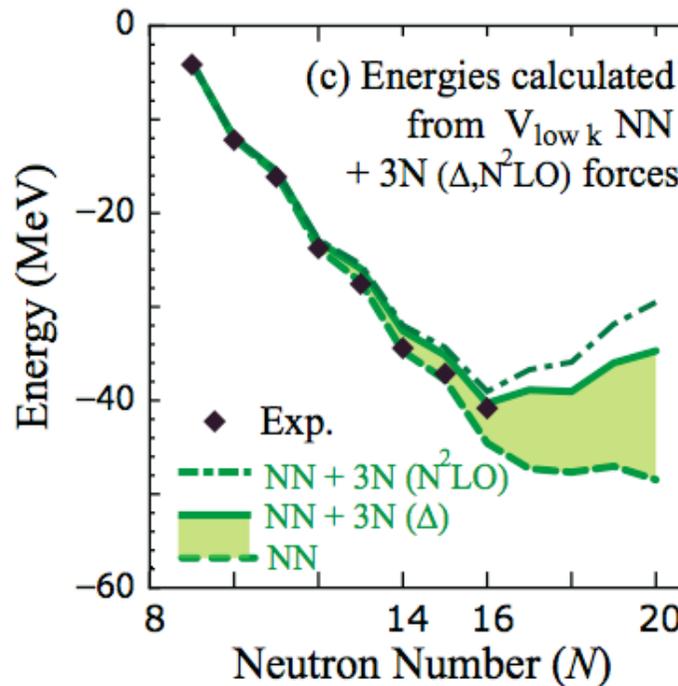
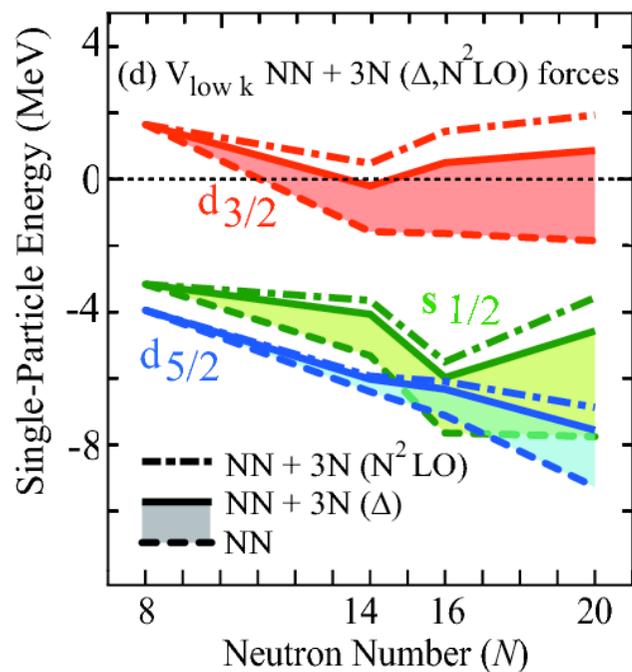
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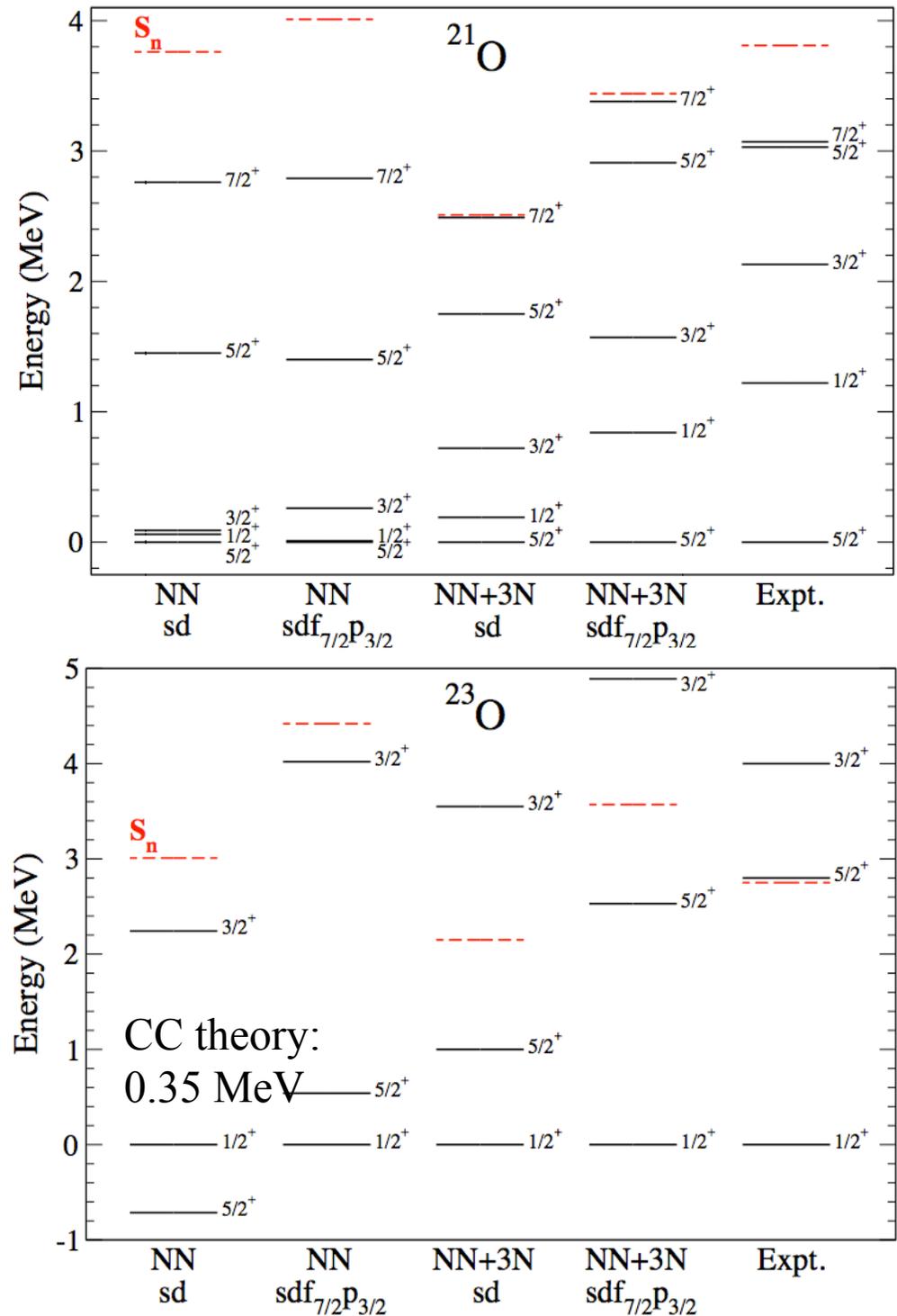
Oxygen spectra

focused on bound excited states

Holt, AS, arXiv:1108.2680.

NN only too compressed

3N contributions and extended valence space are key to reproduce excited states

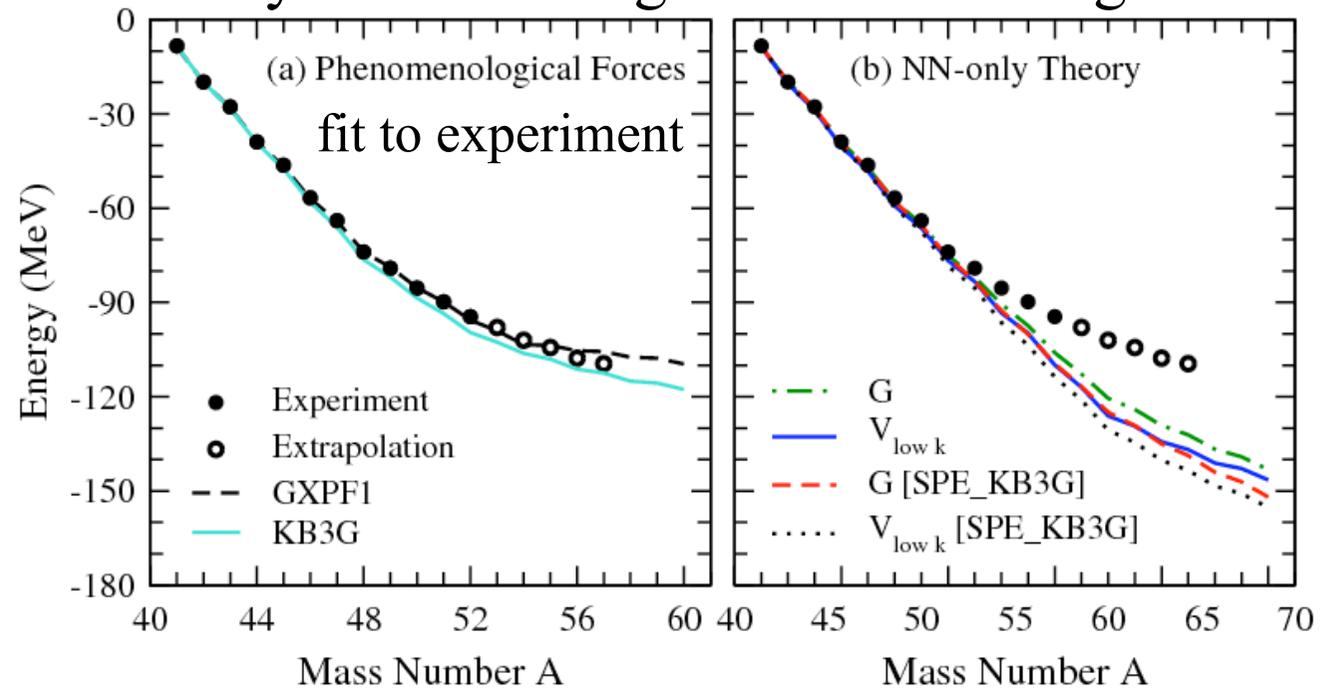


Evolution to neutron-rich calcium isotopes

repulsive 3N contributions also key for calcium ground-state energies

Holt et al., arXiv:1009:5984

mass measured to ^{52}Ca
shown to exist to ^{58}Ca



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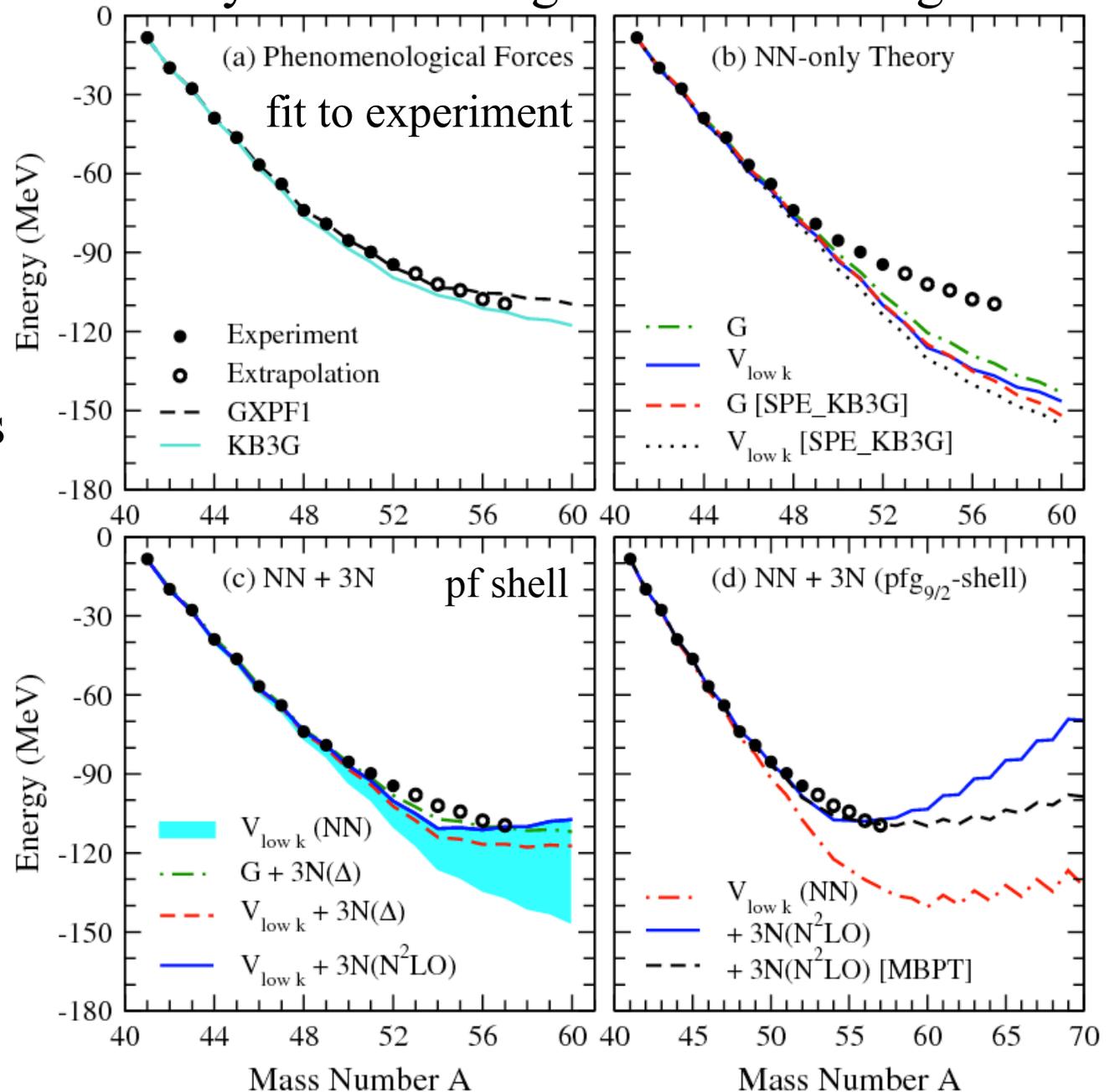
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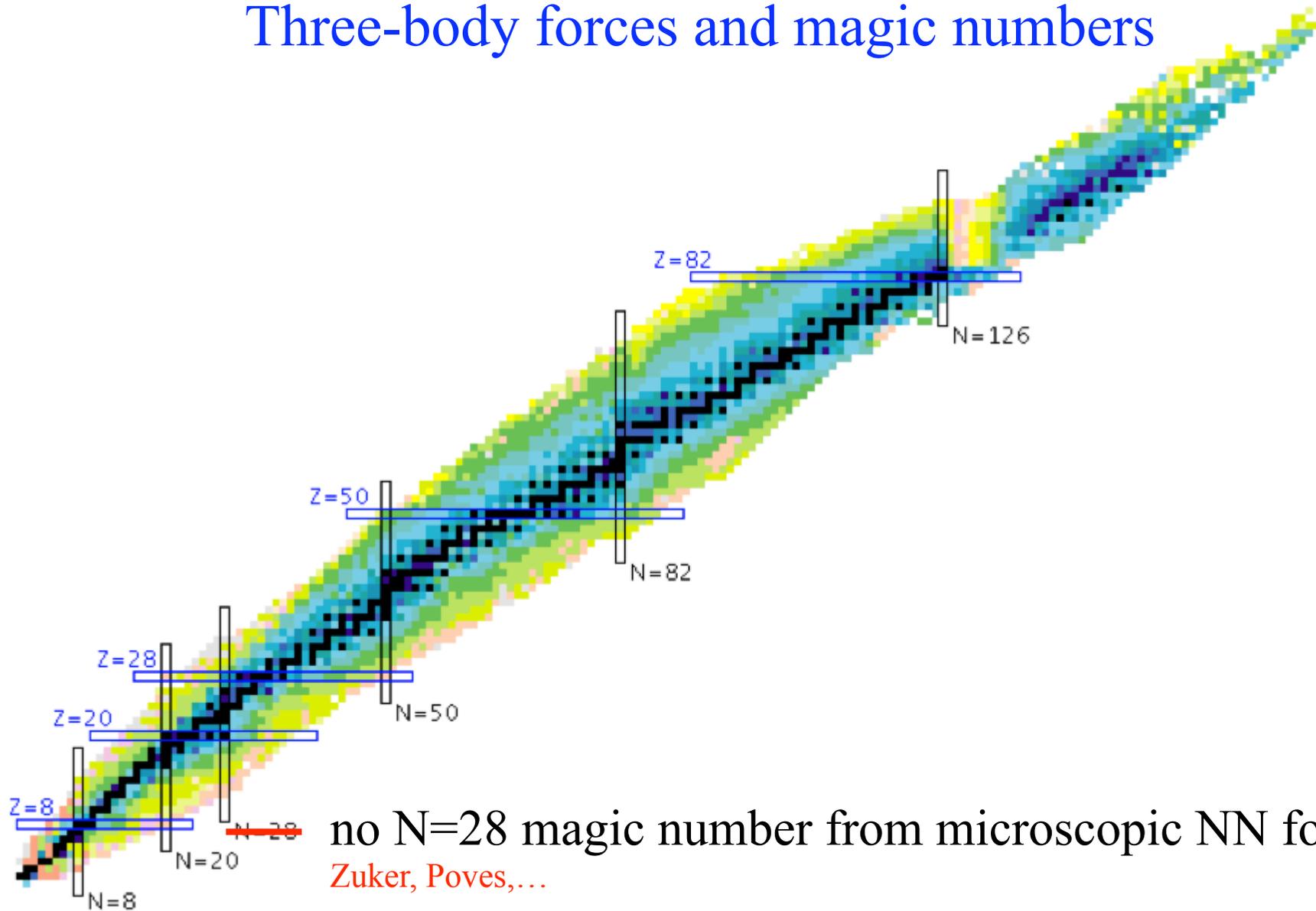
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predict drip-line
around ^{60}Ca ,

continuum contributions
will be key



Three-body forces and magic numbers



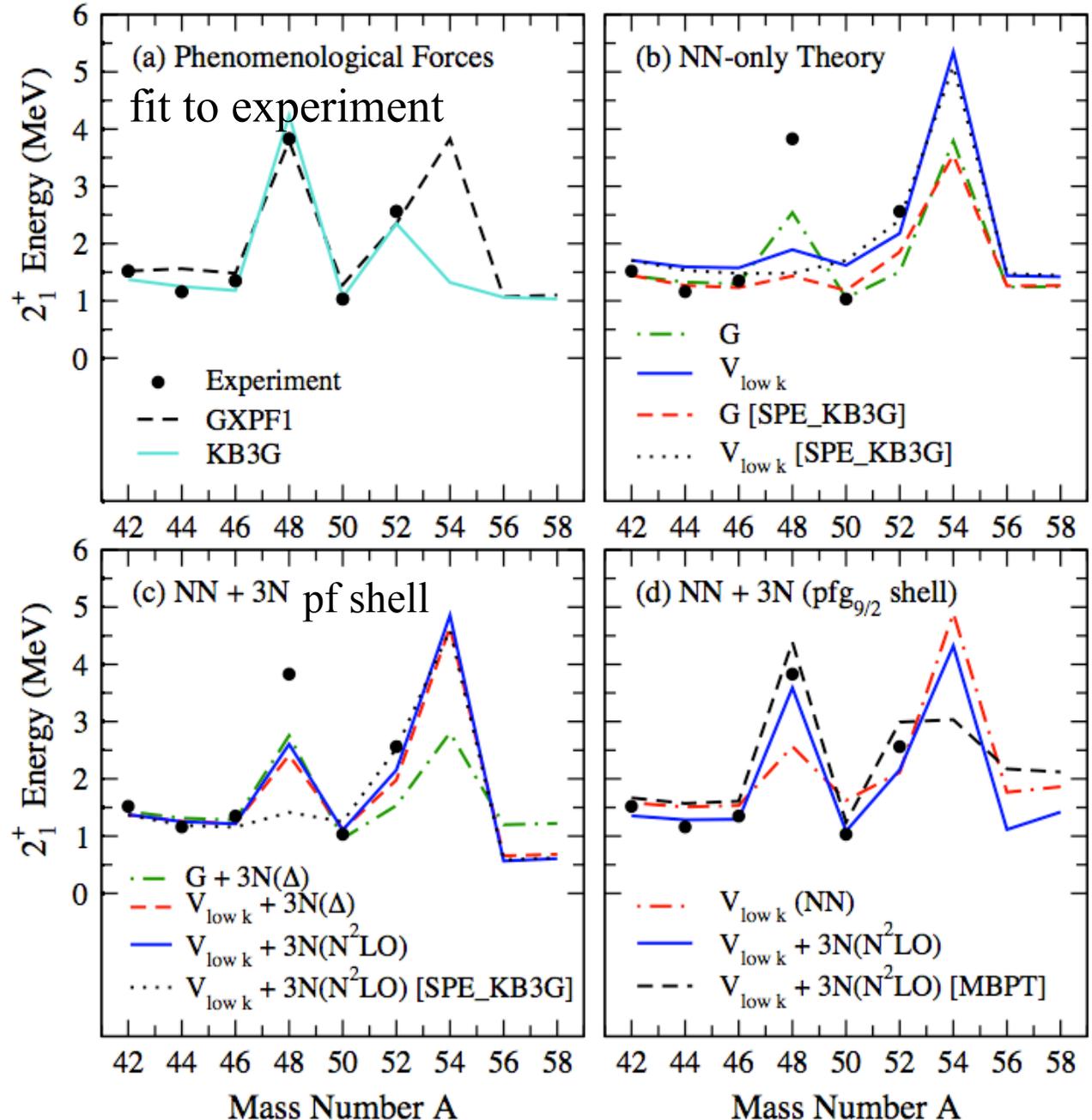
Three-body forces and magic numbers

3N mechanism important for shell structure

Holt et al., arXiv:1009:5984

N=28 shell closure
due to 3N forces
and single-particle
effects (^{41}Ca)

N=34: predict high
 2^+ excitation energy
in ^{54}Ca at 3-5 MeV





Outline

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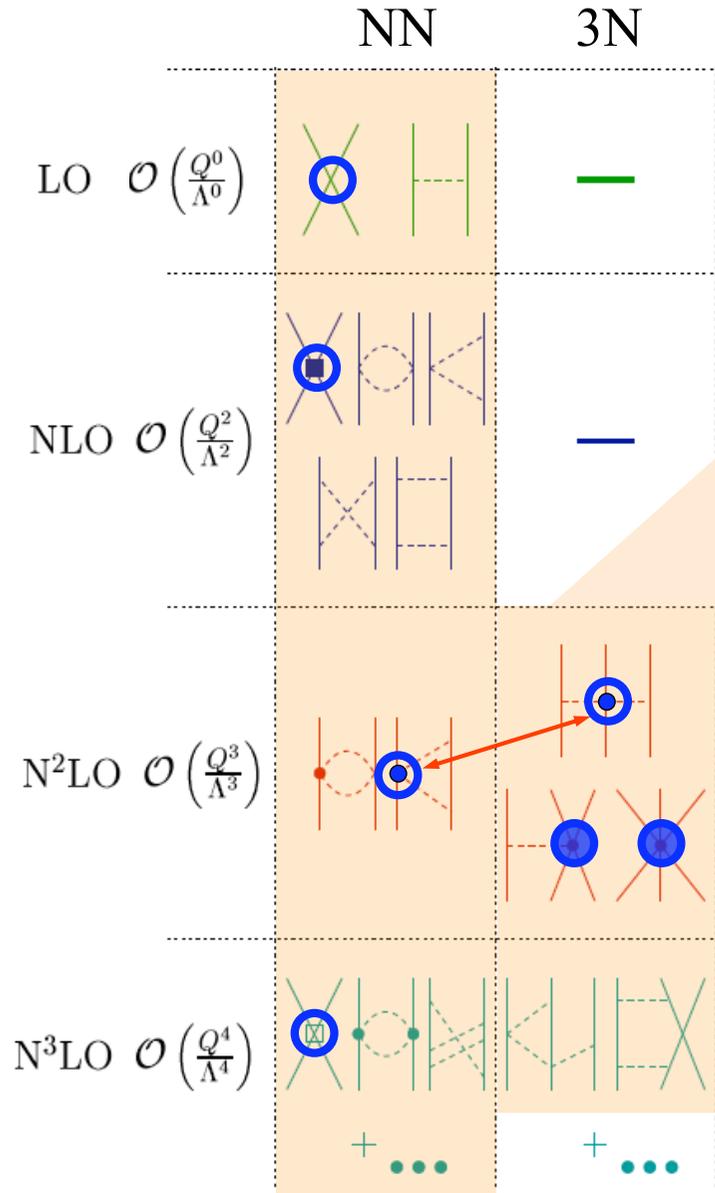
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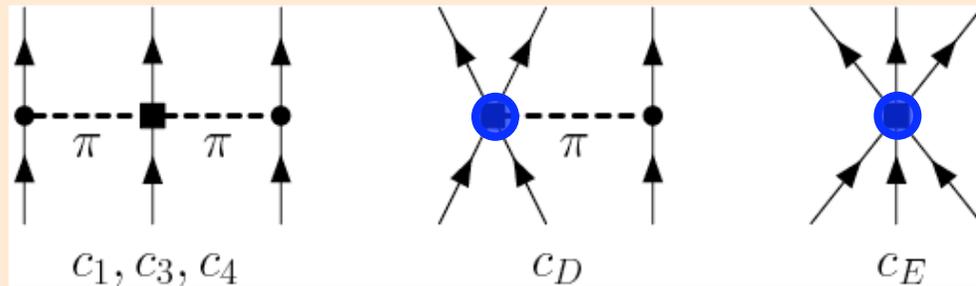
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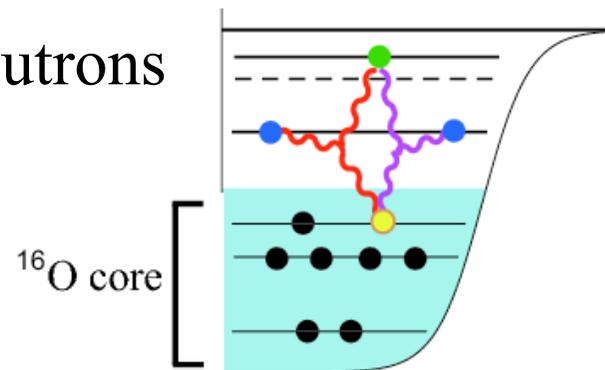
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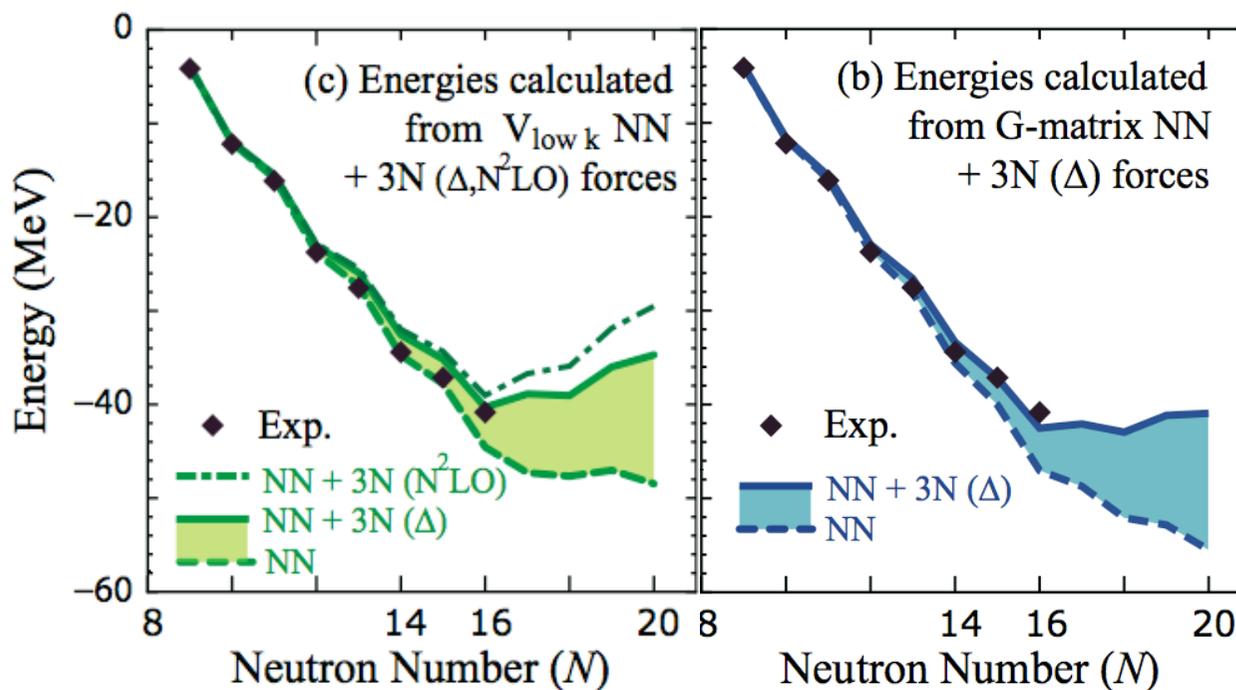
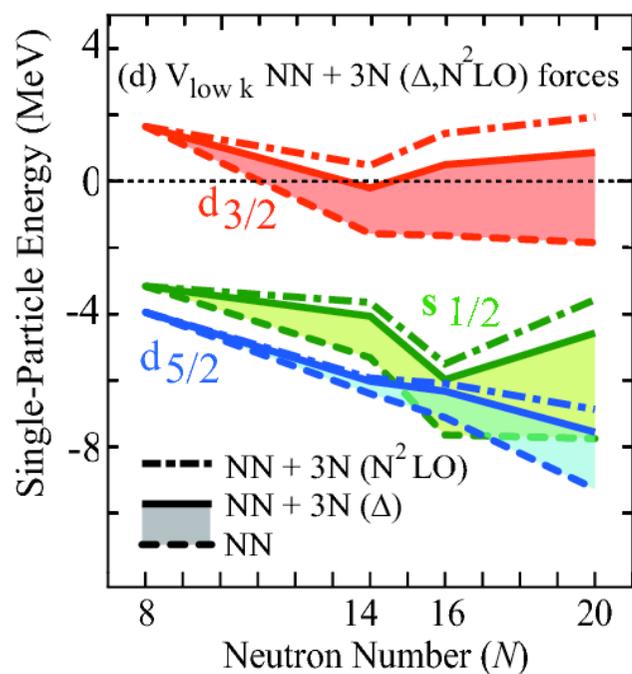
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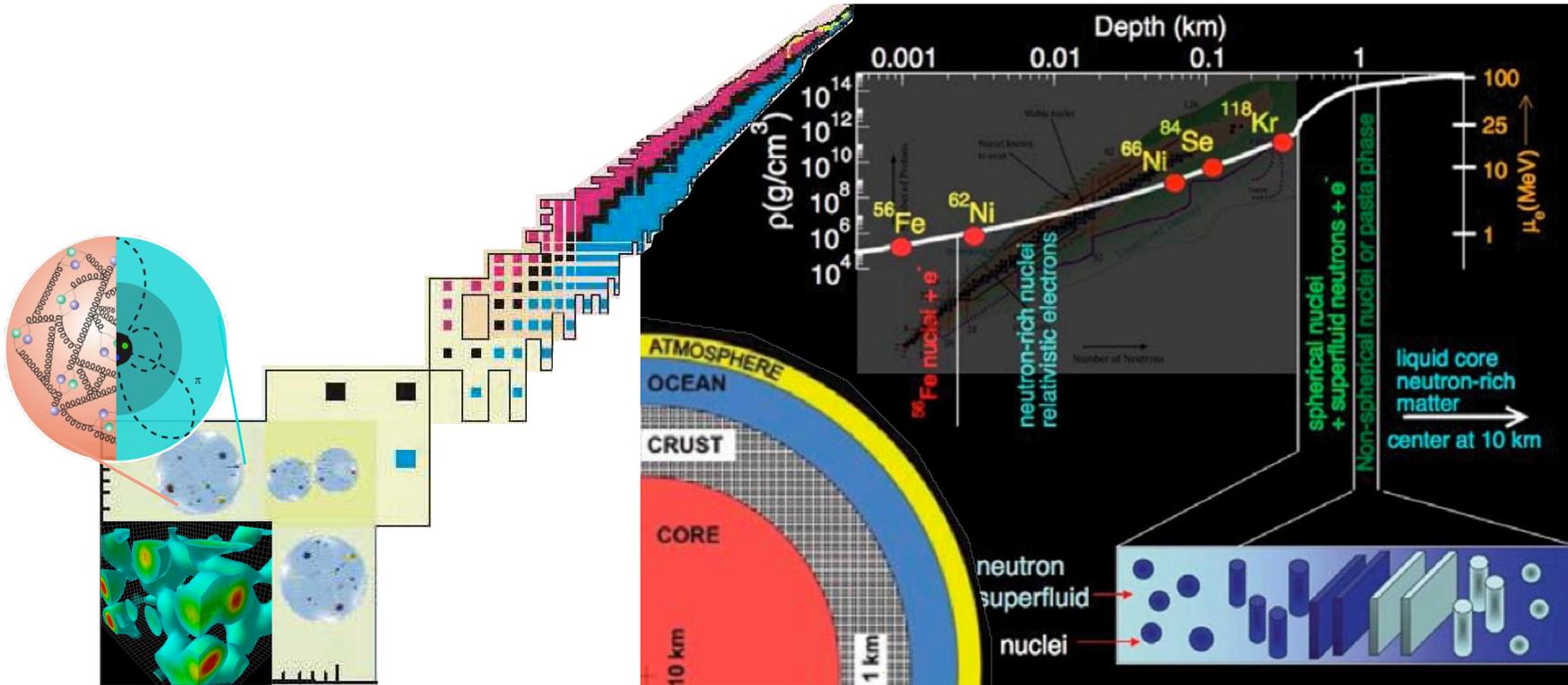


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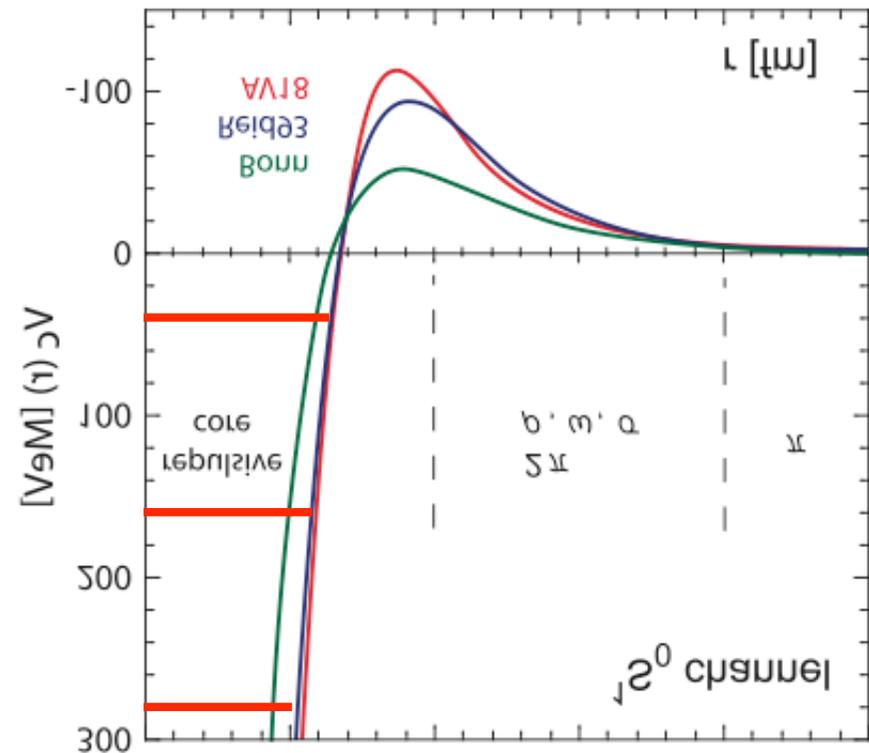
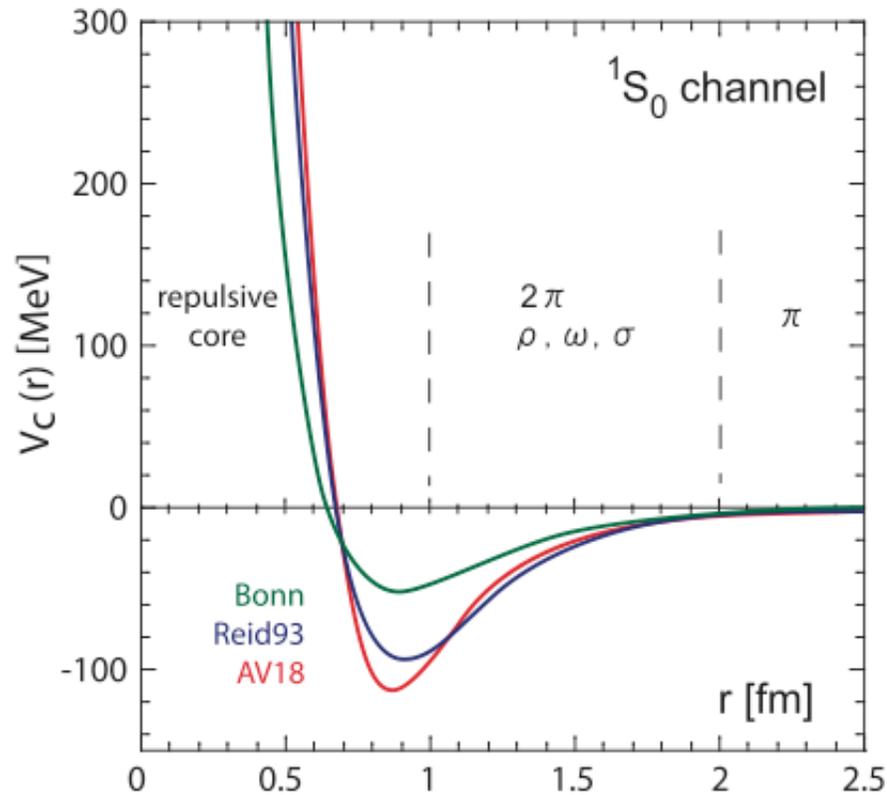
microscopic explanation of the oxygen anomaly Otsuka et al., PRL (2010)

Extreme neutron-rich matter in stars



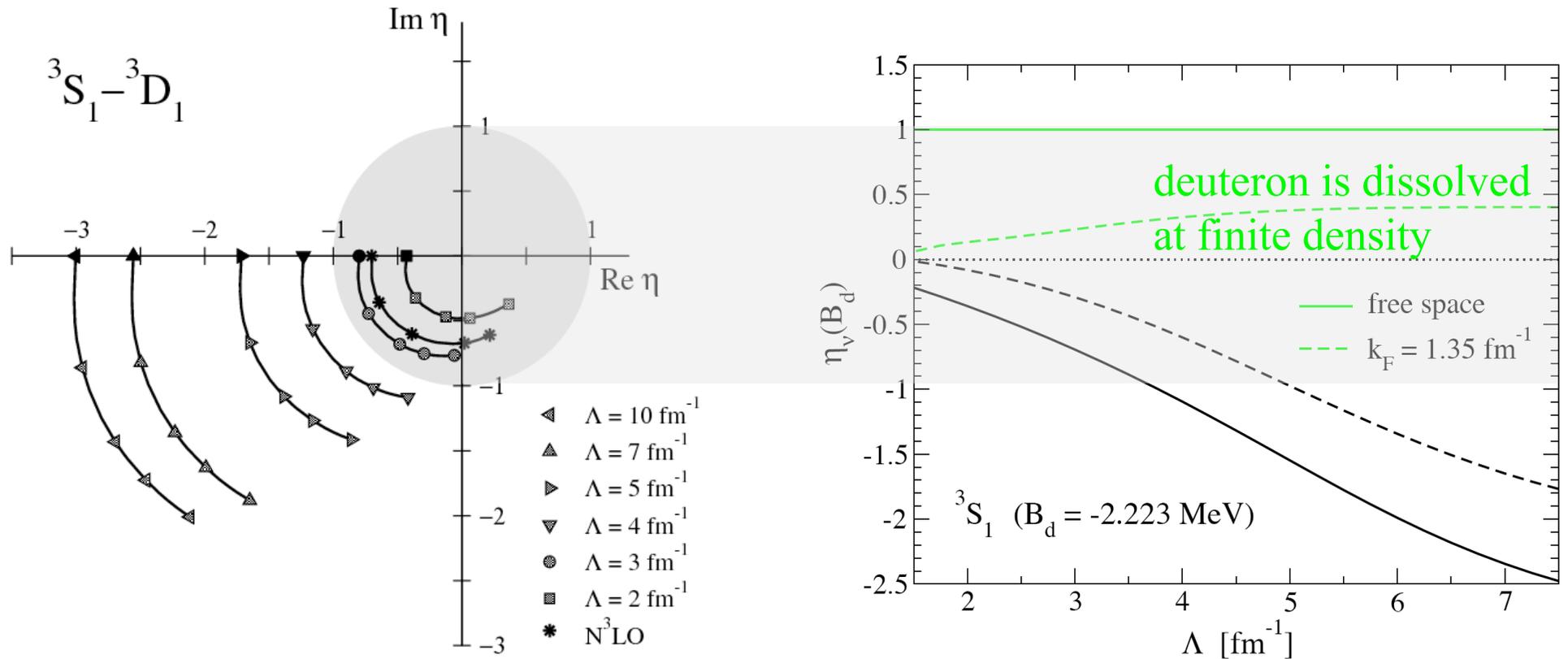
Convergence with low-momentum interactions

large cutoffs lead to **flipped-potential bound states**, even for small $-\lambda V$
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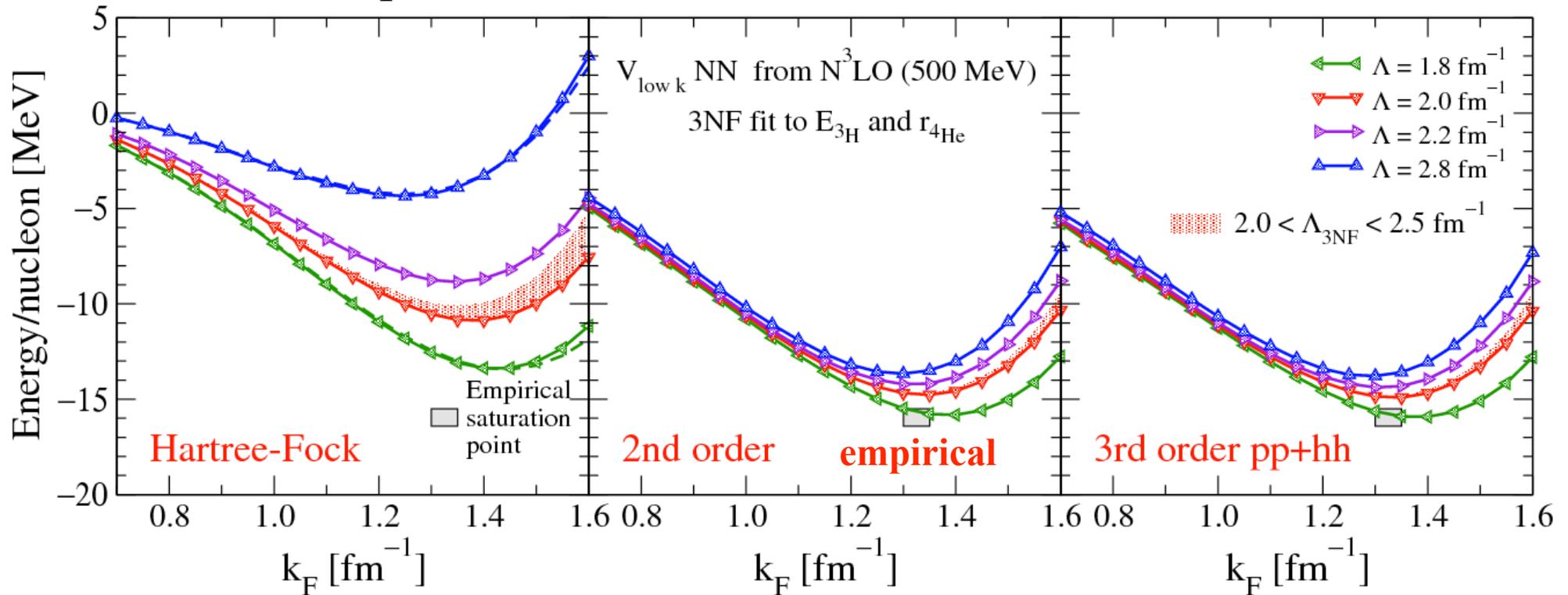


Weinberg eigenvalue analysis: two-body scattering becomes perturbative after RG evolution, except in channels with bound states

EFT and RG leads to improved convergence for nuclei and nuclear matter

Advances in nuclear matter theory

Is nuclear matter perturbative with chiral EFT and RG evolution?



Hebeler, Bogner, Furnstahl, Nogga, AS (2009, 2010)

exciting: empirical saturation with theoretical uncertainties

improved 3N treatment see also Holt, Kaiser, Weise (2010)

input to develop a universal energy density functional for all nuclei

UNEDF SciDAC Collaboration

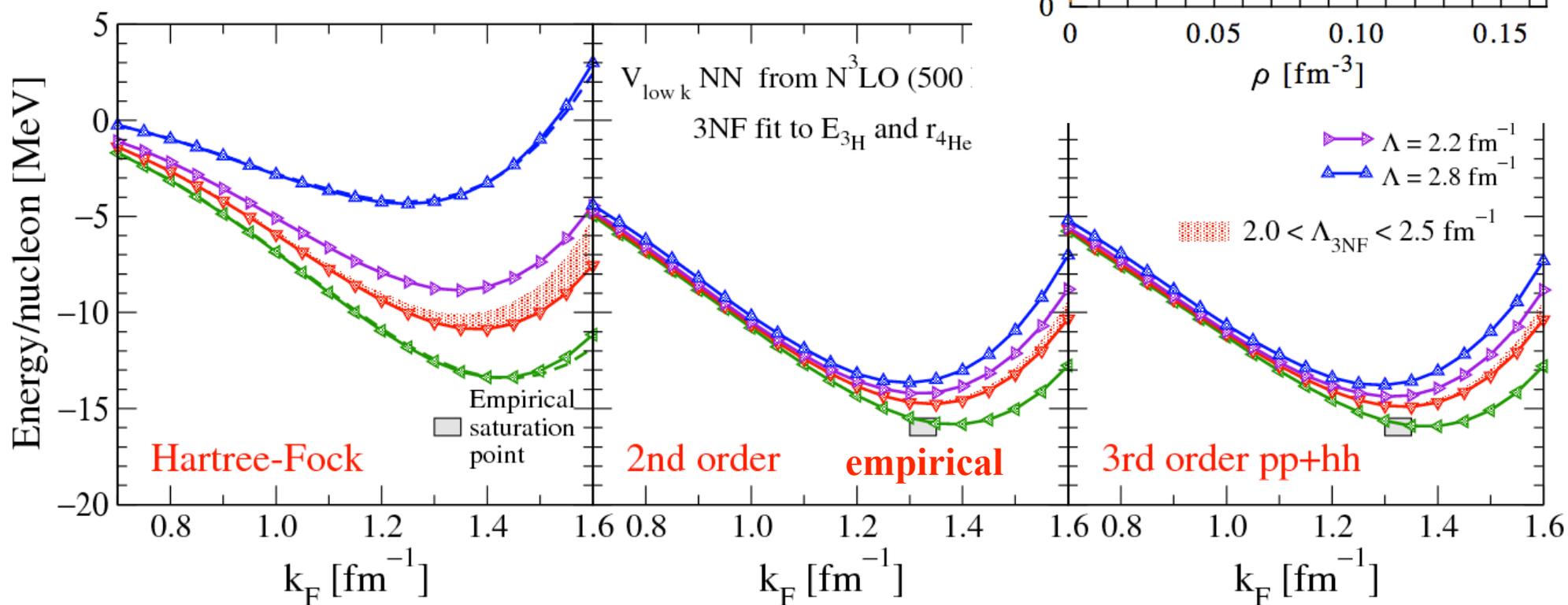
Universal Nuclear Energy Density Functional

Impact of 3N forces on neutron matter

Hebeler, AS (2010); Tolos, Friman, AS (2007)

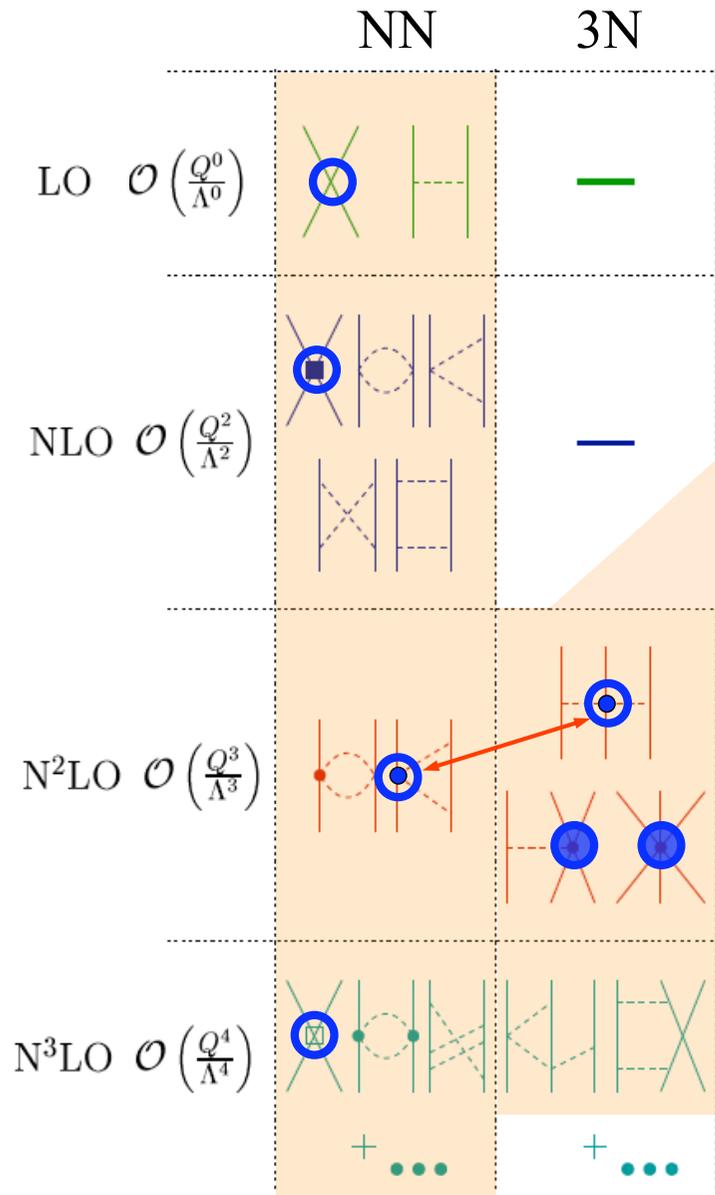
only long-range parts of 3N forces contribute to neutron matter (c_1 and c_3)

neutron matter: many-body forces are predicted to N^3LO !



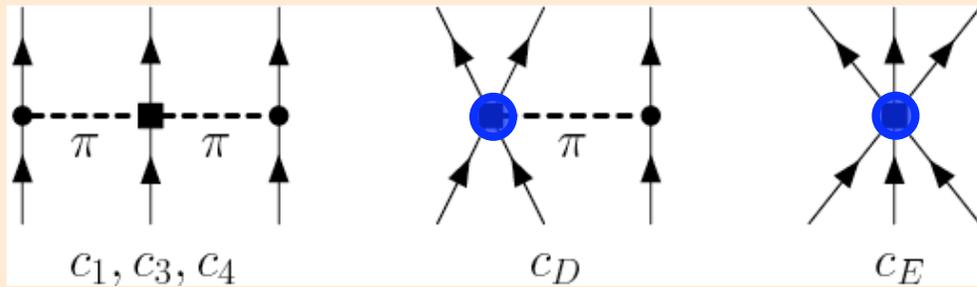
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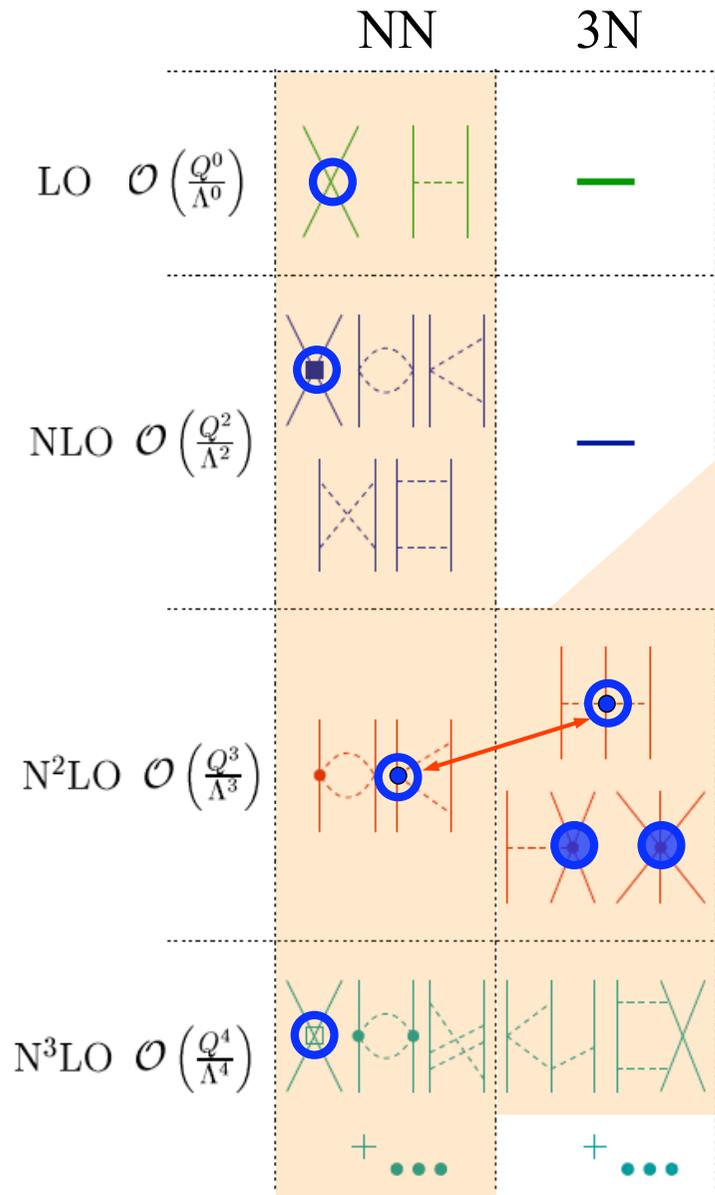
3N,4N: only 2 new couplings to N³LO



Question: Why do the c_D and c_E terms not contribute to neutron matter?

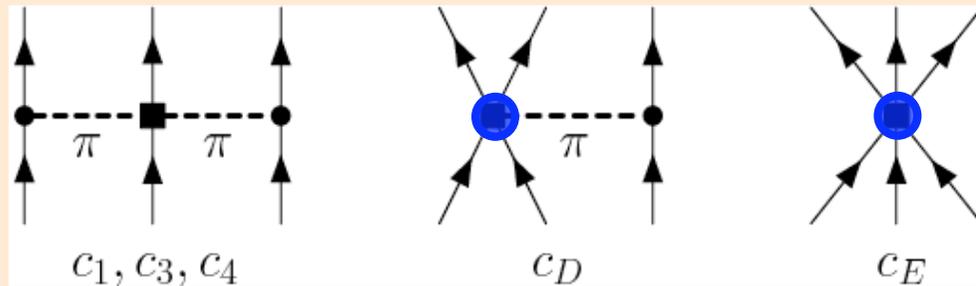
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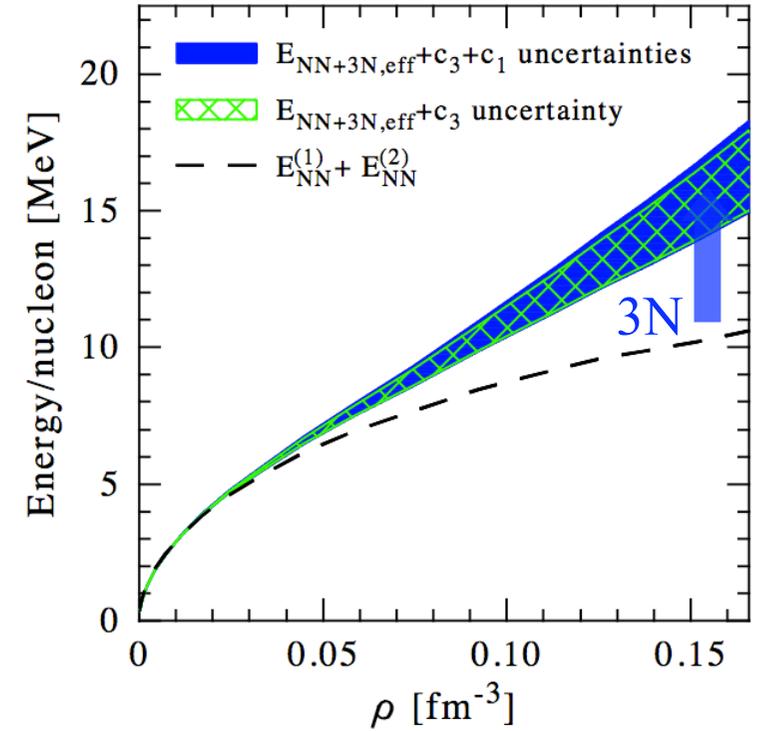
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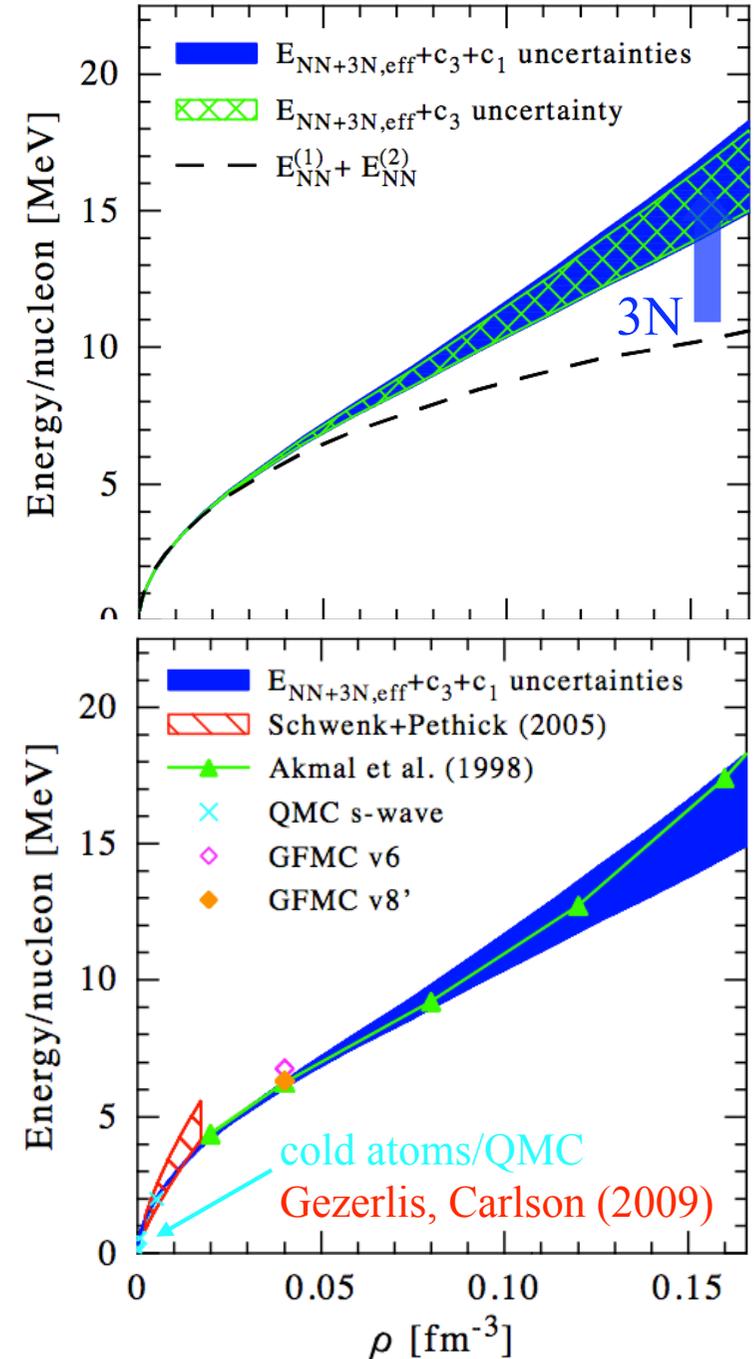


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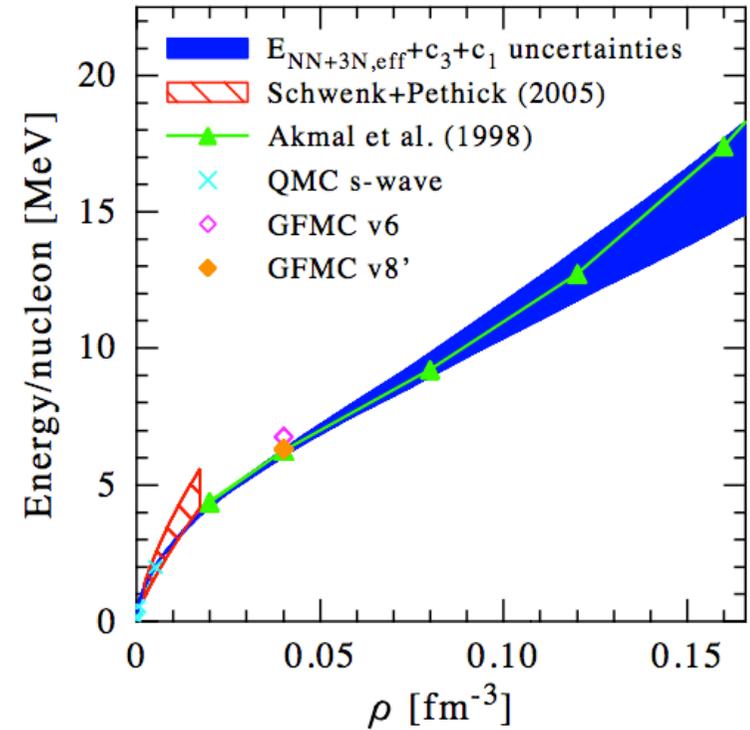
microscopic calculations within band



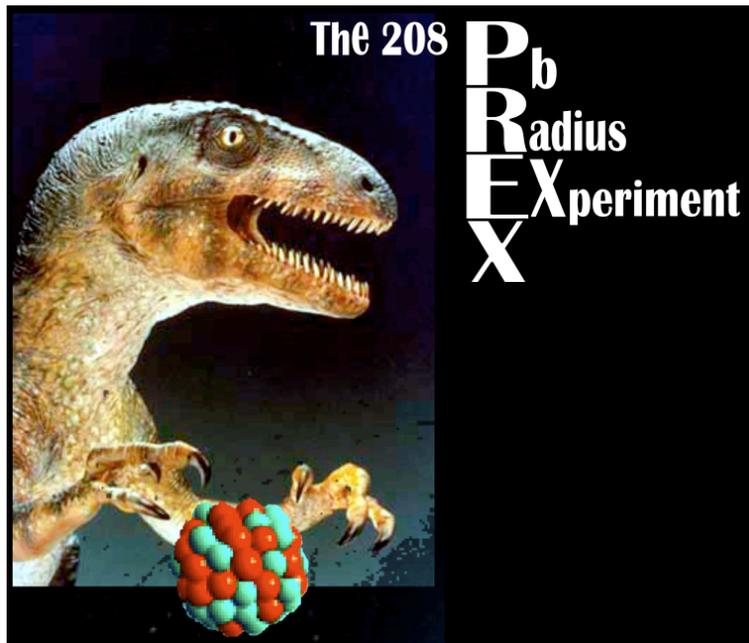
Symmetry energy and neutron skin Hebeler et al. (2010)

neutron matter band predicts range for symmetry energy 30.1-34.4 MeV

c_1 [GeV ⁻¹]	c_3 [GeV ⁻¹]	\bar{S}_2 [MeV]
-0.7	-2.2	30.1
-1.4	-4.8	34.4
NN-only EM		26.5
NN-only EGM		25.6



and neutron skin of ²⁰⁸Pb to 0.17±0.03 fm



compare to ±0.05 fm future PREX goal
first result: 0.34+0.15-0.17 fm

from complete E1 response

0.156+0.025-0.021 fm Tamii et al., PRL (2011).

Discovery of the heaviest neutron star

A two-solar-mass neutron star measured using Shapiro delay

P. B. Demorest¹, T. Pennucci², S. M. Ransom¹, M. S. E. Roberts³ & J. W. T. Hessels^{4,5}

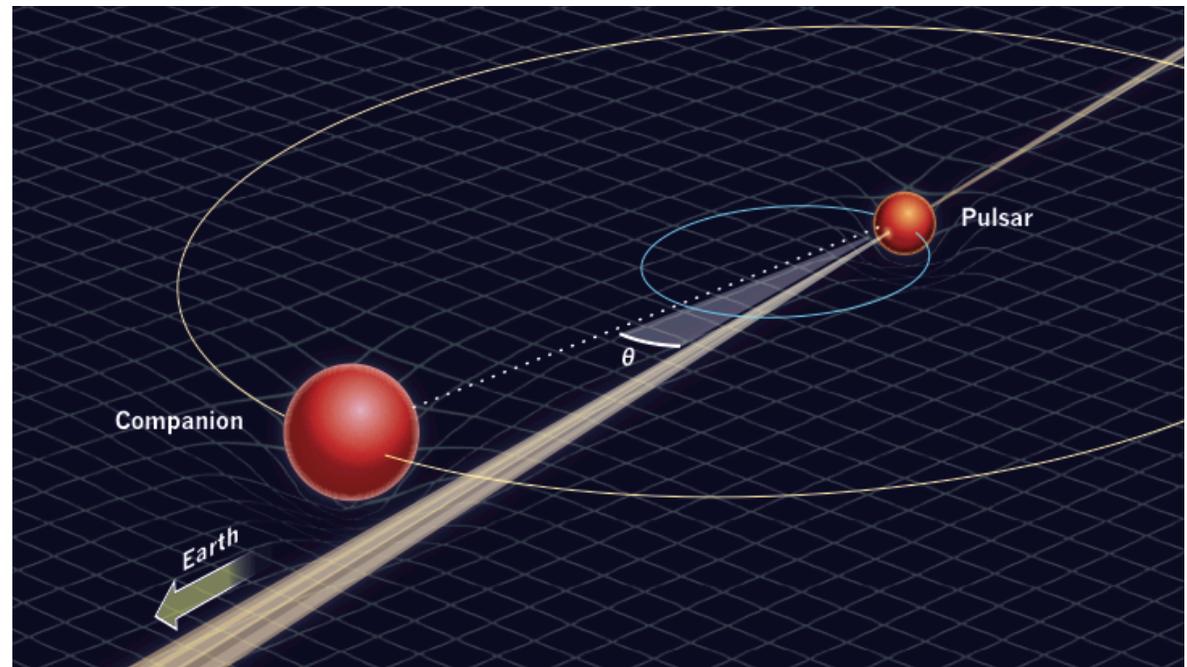
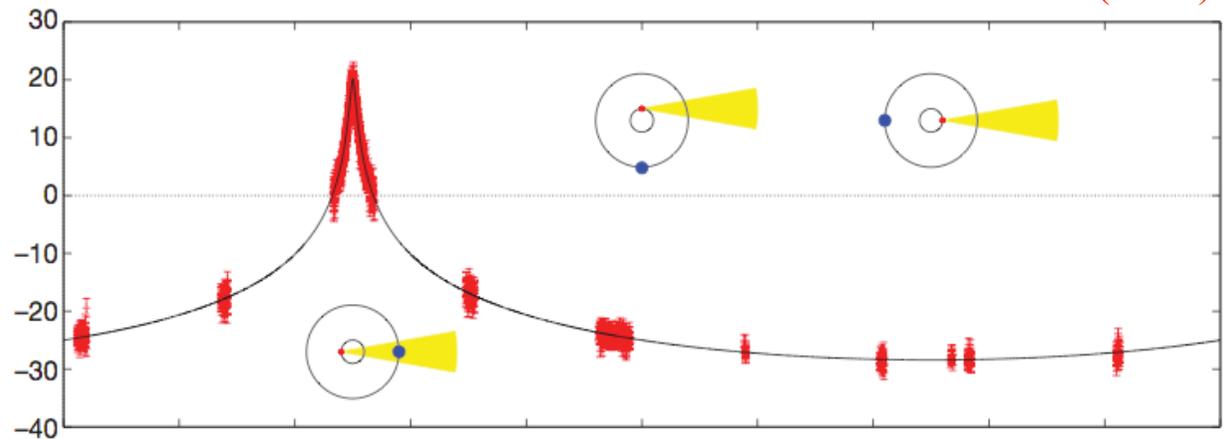
Nature (2010)

direct measurement of
neutron star mass from
increase in signal travel
time near companion

J1614-2230

most edge-on binary
pulsar known (89.17°)
+ massive white dwarf
companion ($0.5 M_{\text{sun}}$)

heaviest neutron star
with $1.97 \pm 0.04 M_{\text{sun}}$



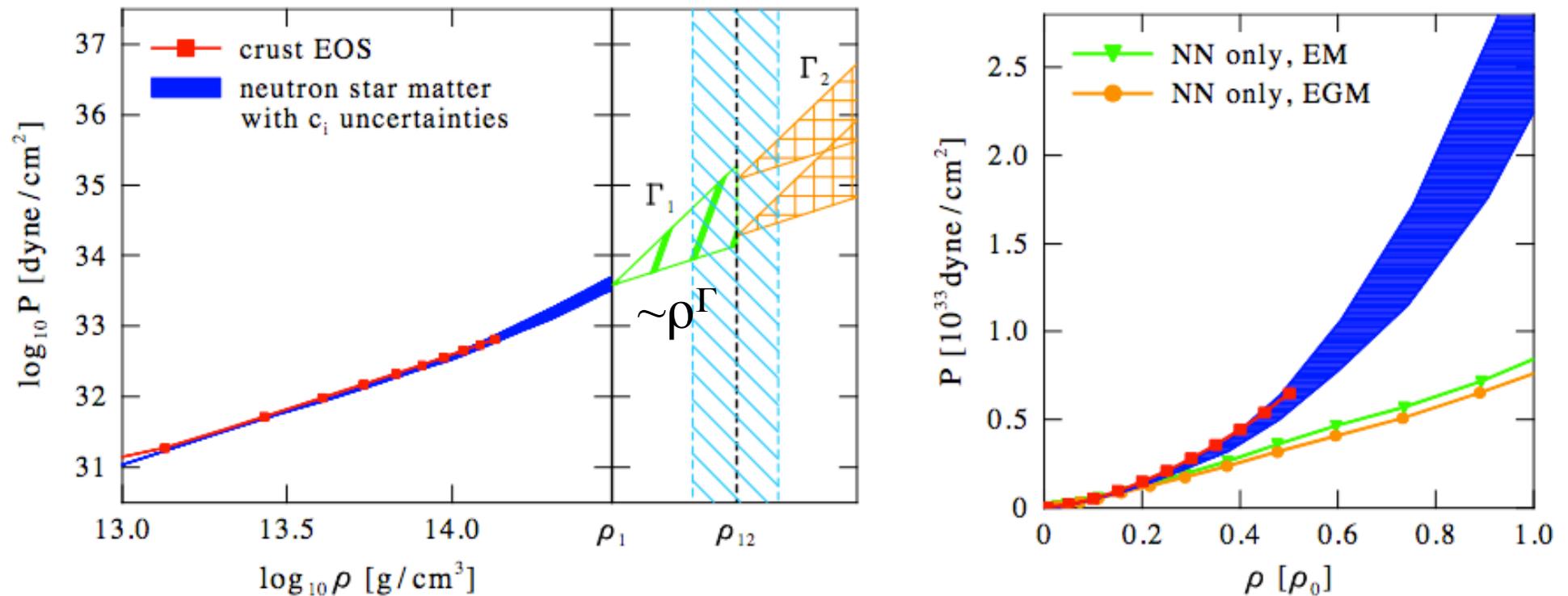
Impact on neutron stars

Neutron star structure determined by Tolman-Oppenheimer-Volkov eqn

$$\frac{dP}{dr} = -\frac{GM\epsilon}{r^2} \left[1 + \frac{P}{\epsilon c^2} \right] \left[1 + \frac{4\pi r^3 P}{Mc^2} \right] \left[1 - \frac{2GM}{c^2 r} \right]^{-1}$$

with equation of state/pressure for neutron-star matter

Hebeler et al., (2010)

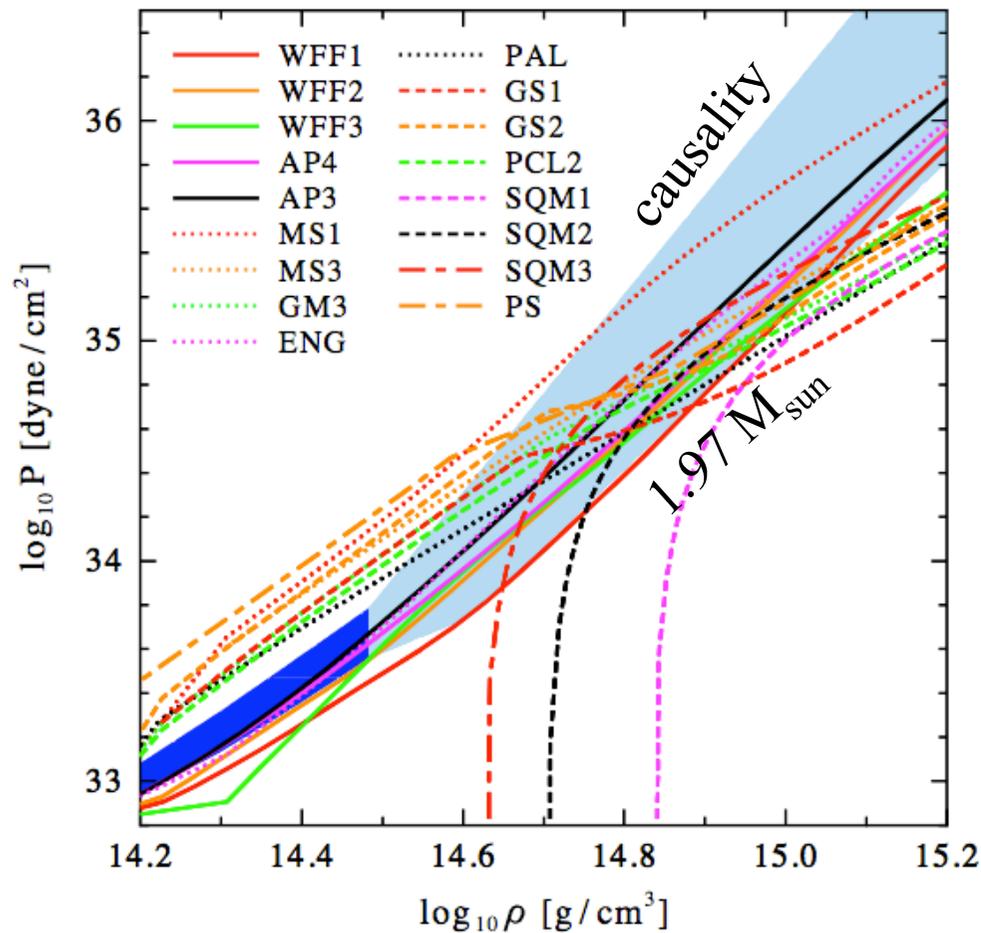


pressure below nuclear densities agrees with standard crust equation of state only after 3N forces are included

extend uncertainty band to higher densities using piecewise polytropes

Pressure of neutron star matter Hebeler et al. (2010)

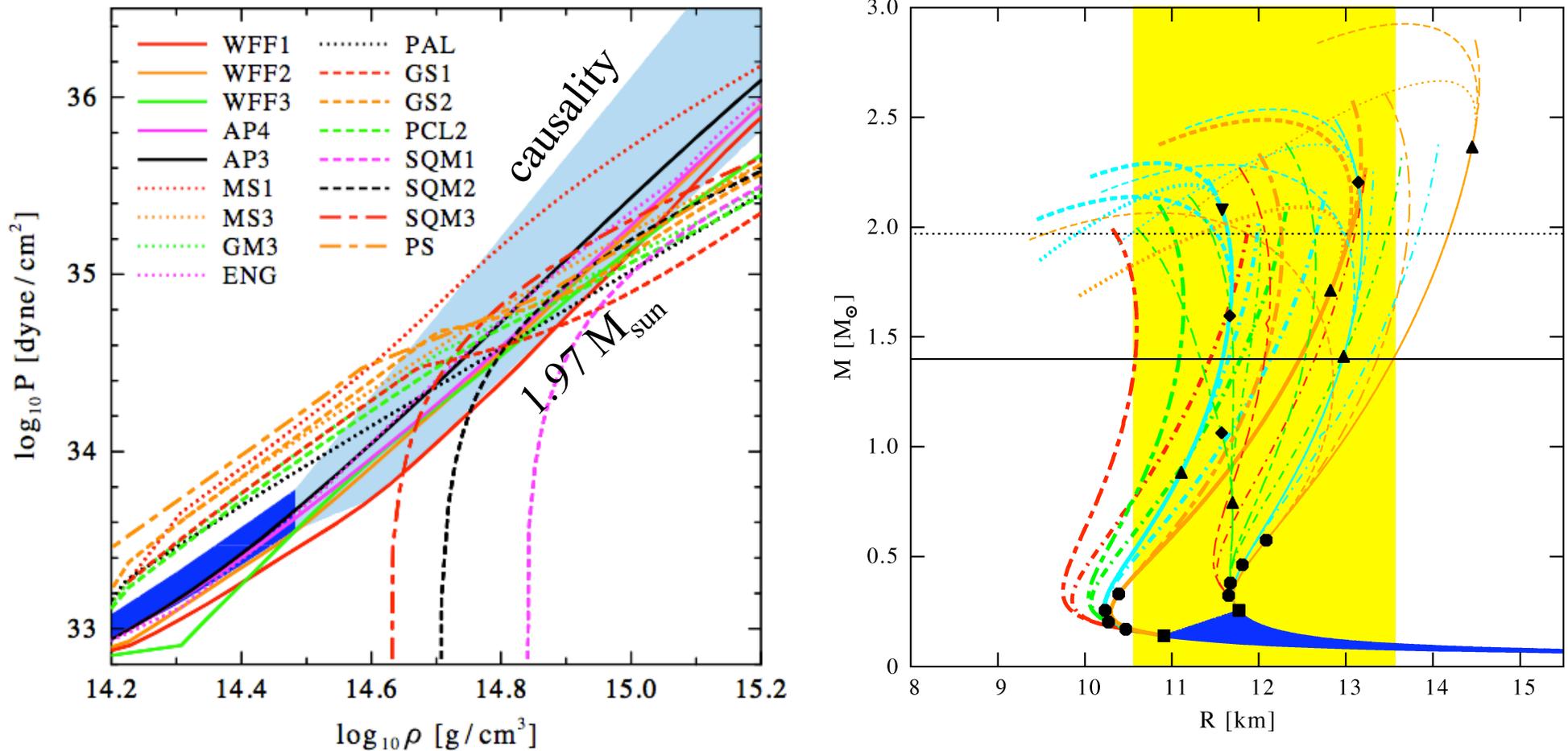
constrain polytropes by causality and require to support $1.97 M_{\text{sun}}$ star



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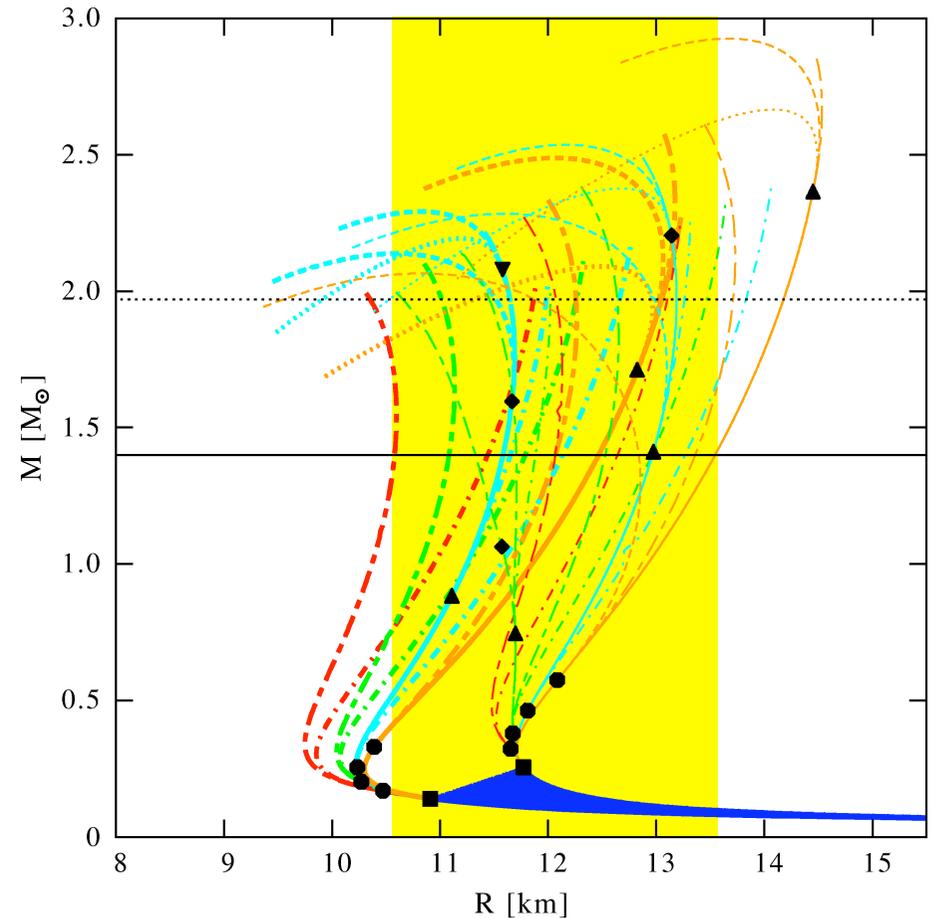
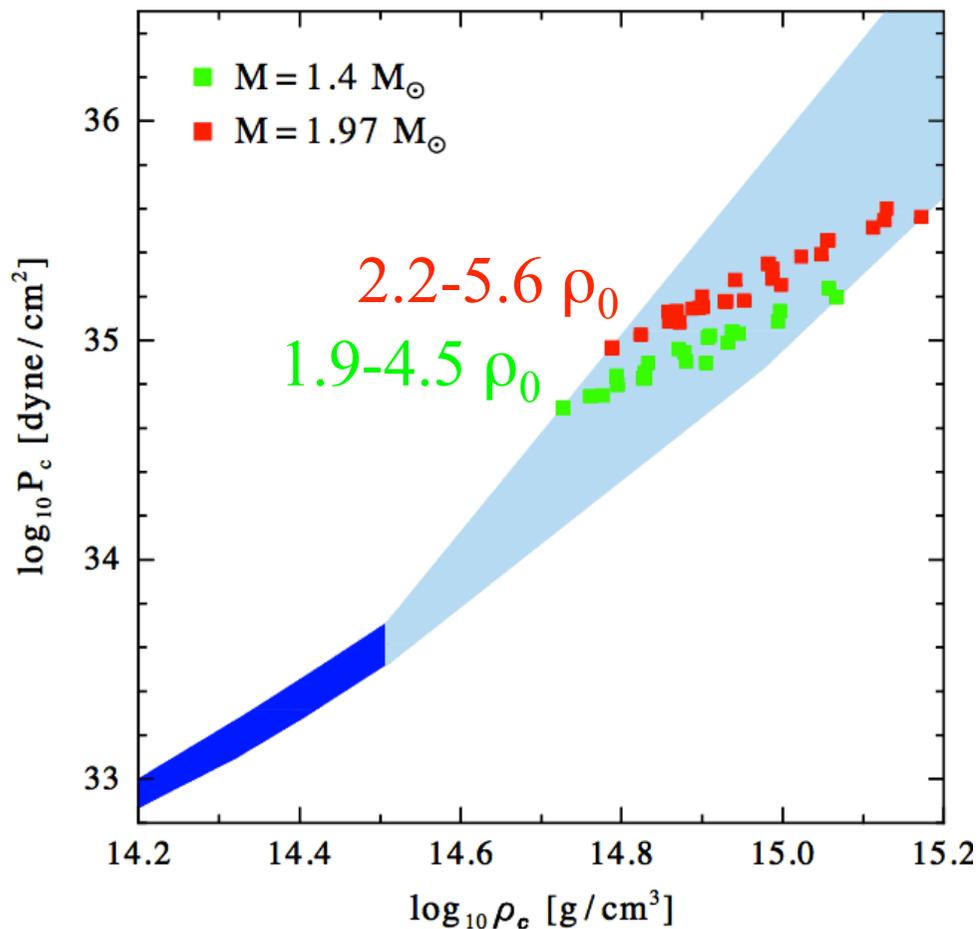


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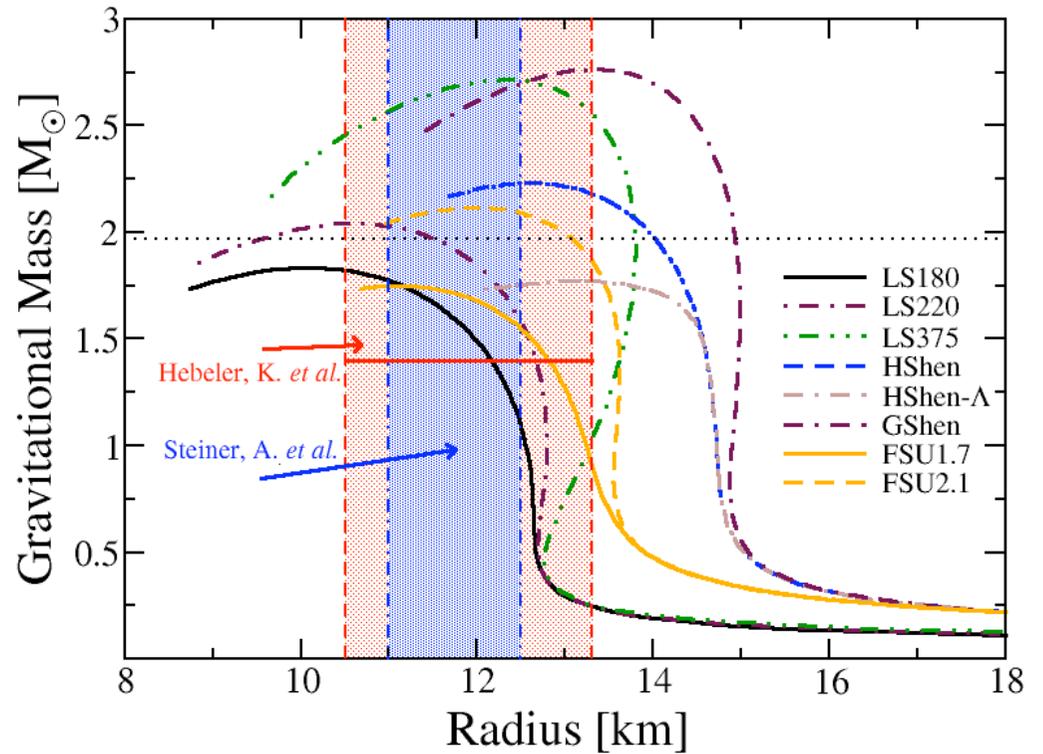
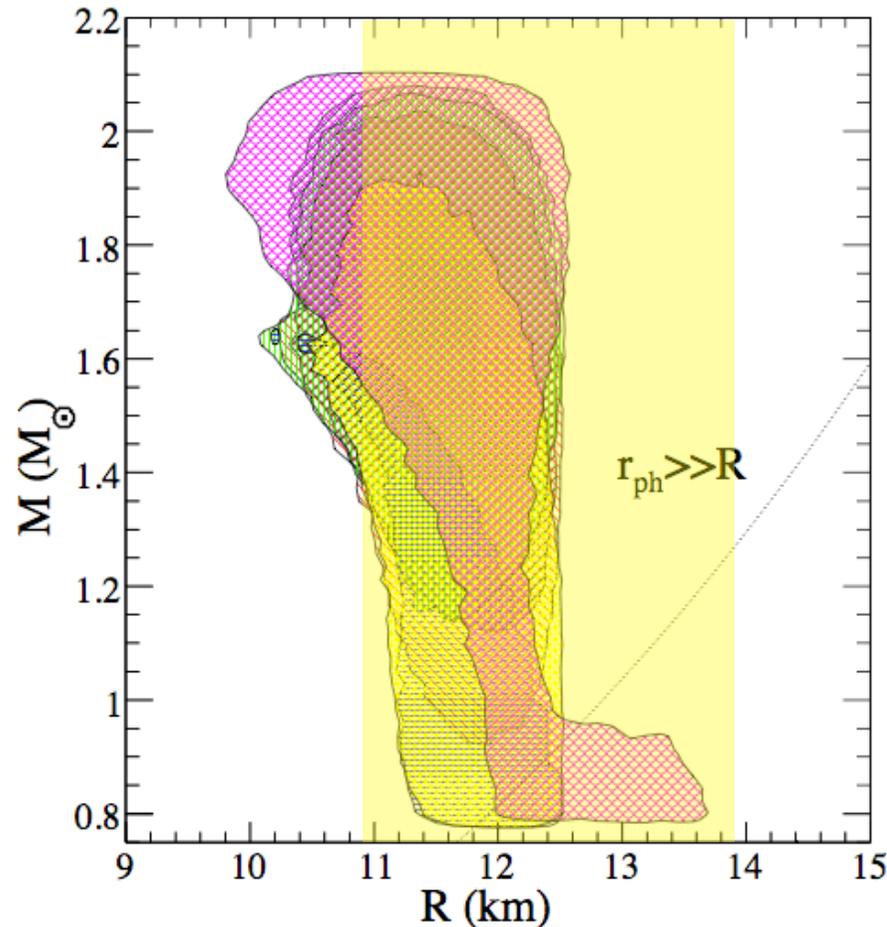


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Comparison to astrophysics

constrain polytropes by causality and require to support $1.97 M_{\text{sun}}$ star



from Evan O' Connor

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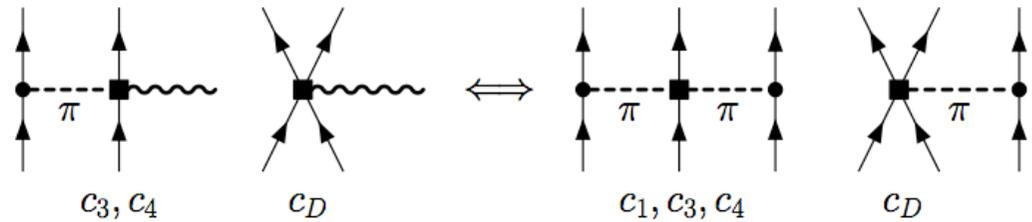
consistent with extraction from X-ray burst sources [Steiner et al., ApJ \(2010\)](#)

provides important constraints for EOS for core-collapse supernovae

Chiral EFT for electroweak transitions Menendez, Gazit, AS (2011).

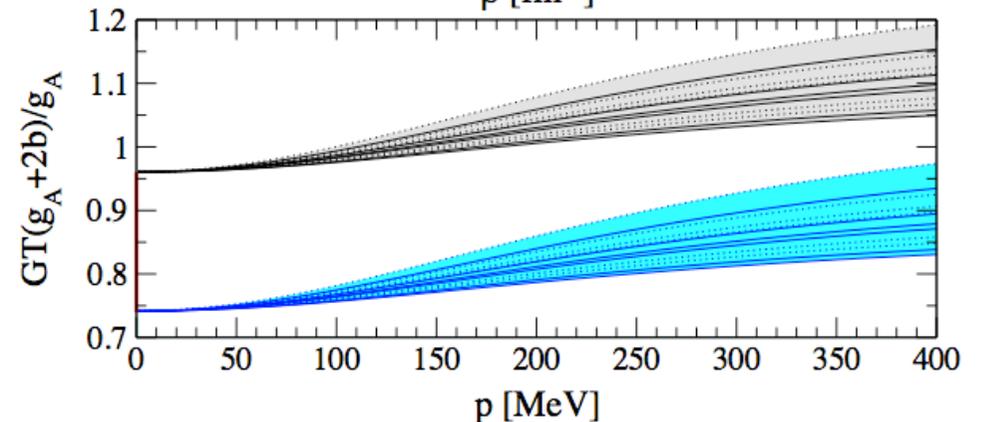
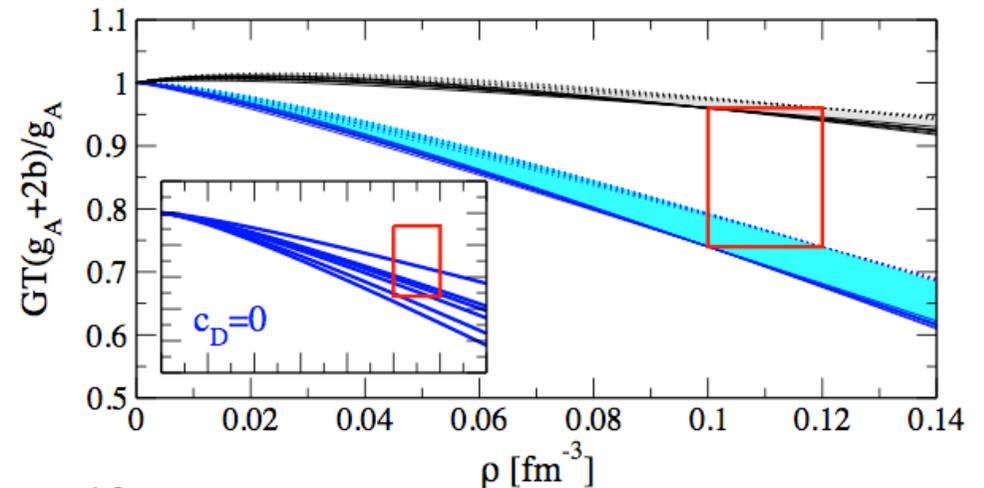
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two-body currents determined
by NN, 3N couplings to N^3 LO
Park et al., Phillips,...



explains part of quenching of g_A
(dominated by long-range part)

+ predict momentum dependence
(weaker quenching for larger p)



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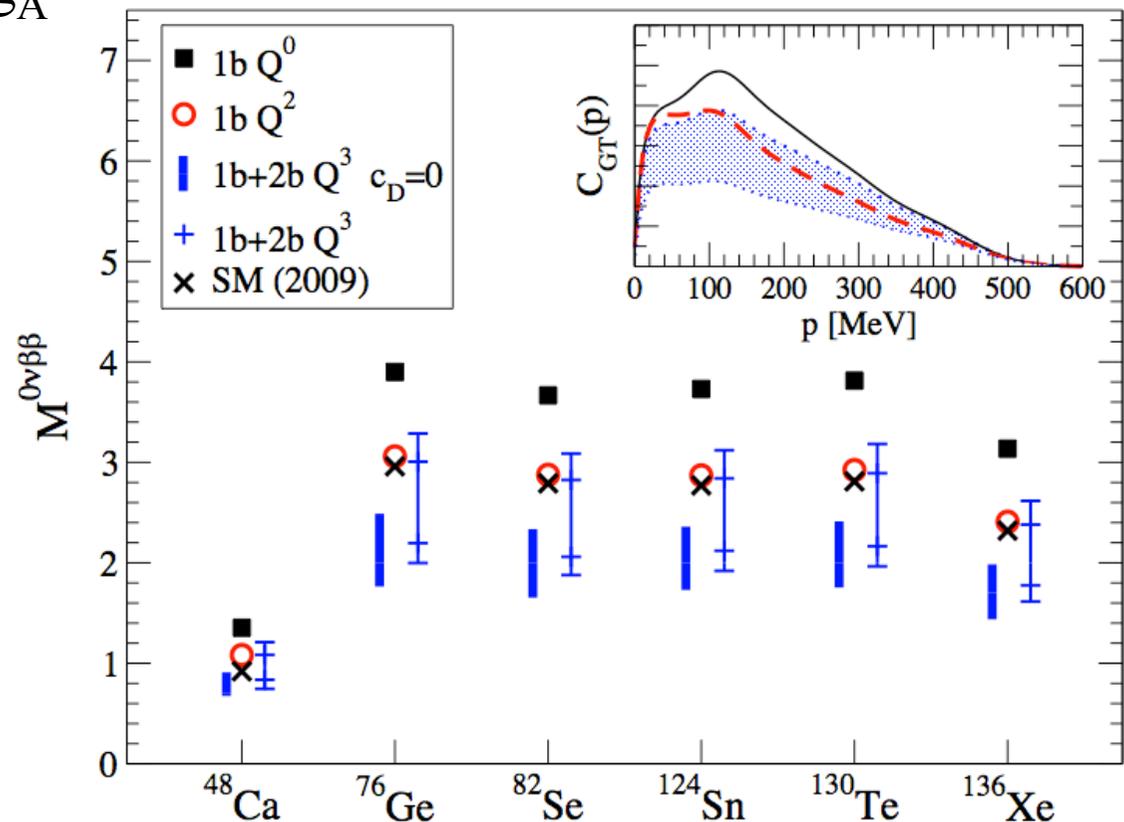
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Thanks to collaborators!



T. Krüger, J. Menendez, V. Soma, I. Tews



S.K. Bogner



R.J. Furnstahl,
K. Hebeler



A. Nogga

J.D. Holt

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Summary

Exciting era with advances on many fronts:
development of effective field theory and the renormalization group

enables a unified description from nuclei to matter in astrophysics

3N forces are a frontier for neutron-rich nuclei/matter:

key to explain why ^{24}O is the heaviest oxygen isotope

Ca isotopes and $N=28$ magic number, key for neutron-rich nuclei

dominant uncertainty of neutron (star) matter below nuclear densities,
constraints on neutron star radii

exciting interactions with experiments and observations!