

What is New in Quarkonium Production at RHIC and LHC ?

Pengfei ZHUANG, Tsinghua University, Beijing

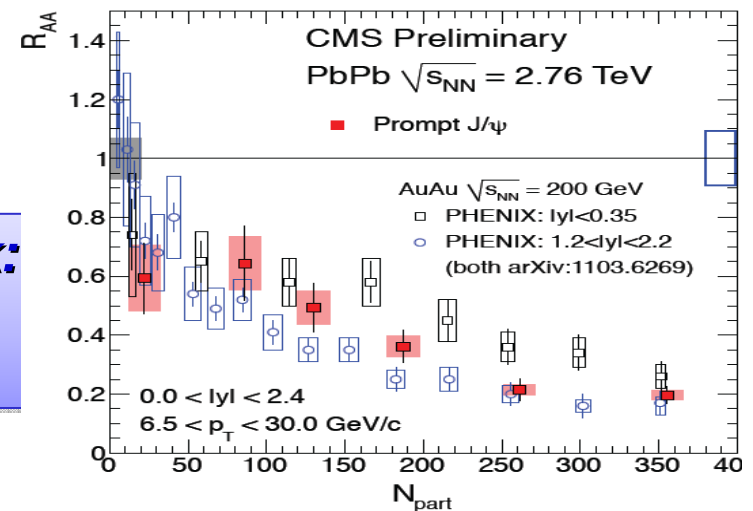
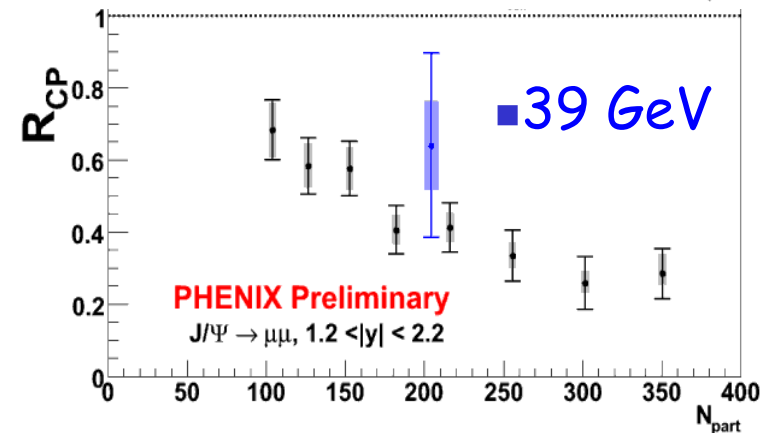
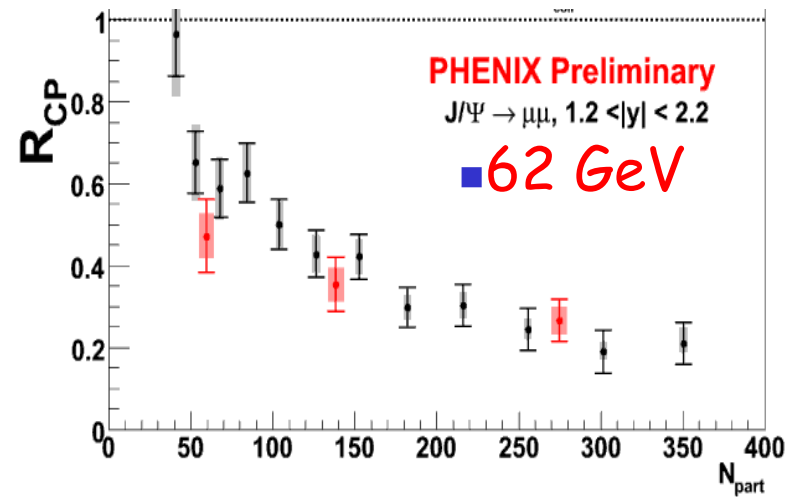
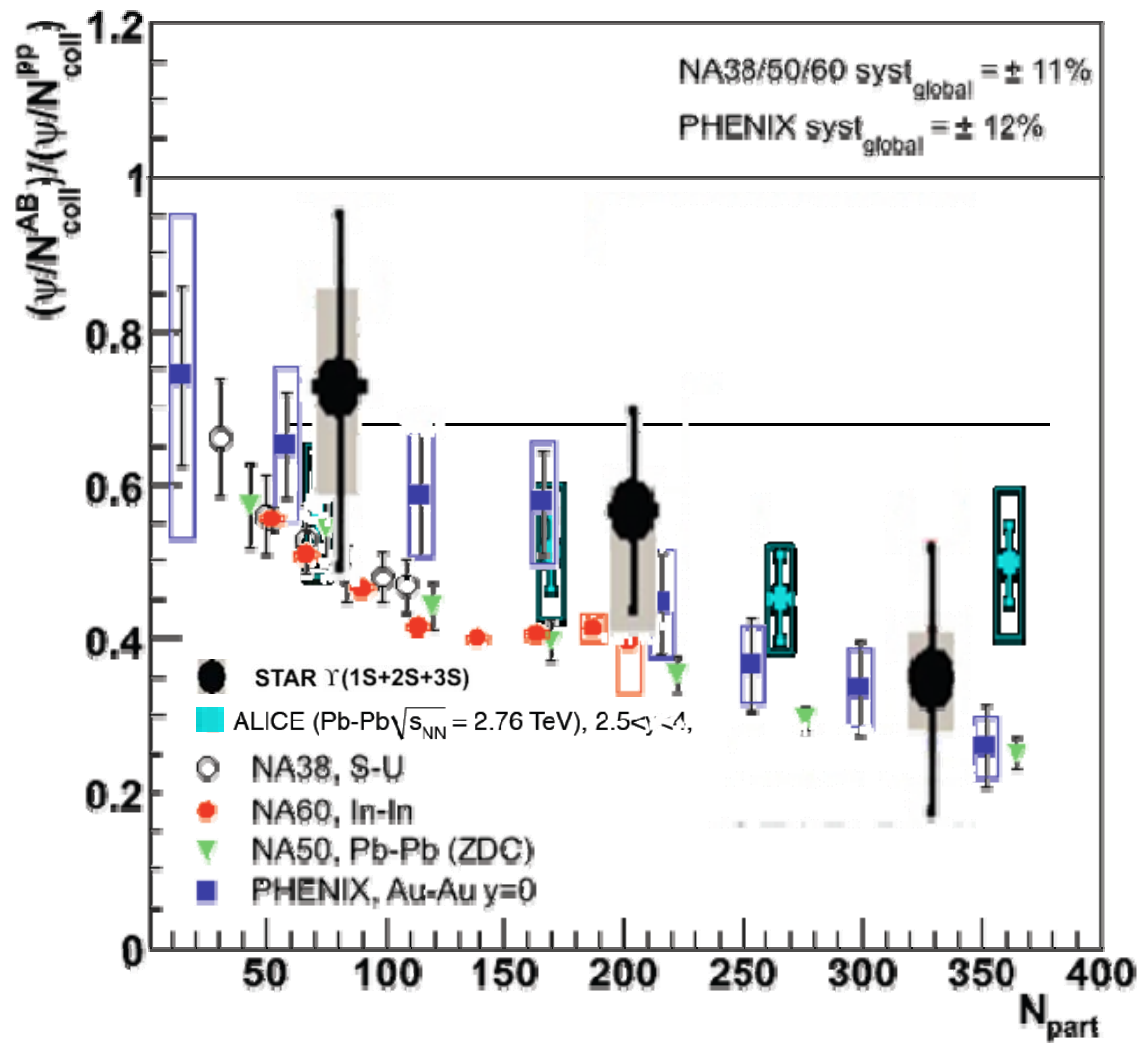
- ***Introduction***
- ***A Full Transport Approach***
- ***P_t Distribution and V_2 Distribution***
- ***Υ Production and Heavy Quark Potential***
- ***Summary***

Thanks to Yunpeng LIU, Kai ZHOU and Dr. Nu XU

EMMI Workshop: Quarkonia in Deconfined Matter

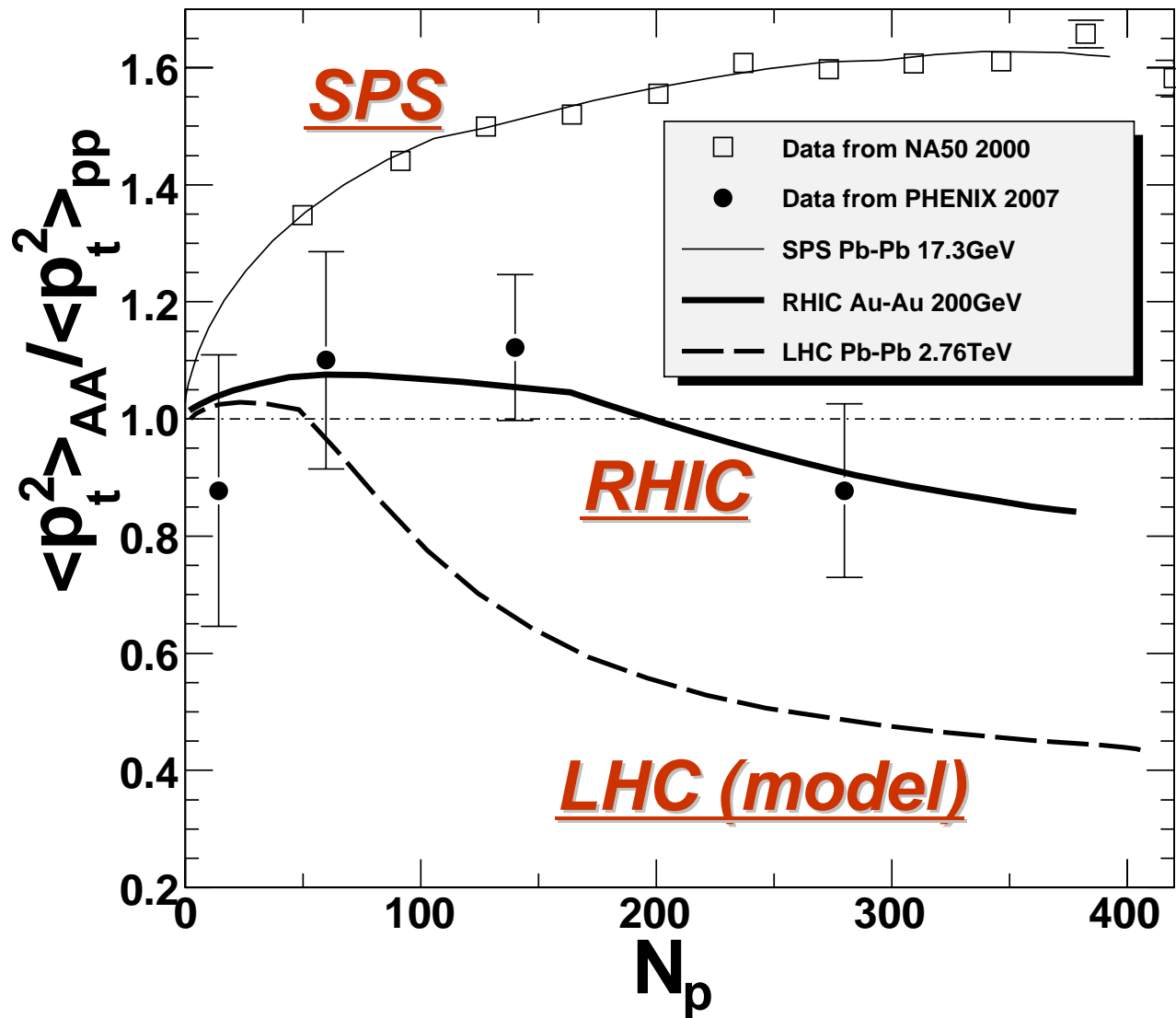
Acitrezza, Italy, September 28-30, 2011

Introduction



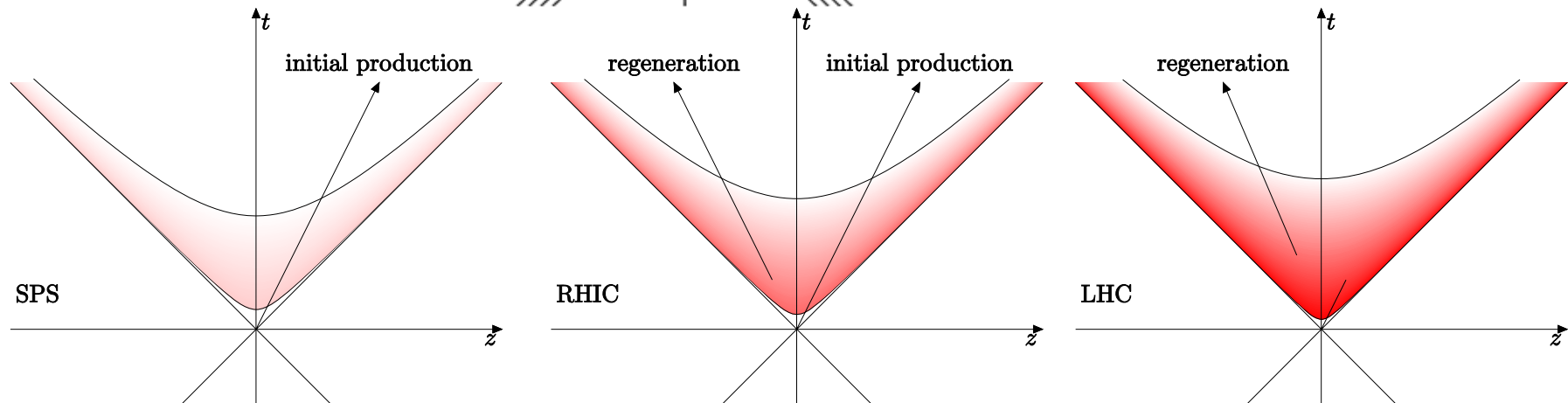
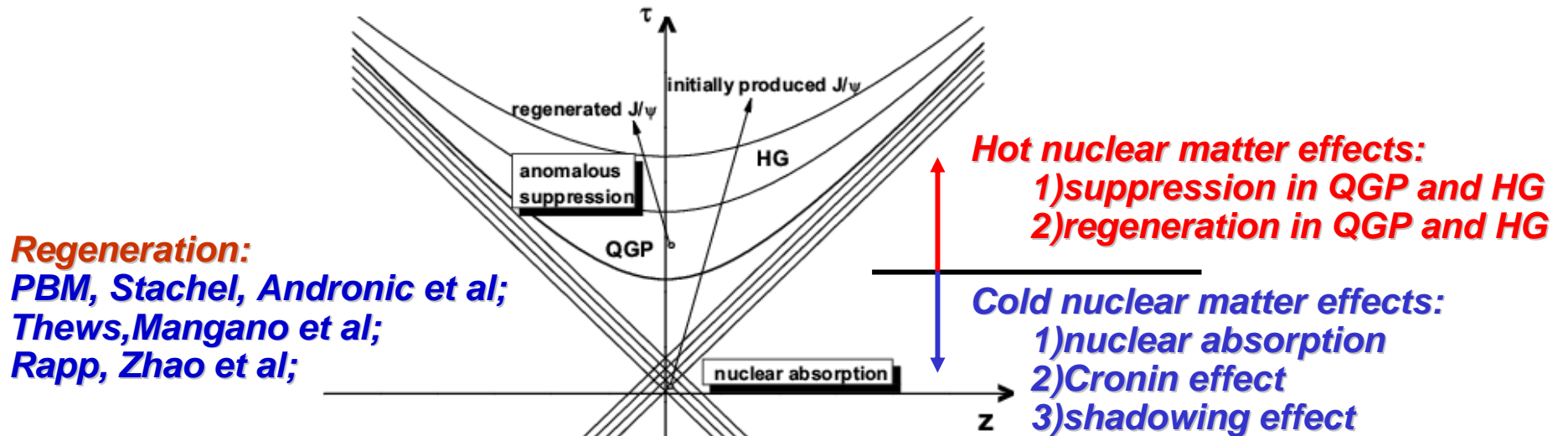
SQM2011 experimental summary by K.Safarik:
overall suppression of J/ψ is nearly identical
between SPS, RHIC & LHC !

However, the transverse momentum distribution is sensitive to the quarkonium production mechanism !



A Full Transport Approach

Quarkonium Production Picture



screening only, screening + regeneration

- initial production controls high pt region
- regeneration becomes important at low pt

A Full Transport Approach for Quarkonia in HIC

Yan, Zhu, Xu, Zhuang, 2005-2011

• QGP hydrodynamics

$$\begin{aligned}\partial_\tau E + \nabla \cdot \mathbf{M} &= -(E + p)/\tau, \\ \partial_\tau M_x + \nabla \cdot (M_x \mathbf{v}) &= -M_x/\tau - \partial_x p, \\ \partial_\tau M_y + \nabla \cdot (M_y \mathbf{v}) &= -M_y/\tau - \partial_y p, \\ \partial_\tau R + \nabla \cdot (R\mathbf{v}) &= -R/\tau\end{aligned}$$

+ equation of state (ideal gas or strongly coupled matter from lattice)

• quarkonium transport equations ($\Psi = J/\psi, \psi', \chi_c$)

$$\partial f_\Psi / \partial \tau + \mathbf{v}_\Psi \cdot \nabla f_\Psi = -\alpha_\Psi f_\Psi + \beta_\Psi. \quad \alpha: \text{suppression} \quad \beta: \text{regeneration}$$

$$\alpha_\Psi(\mathbf{p}_t, \mathbf{x}_t, \tau | \mathbf{b}) = \frac{1}{2E_\Psi} \int \frac{d^3 \mathbf{p}_g}{(2\pi)^3 2E_g} W_{g\Psi}^{c\bar{c}}(s) f_g(\mathbf{p}_g, \mathbf{x}_t, \tau) \Theta(T(\mathbf{x}_t, \tau | \mathbf{b}) - T_c),$$

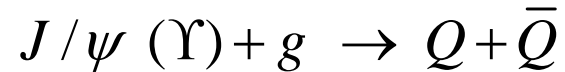
$$\begin{aligned}\beta_\Psi(\mathbf{p}_t, \mathbf{x}_t, \tau | \mathbf{b}) &= \frac{1}{2E_\Psi} \int \frac{d^3 \mathbf{p}_g}{(2\pi)^3 2E_g} \frac{d^3 \mathbf{p}_c}{(2\pi)^3 2E_c} \frac{d^3 \mathbf{p}_{\bar{c}}}{(2\pi)^3 2E_{\bar{c}}} W_{c\bar{c}\Psi}^{g\Psi}(s) f_c(\mathbf{p}_c, \mathbf{x}_t, \tau | \mathbf{b}) f_{\bar{c}}(\mathbf{p}_{\bar{c}}, \mathbf{x}_t, \tau | \mathbf{b}) \\ &\times (2\pi)^4 \delta^{(4)}(p + p_g - p_c - p_{\bar{c}}) \Theta(T(\mathbf{x}_t, \tau | \mathbf{b}) - T_c),\end{aligned}$$

• analytic solution

$$\begin{aligned}f_\Psi(\mathbf{p}_t, \mathbf{x}_t, \tau | \mathbf{b}) &= f_\Psi(\mathbf{p}_t, \mathbf{x}_t - \mathbf{v}_\Psi(\tau - \tau_0), \tau_0 | \mathbf{b}) e^{-\int_{\tau_0}^{\tau} d\tau' \alpha_\Psi(\mathbf{p}_t, \mathbf{x}_t - \mathbf{v}_\Psi(\tau - \tau'), \tau' | \mathbf{b})} \\ &+ \int_{\tau_0}^{\tau} d\tau' \beta_\Psi(\mathbf{p}_t, \mathbf{x}_t - \mathbf{v}_\Psi(\tau - \tau'), \tau' | \mathbf{b}) e^{-\int_{\tau'}^{\tau} d\tau'' \alpha_\Psi(\mathbf{p}_t, \mathbf{x}_t - \mathbf{v}_\Psi(\tau - \tau''), \tau'' | \mathbf{b})}.\end{aligned}$$

• initial production $f(\vec{p}, \vec{x}, t_0)$ including CNM, correlation between the cold and hot nuclear matter effect

Dissociation Cross Section



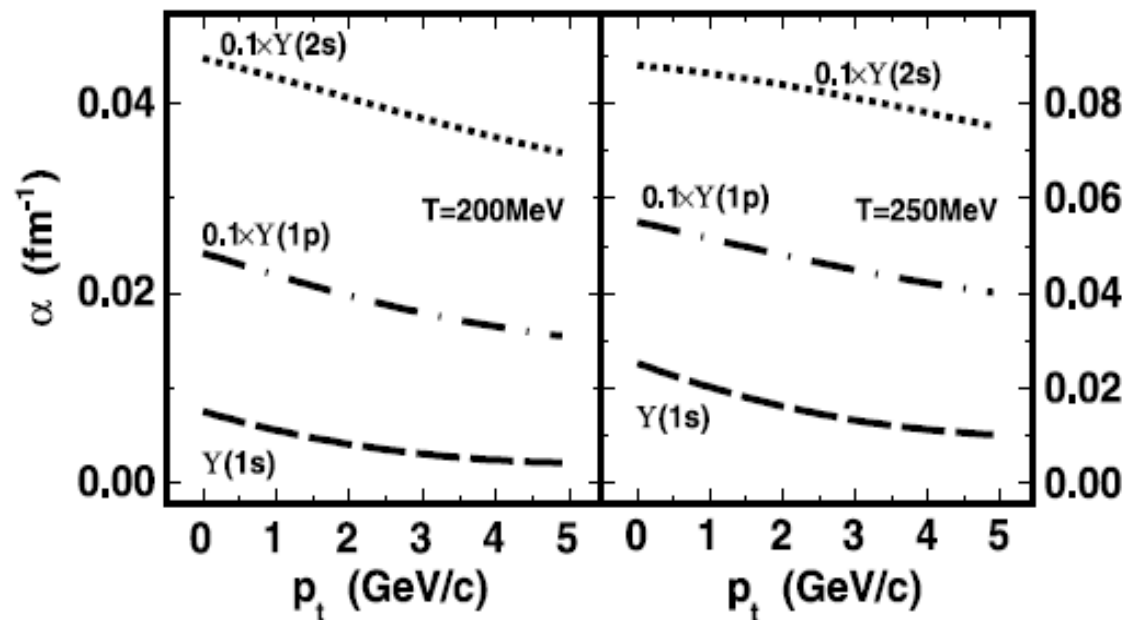
- **gluon dissociation cross section described by OPE (Bhanot, Peskin, 1999):**

$$\sigma(p_\psi, p_g)$$

- **at finite temperature, we use the classical relation**

$$\sigma(p_\psi, p_g, T) \propto \frac{\langle r^2 \rangle(T)}{\langle r^2 \rangle(0)} \sigma(p_\psi, p_g) \quad \langle r^2 \rangle(T) \text{ is calculated through the Schroedinger equation}$$

- **dissociation rate of Υ at a fixed medium velocity $v=0.5$ and for $V=U$:**



- 1) **T -dependence of the differential cross section is still an open problem;**
- 2) **we did not consider quasi-free processes which may play an important role at high $T < T_d$.**

- **regeneration rate is determined by the detailed balance**

Dissociation Length $\langle r^2 \rangle(T)$ in Potential Model

Lattice calculated free energy F for a pair of $Q\bar{Q}$

two limits of the potential:

$$V(r,T) = \begin{cases} F, & \text{slow dissociation} \\ U = F + TS, & \text{rapid dissociation} \end{cases}$$

Schroedinger equation at finite T :

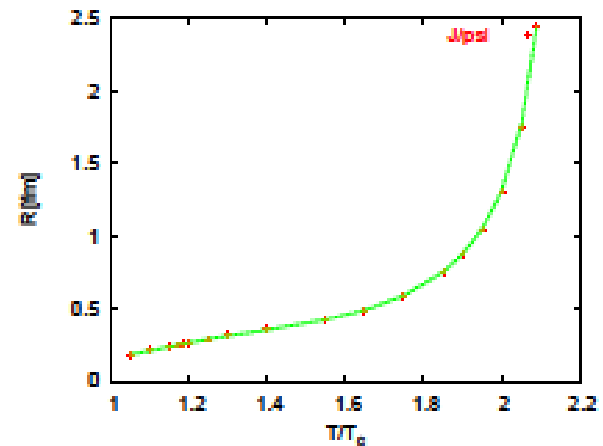
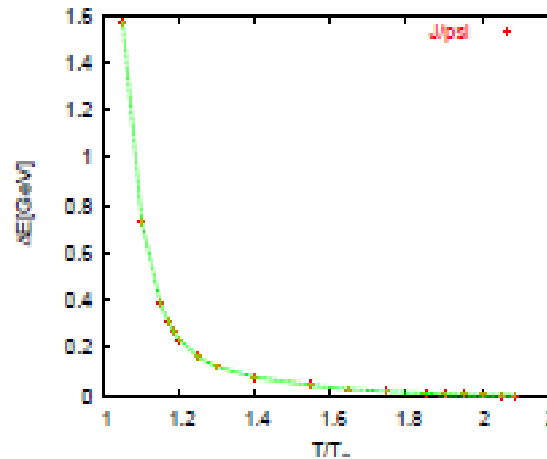
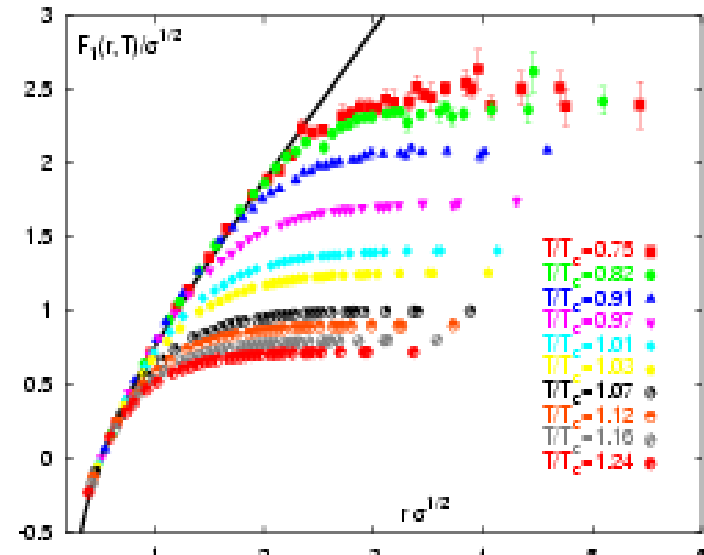
average distance $\langle r \rangle(T)$

binding energy $\varepsilon(T)$

dissociation temperature

$$\langle r \rangle(T_D) \rightarrow \infty, \quad \varepsilon(T_D) \rightarrow 0$$

for $V=U$ (Satz et al)

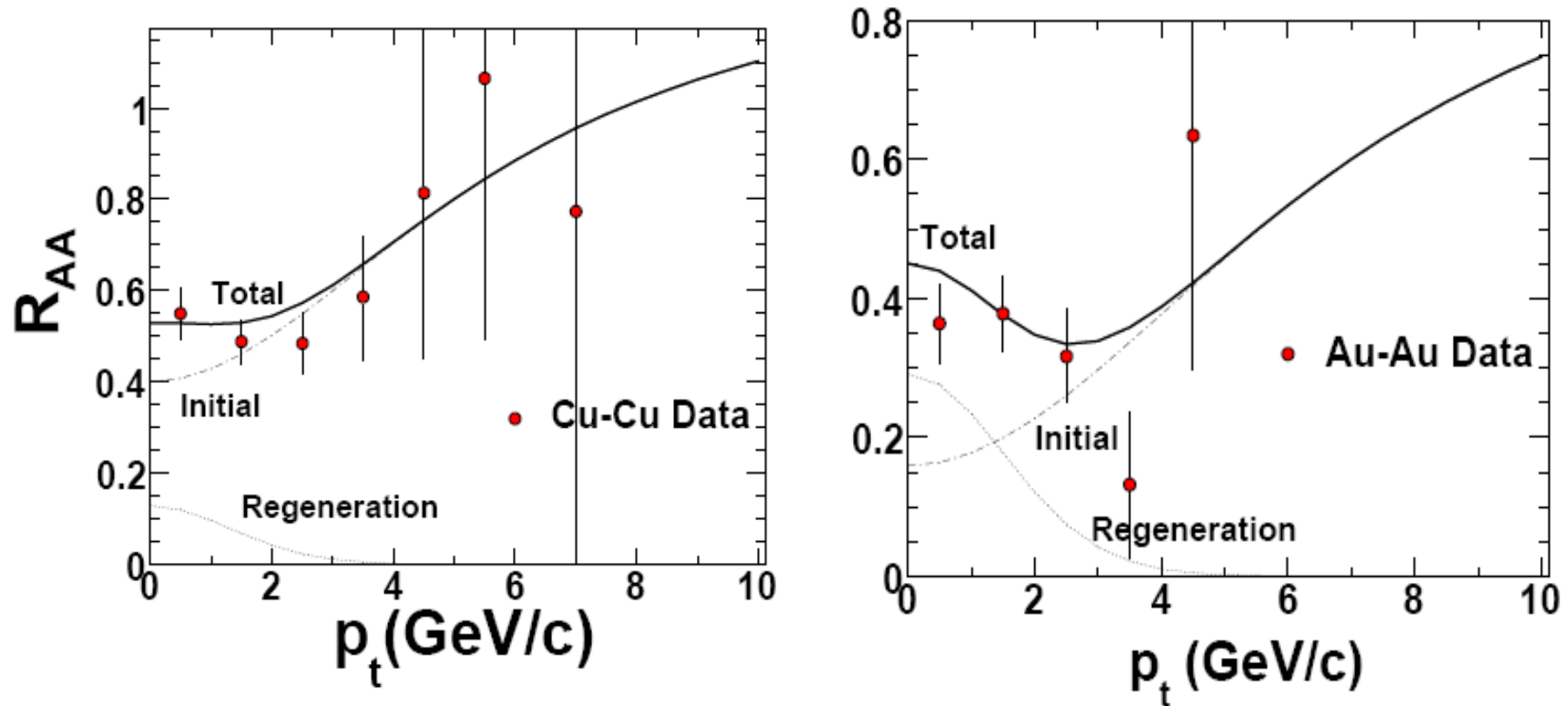


state	J/ ψ (1S)	χ_c (1P)	ψ' (2S)	Υ (1S)	χ_b (1P)	Υ (2S)	χ_b (2P)	Υ (3S)
T_d/T_c	2.10	1.16	1.12	> 4.0	1.76	1.60	1.19	1.17

***P_t Distribution and V_2
Distribution***

$J/\psi R_{AA}(p_t)$ at RHIC

Liu, Qu, Xu, Zhuang: arXiv:0901.2757, PLB2009



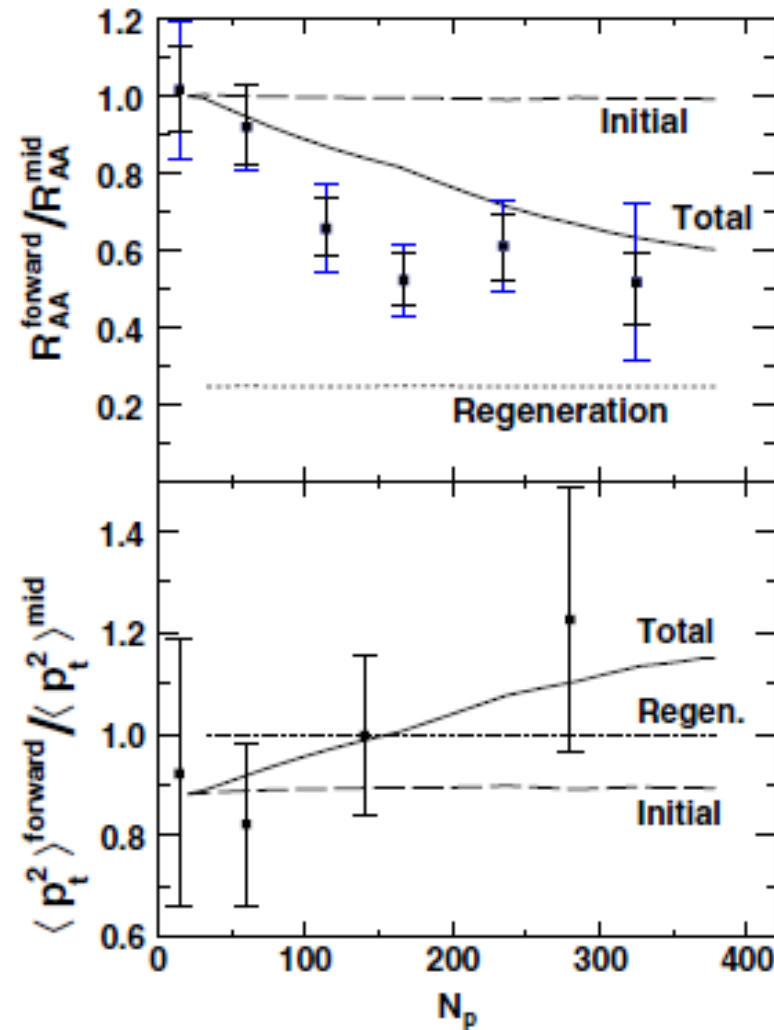
the competition between initial production and regeneration leads to a minimum, a signature for the coexistence of both production mechanisms.

$$\tau_0 = 0.6 \text{ fm}, \quad T_0 = 344 \text{ MeV},$$

$$\sigma_{pp}^{J/\psi} = 0.74 \mu\text{b}, \quad \sigma_{pp}^{c\bar{c}} = 0.12 \text{ nb} \quad (\text{PHENIX pp data) at mid rapidity}$$

J/ψ Rapidity Dependence at RHIC

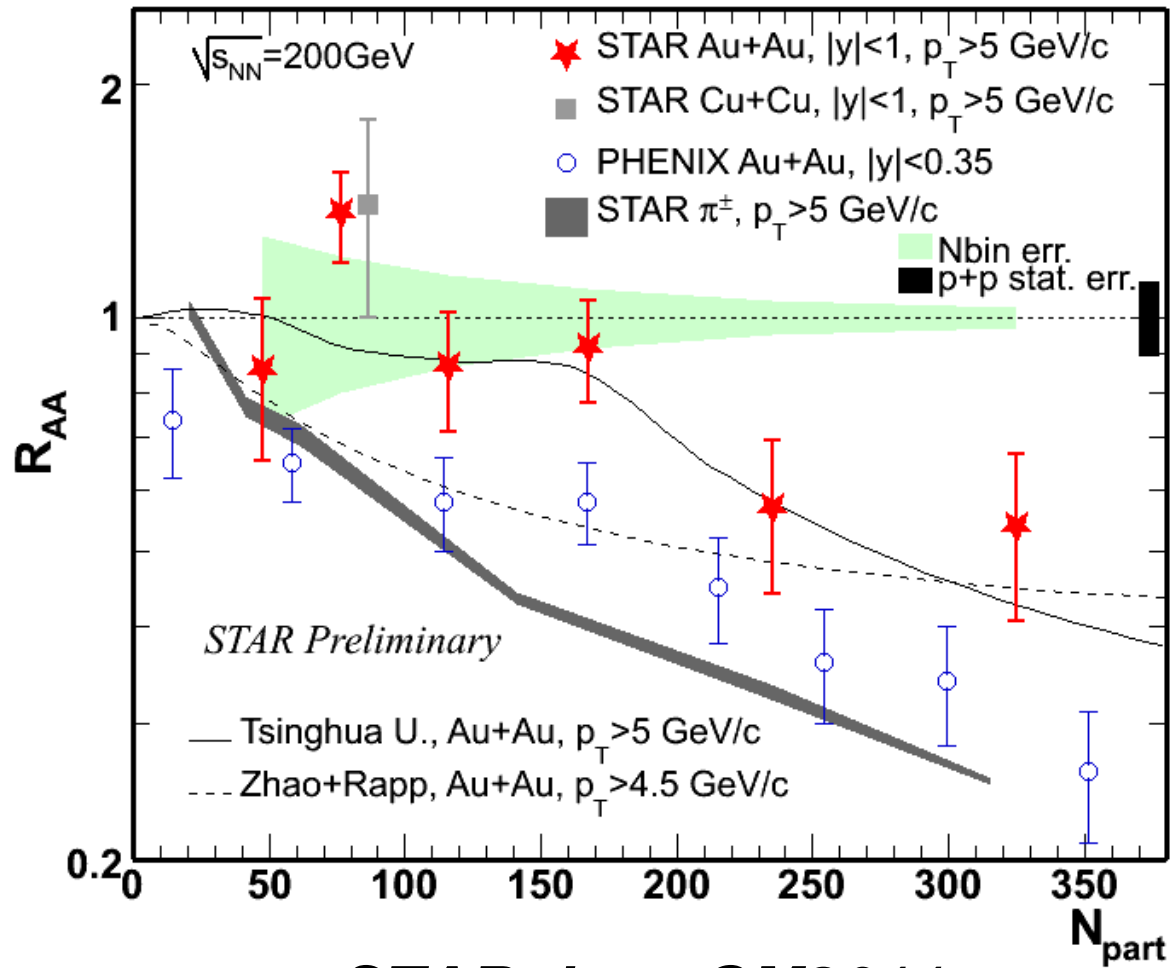
Liu, Xu, Zhuang: arXiv:0907.2723, JPG2010



less regeneration in forward rapidity explains the two puzzles naturally.

$$\sigma_{pp}^{J/\psi} = 0.42 \mu\text{b}, \quad \sigma_{pp}^{c\bar{c}} = 0.04 \text{nb} \text{ at forward rapidity}$$

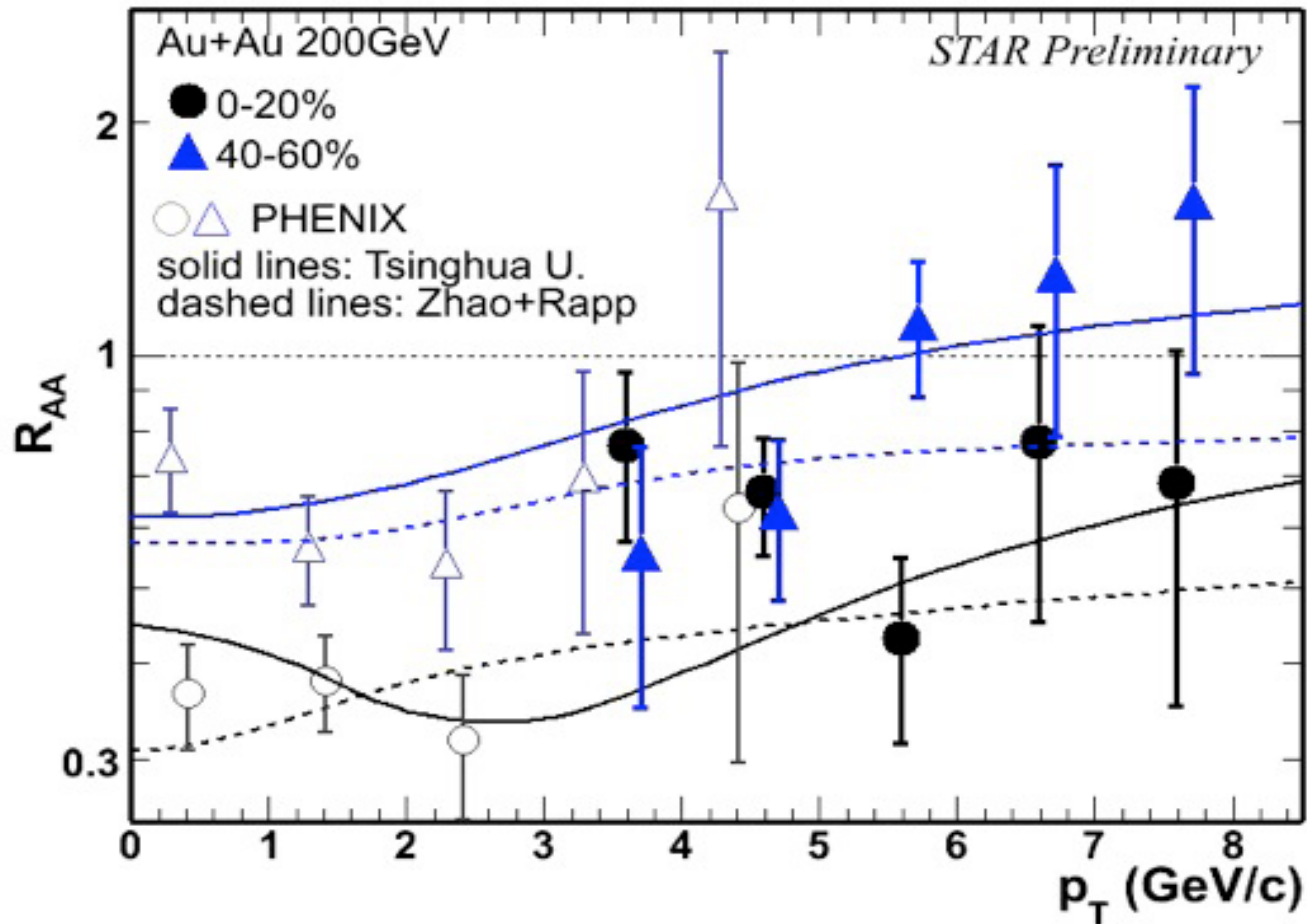
$J/\psi R_{AA}(N_p)$ at high p_T at RHIC



STAR data, QM2011

high p_T particles can survive in hot medium.

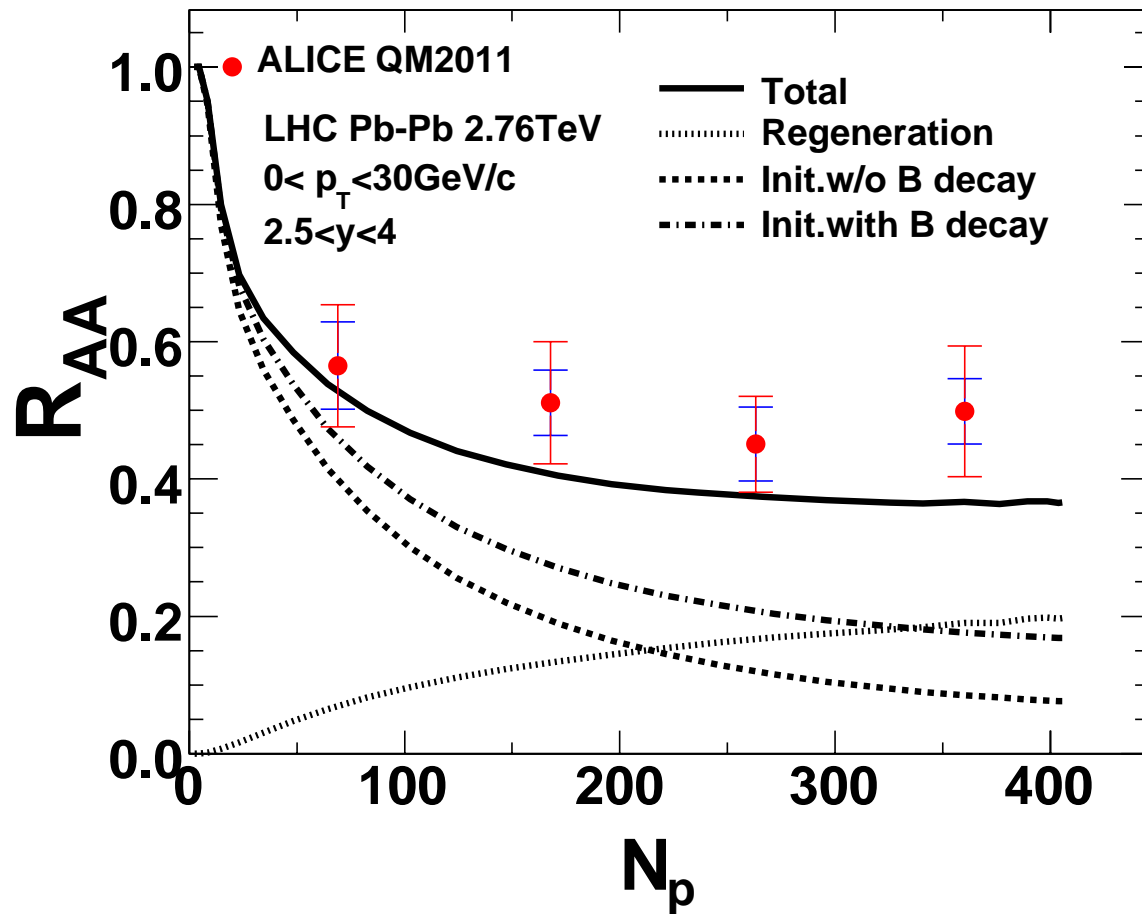
centrality dependence of $J/\psi R_{AA}(p_t)$ at RHIC



STAR data, QM2011

more suppression in central collisions

$J/\psi R_{AA}(N_p)$ at LHC

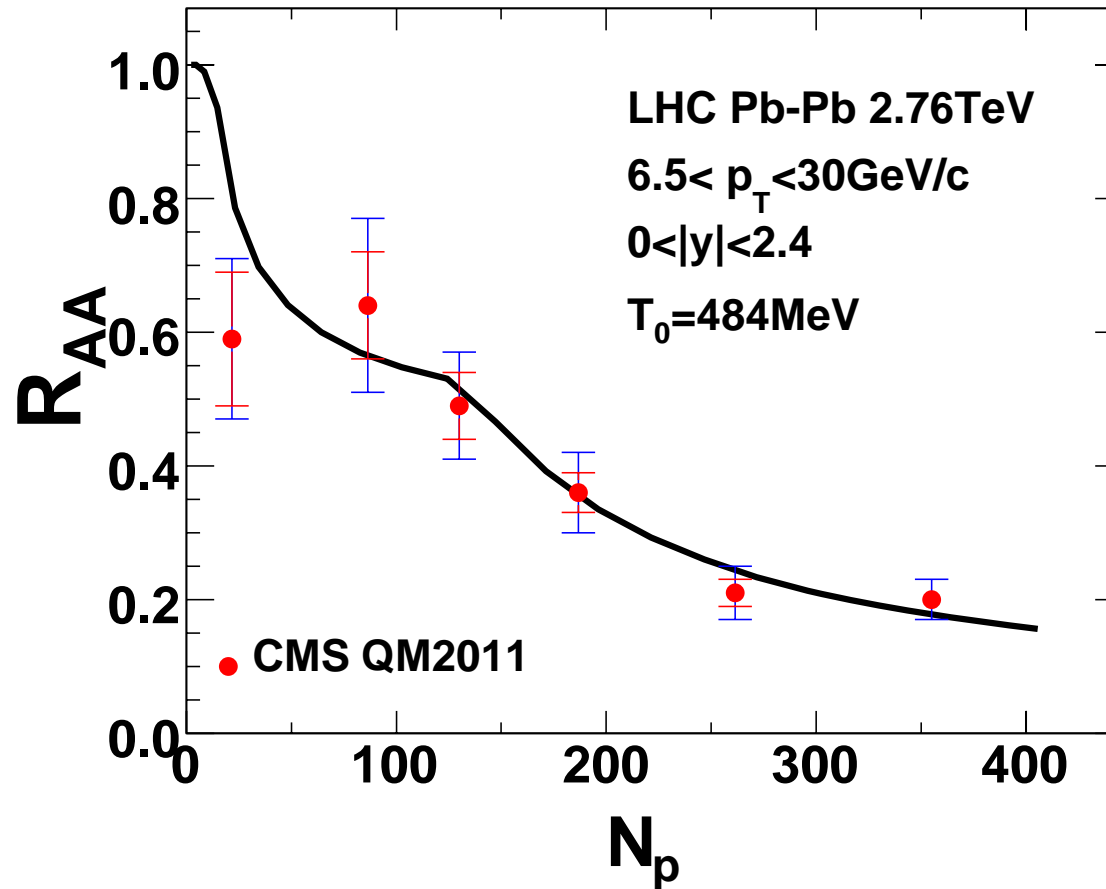


more regeneration at LHC $\rightarrow R_{AA}^{ALICE} > R_{AA}^{PHENIX}$

$$\tau_0 = 0.6 \text{ fm}, \quad T_0 = 430 \text{ MeV (Hirano, Heinz),}$$

$$\sigma_{pp}^{J/\psi} = 2.33 \mu\text{b (arXiv:1107.0137)}, \quad \sigma_{pp}^{c\bar{c}} = 3.45 \text{ nb (total, QM2011 talk by Dainese)}$$

$J/\psi R_{AA}(N_p)$ at high p_T at LHC

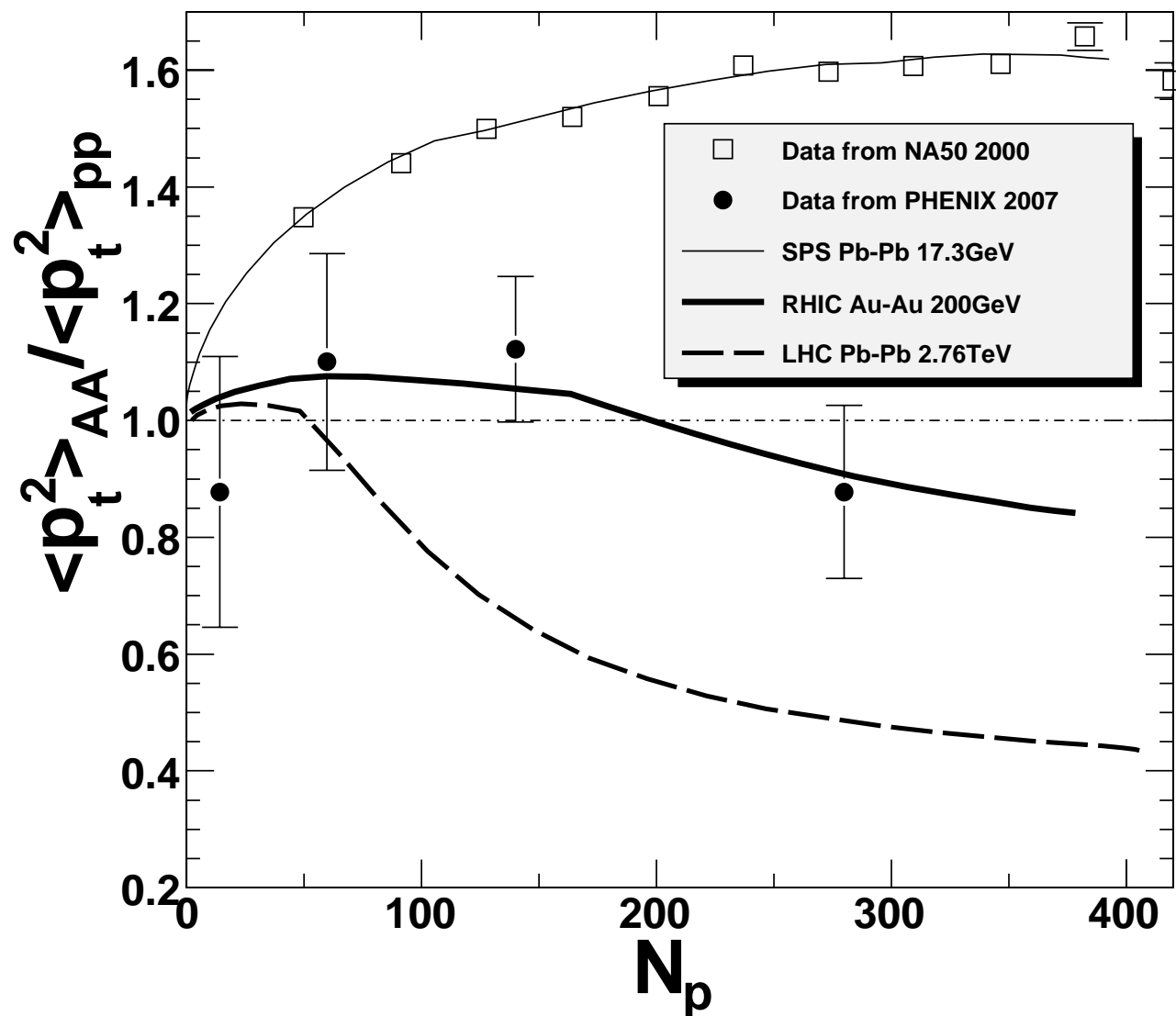


more regeneration at low p_T $\rightarrow R_{AA}^{ALICE} > R_{AA}^{CMS}$

$$\tau_0 = 0.6 \text{ fm}, \quad T_0 = 484 \text{ MeV (Hirano, Heinz),}$$

$$d\sigma_{pp}^{J/\psi} / dy = 3.5 \text{ } \mu\text{b (arXiv:1107.0137)}$$

Averaged Transverse Momentum

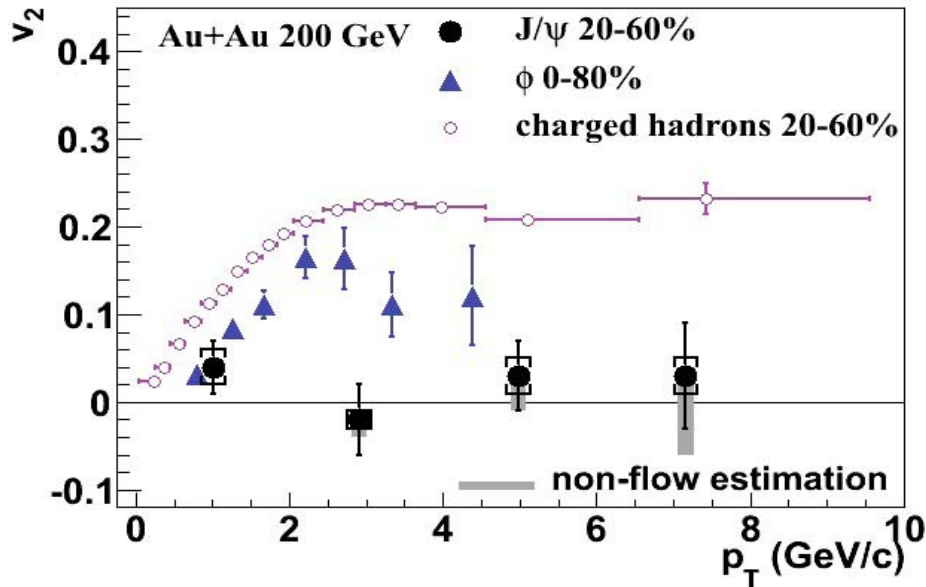


SPS:
Cronin effect

RHIC:
competition
between the two
sources

LHC:
dominant
regeneration

J/psi elliptic flow at RHIC

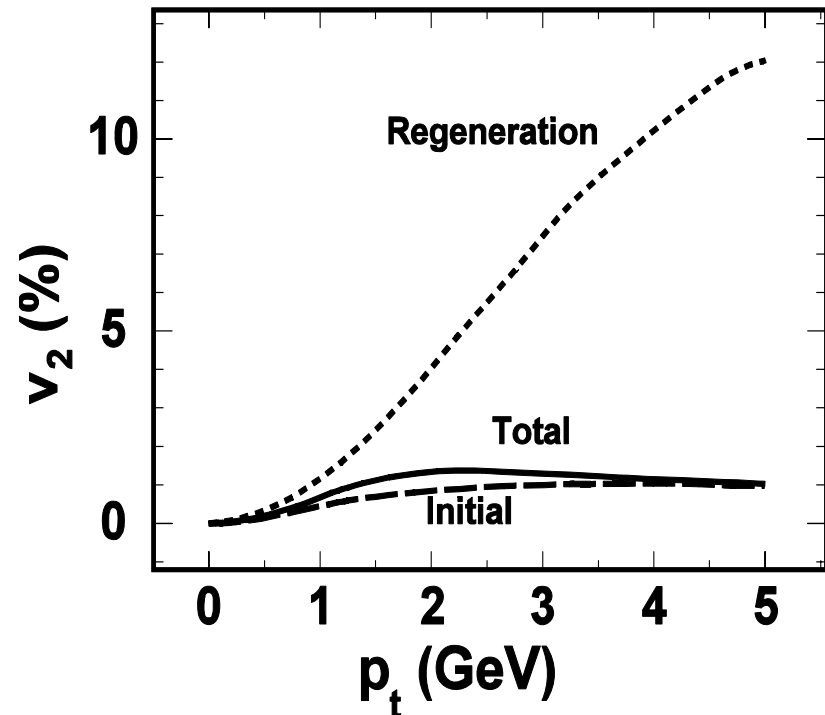


STAR data, QM2011

our prediction:

Yan, Xu, Zhuang: [nucl-th/0608010](#), PRL2006

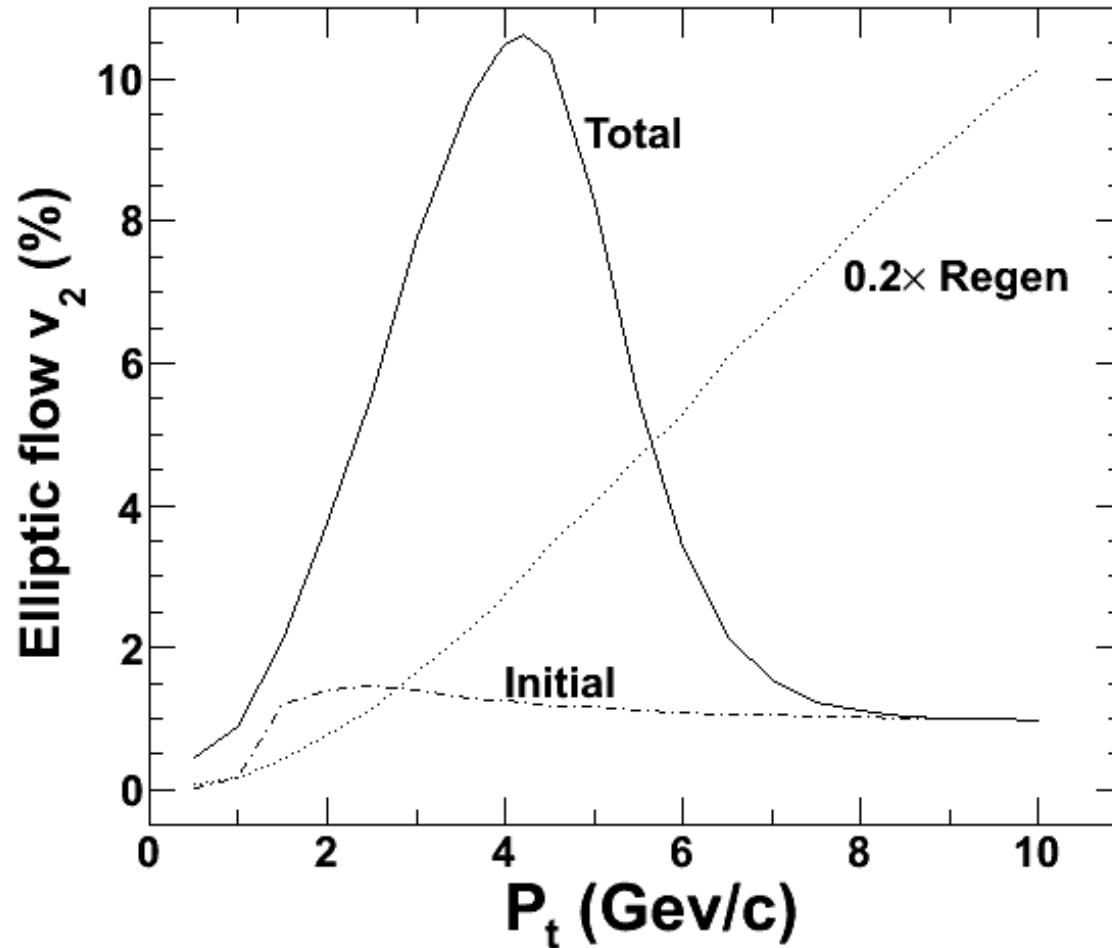
impact parameter $b=7.8$ fm



no J/psi v_2 at RHIC, but remarkable v_2 at LHC !

J/psi elliptic flow at LHC

our prediction at $\sqrt{s} = 5.5$ TeV, $b=7.8$ fm (NPA2010) :



remarkable v_2 at LHC !

γ Production and Heavy Quark Potential

Y, a Cleaner Probe at RHIC

J/ψ :

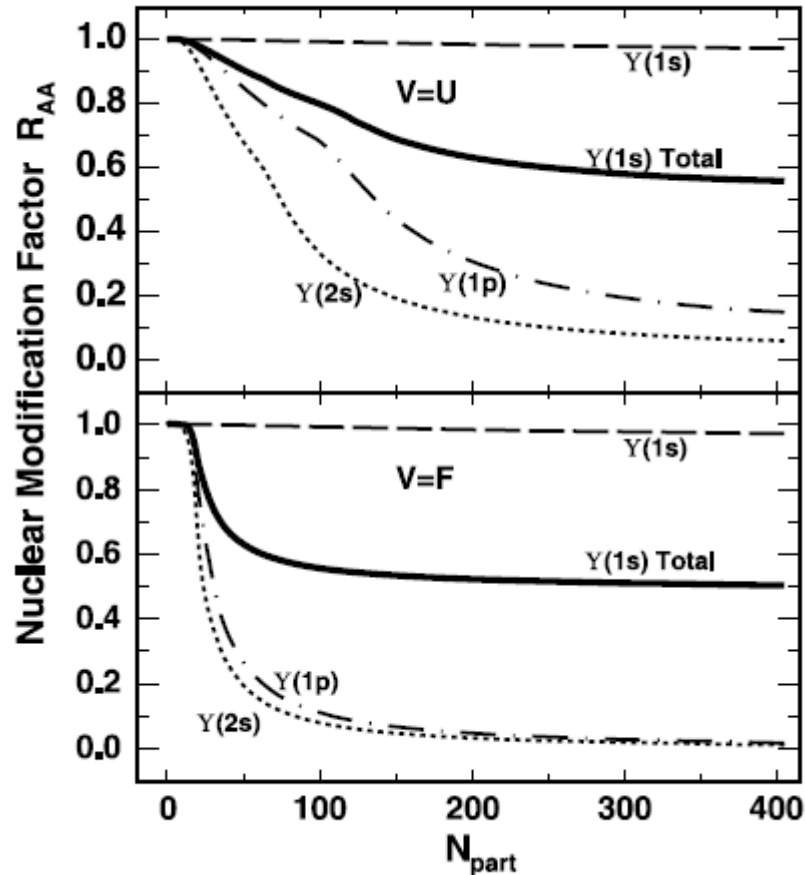
the production and suppression mechanisms are complicated: there are primordial production and nuclear absorption in the initial state and regeneration and anomalous suppression during the evolution of the hot medium.

Y:

- 1) the regeneration can be safely neglected;*
- 2) there is almost no feed-down for Y ;*
- 3) weaker CNM effect*

Y at RHIC: $R_{AA}(N_p)$

Liu, Chen, Xu, Zhuang: arXiv:1009.2585, PLB2011

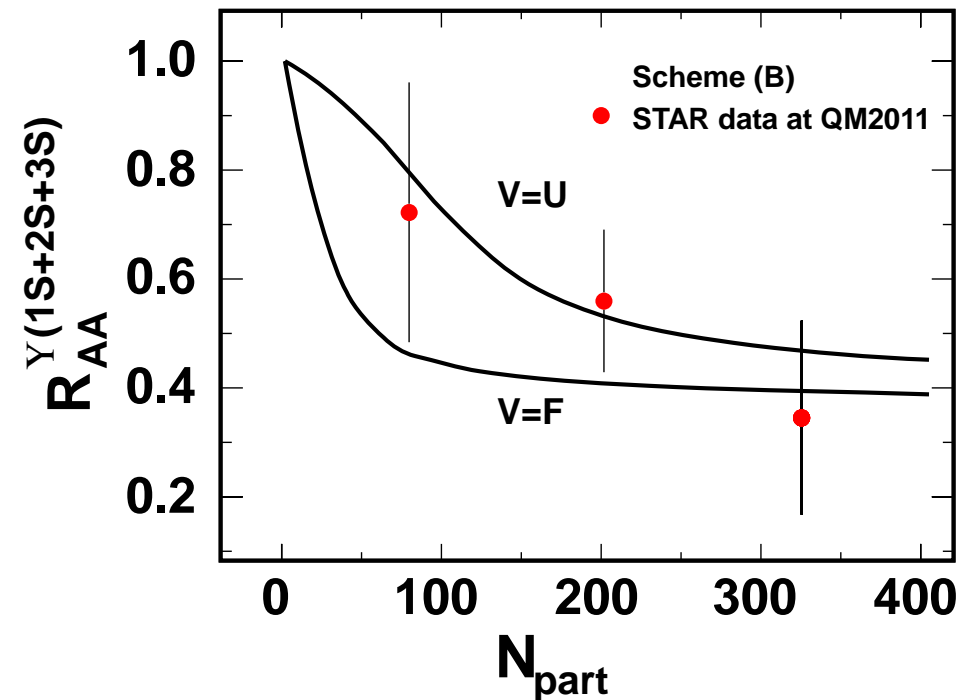


for minimum bias events:

PHENIX data: $R_{AA} < 0.64$ (NPA2009)

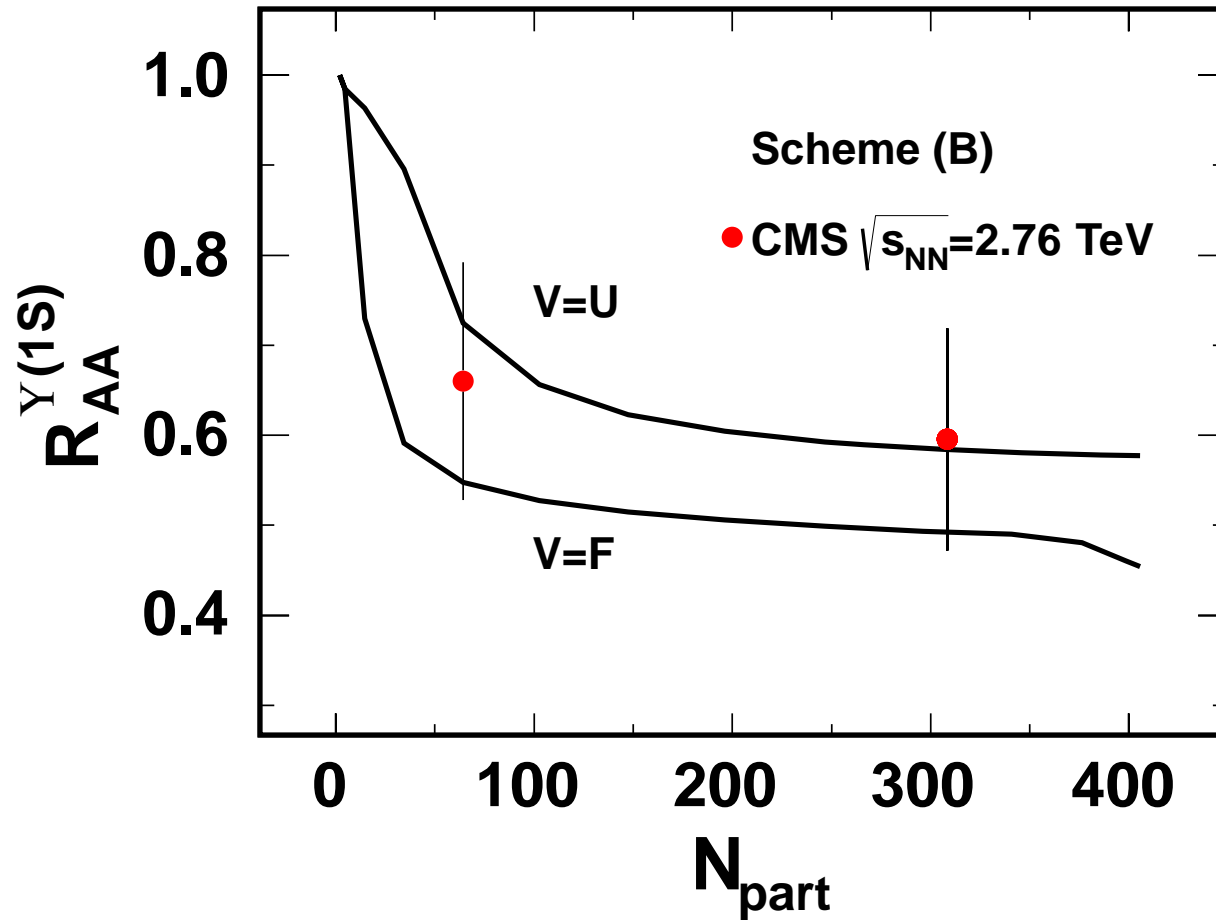
our result: $R_{AA} = 0.63$ for $V=U$

$R_{AA} = 0.53$ for $V=F$



• from the comparison with data, V is close to U .

Y at LHC: $R_{AA}(N_p)$



$$\sigma_{pp}^Y = 14 \mu b, \quad \sigma_{pp}^{b\bar{b}} = 43 nb$$

• again, V is close to U .

Conclusions:

- *pt dependence is more sensitive to the production and suppression mechanism.*
- *regeneration is important at RHIC and LHC.*
- *competition between initial production and regeneration can explain systematically the data from SPS to LHC.*
- *Upsilon production at RHIC and LHC supports $V=U$.*

Uncertainty analysis:

pp collision, shadowing effect, EoS, time scales,

Suggestions:

- *measure D - D bar correlation at LHC (Zhu, Bleicher, Huang, Schweda, Stoecker, Xu, Zhuang, PLB2007, Zhu, Xu, Zhuang, PRL2008)*
- *measure J/ψ - D correlation at LHC (since both are from the same source)*
- *measure quarkonium v_2 at LHC (which is very sensitive to the production and suppression mechanisms).*

BACKUP

Quarkonium in Vacuum

State	J/ψ (1S)	χ_c (1P)	ψ' (2S)
m (GeV/c ²)	3.10	3.53	3.68
r_0 (fm)	0.50	0.72	0.90

State	Υ (1S)	χ_b (1P)	Υ' (2S)	χ_b' (2P)	Υ'' (3S)
m (GeV/c ²)	9.46	9.99	10.02	10.26	10.36
r_0 (fm)	0.28	0.44	0.56	0.68	0.78

Contribution to the observed ground state $\Upsilon(1S)$

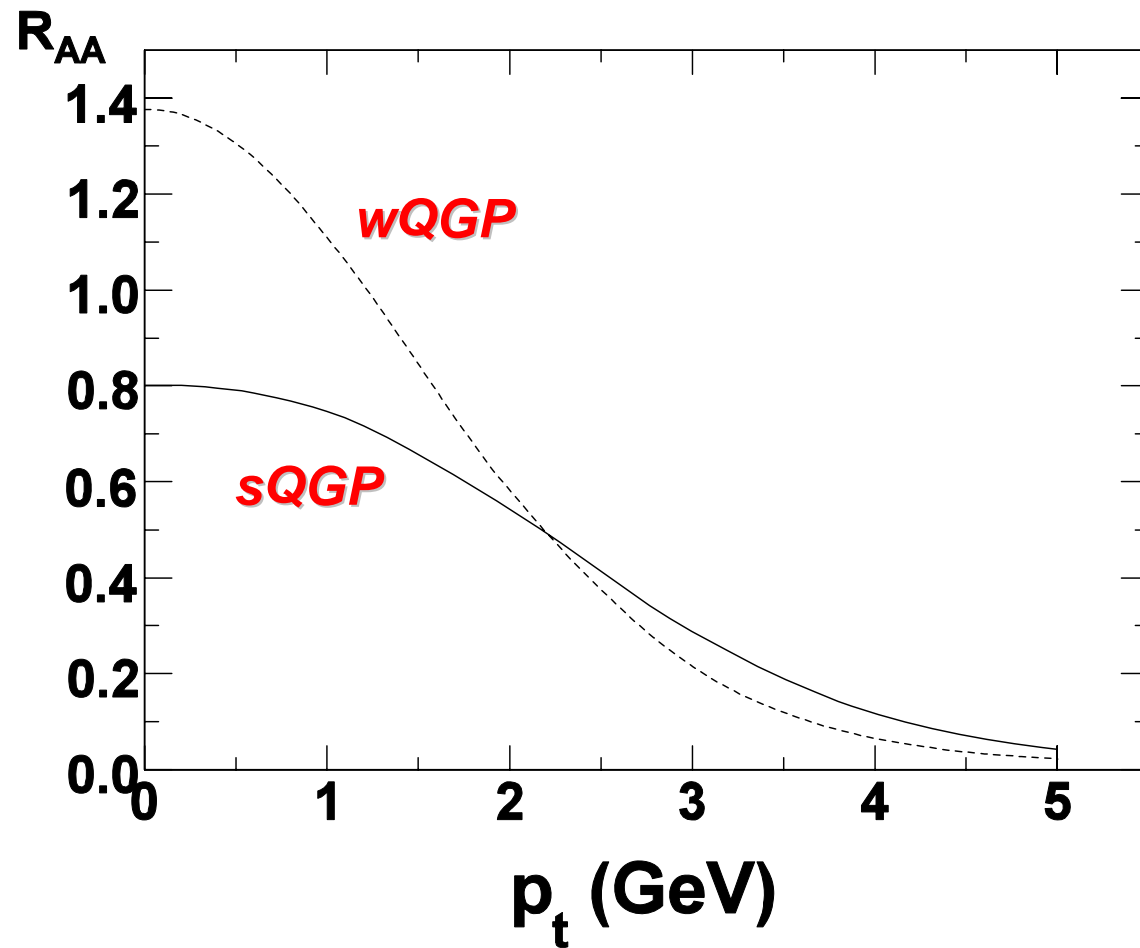
$\Upsilon(1S)$	$\Upsilon(1P)$	$\Upsilon(2S)$	$\Upsilon(2P)$	$\Upsilon(3S)$
51%	27%	11%	10%	1%

Contribution to the observed ground state J/ψ

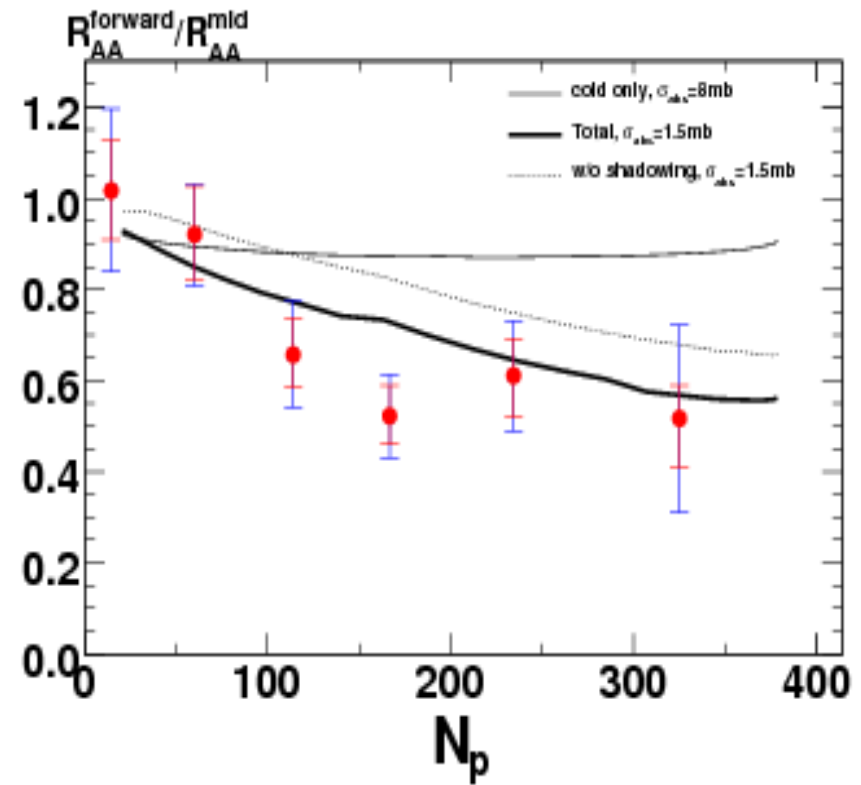
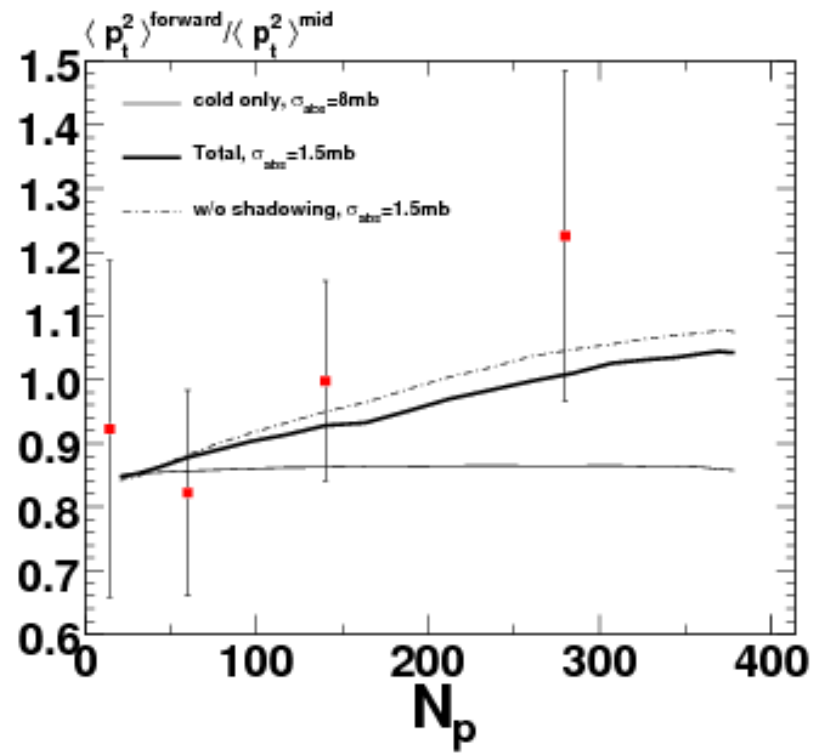
J/ψ	χ_c	ψ'
60%	30%	10%

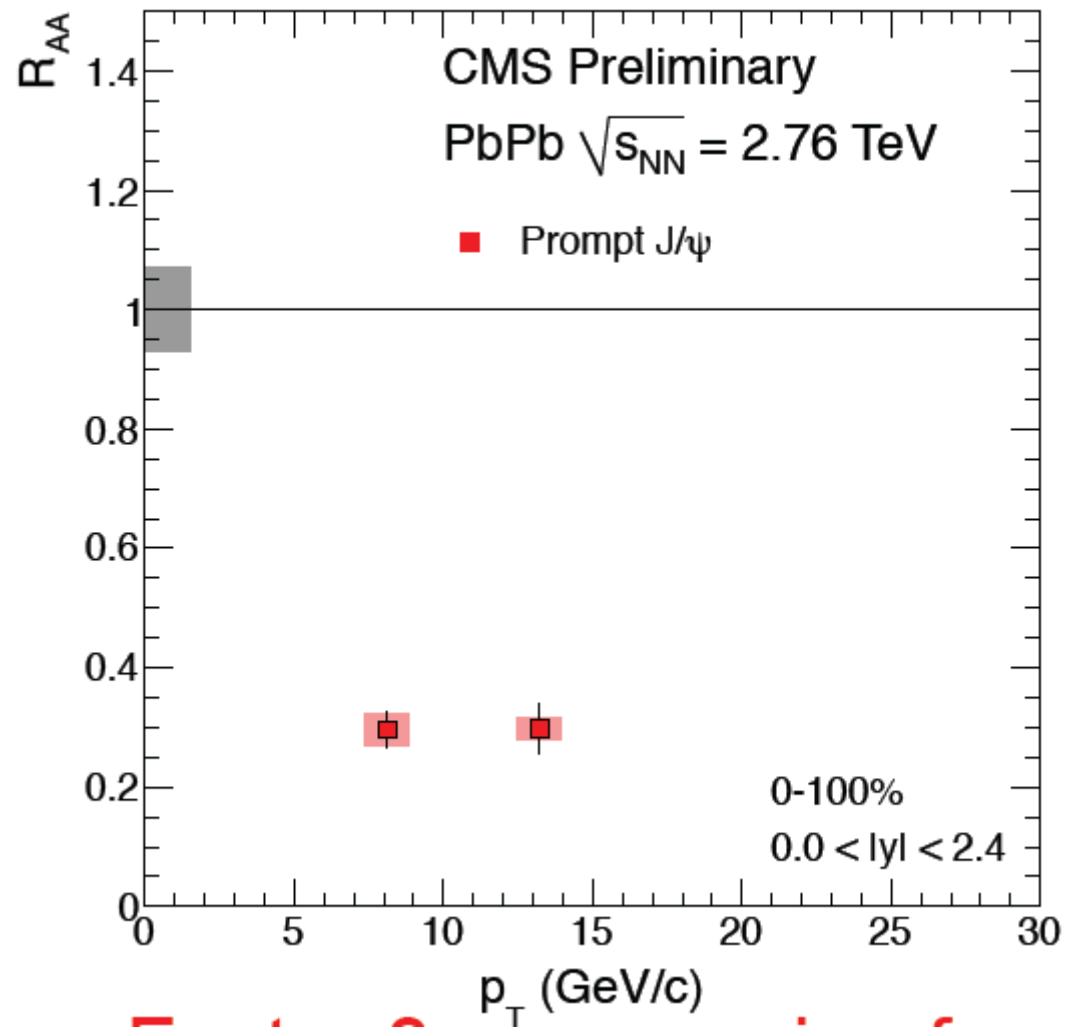
Dependence on EoS

J/Psi Pt distribution at LHC where EoS plays an essential role!



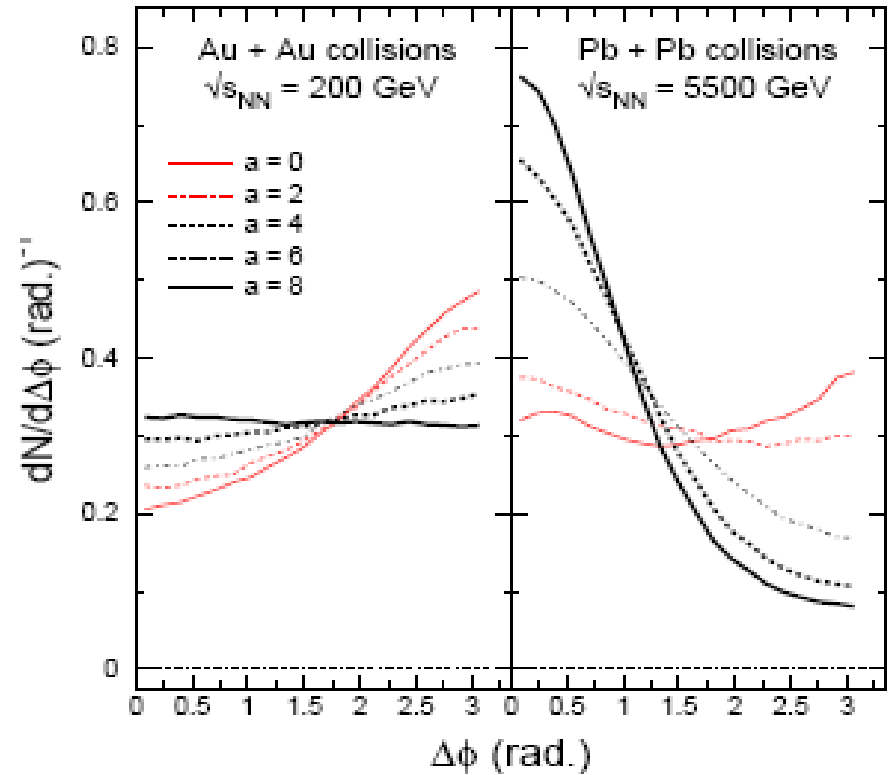
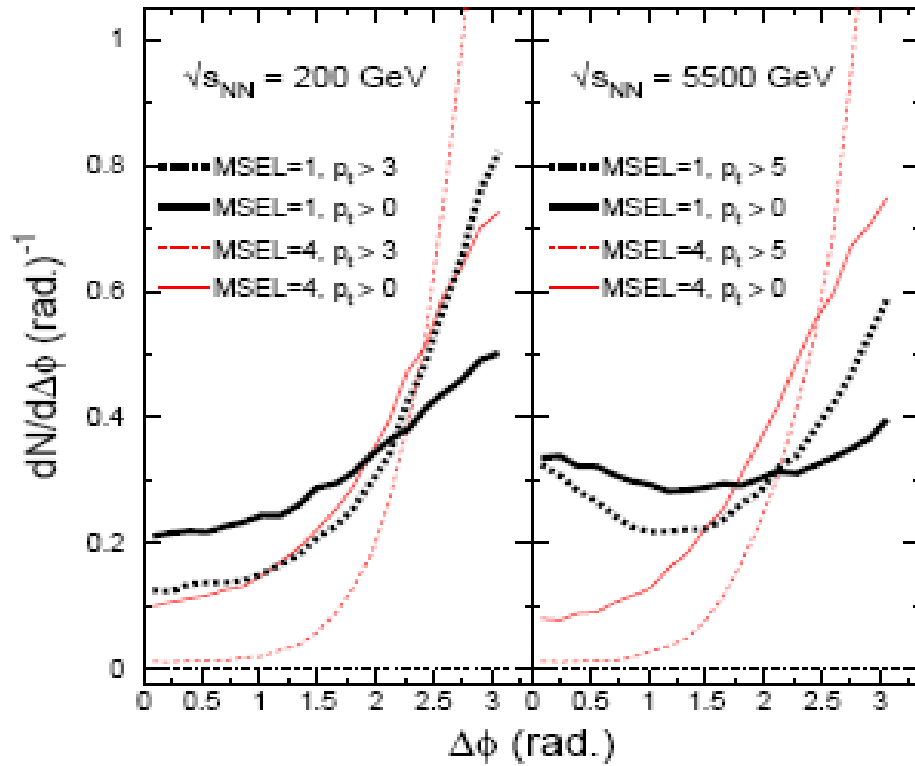
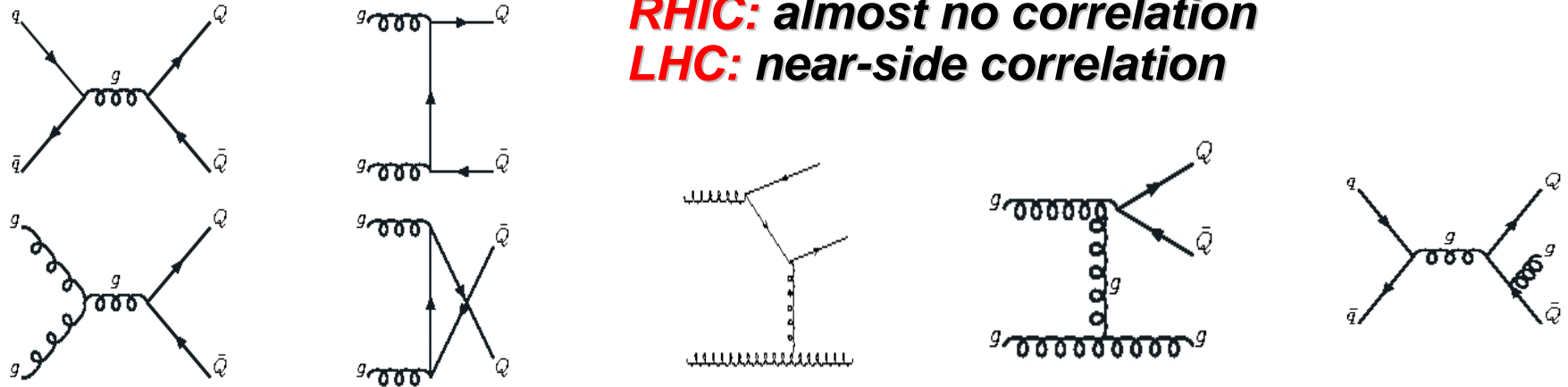
Only Cold Nuclear Matter Effect ?





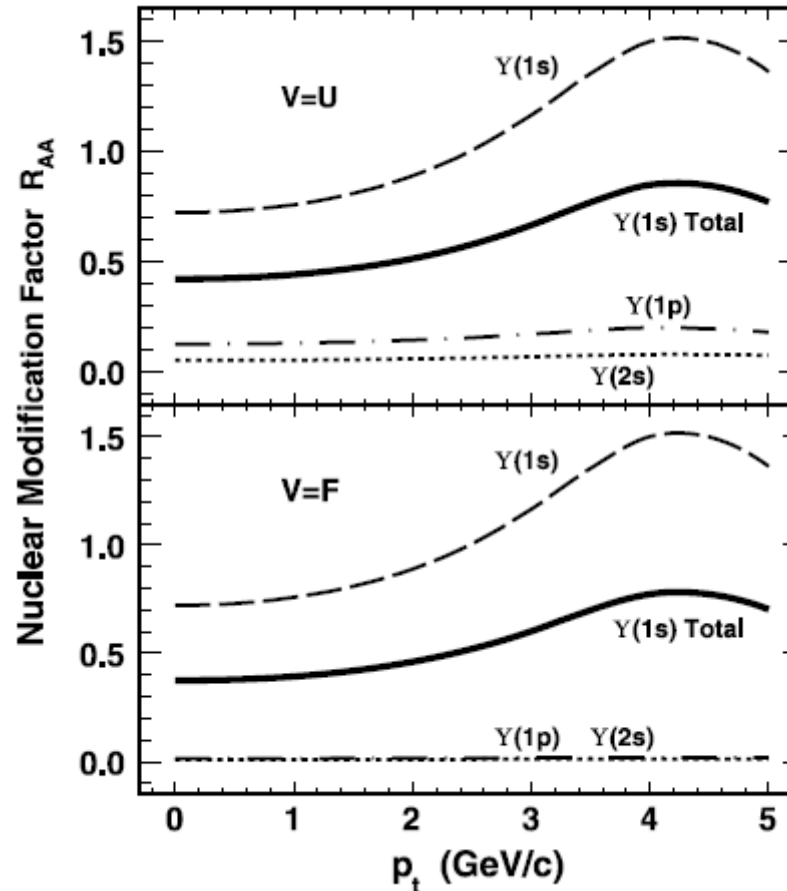
DD Correlation

NN collisions: back-to-back correlation
RHIC: almost no correlation
LHC: near-side correlation



Y at RHIC: $R_{AA}(p_t)$

Liu, Chen, Xu, Zhuang: arXiv:1009.2585,PLB2011



central Au+Au at $\sqrt{s}=200$
GeV

- **strong Cronin effect for survived ground state**
- **no p_t dependence for disappeared excited states**

Y at RHIC: $\langle p_t^2 \rangle(N_p)$

relation between Υ at RHIC and J/ψ at SPS:
no Υ regeneration at RHIC and no J/ψ regeneration at SPS

- $T_D^{\Upsilon(1s)} = 4T_c > T_{RHIC}$ **no $\Upsilon(1s)$ suppression at RHIC**
- $T_D^{J/\psi} = 2T_c > T_{SPS}$ **no J/ψ suppression at SPS**

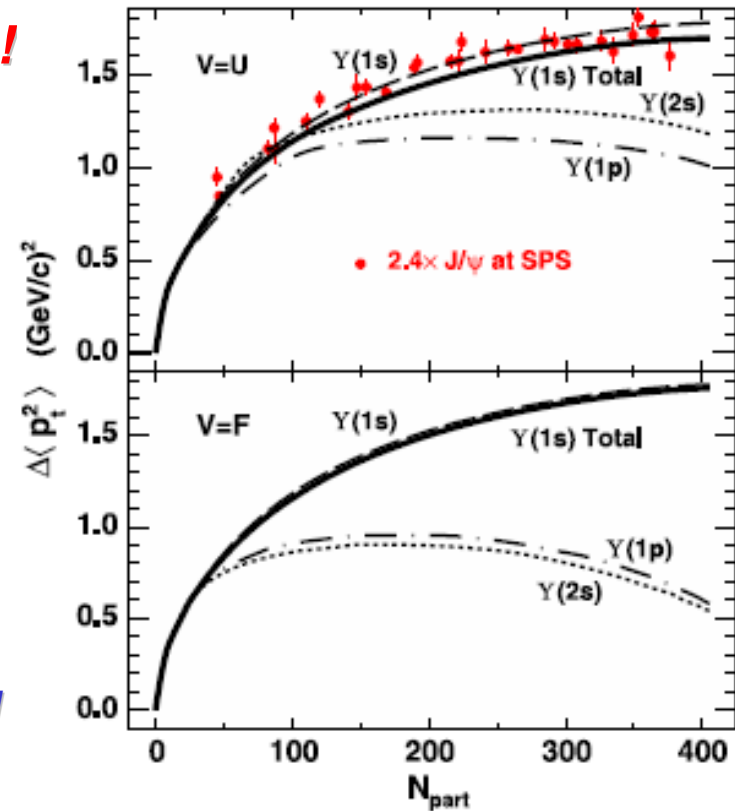
both are controlled by the Cronin effect !

$$\Delta \langle p_t^2 \rangle = \langle p_t^2 \rangle_{AA} - \langle p_t^2 \rangle_{pp} = a_{gN} L$$

$$\Delta \langle p_t^2 \rangle_{\Upsilon}^{RHIC} = \frac{a_{gN}^{RHIC} R_{Au}}{a_{gN}^{SPS} R_{Pb}} \Delta \langle p_t^2 \rangle_{J/\psi}^{SPS} = 2.4 \Delta \langle p_t^2 \rangle_{J/\psi}^{SPS}$$

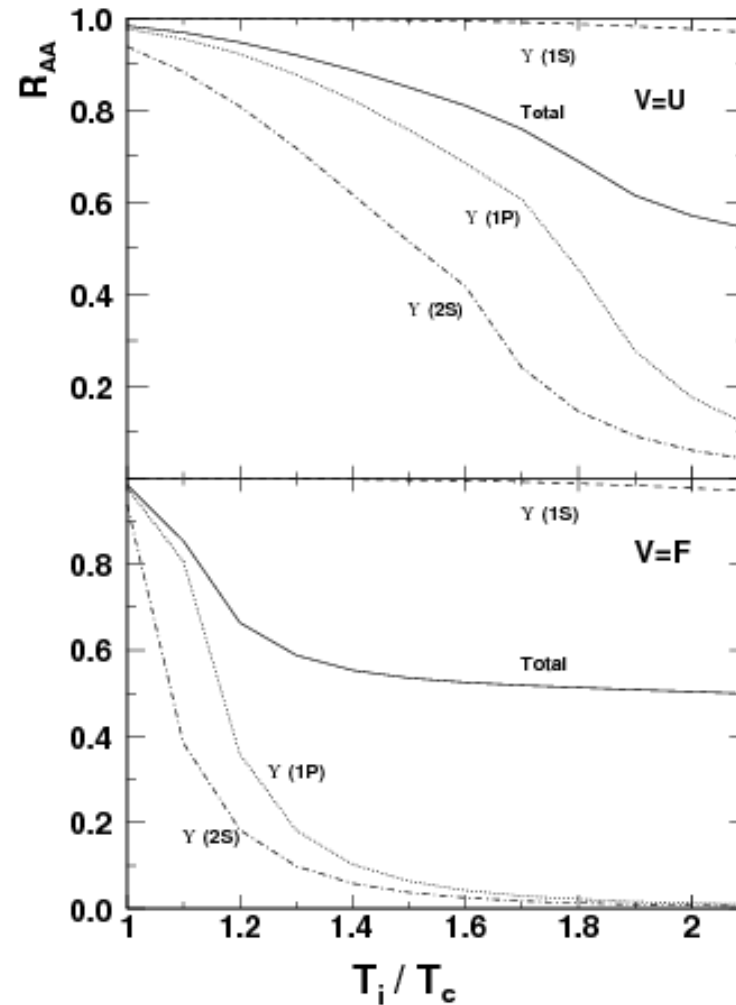
Au+Au at $\sqrt{s}=200$ GeV

Liu, Chen, Xu, Zhuang: arXiv:1009.2585, PLB2011



Measuring RHIC Temperature by Excited γ States

initial temperature dependence of R_{AA}

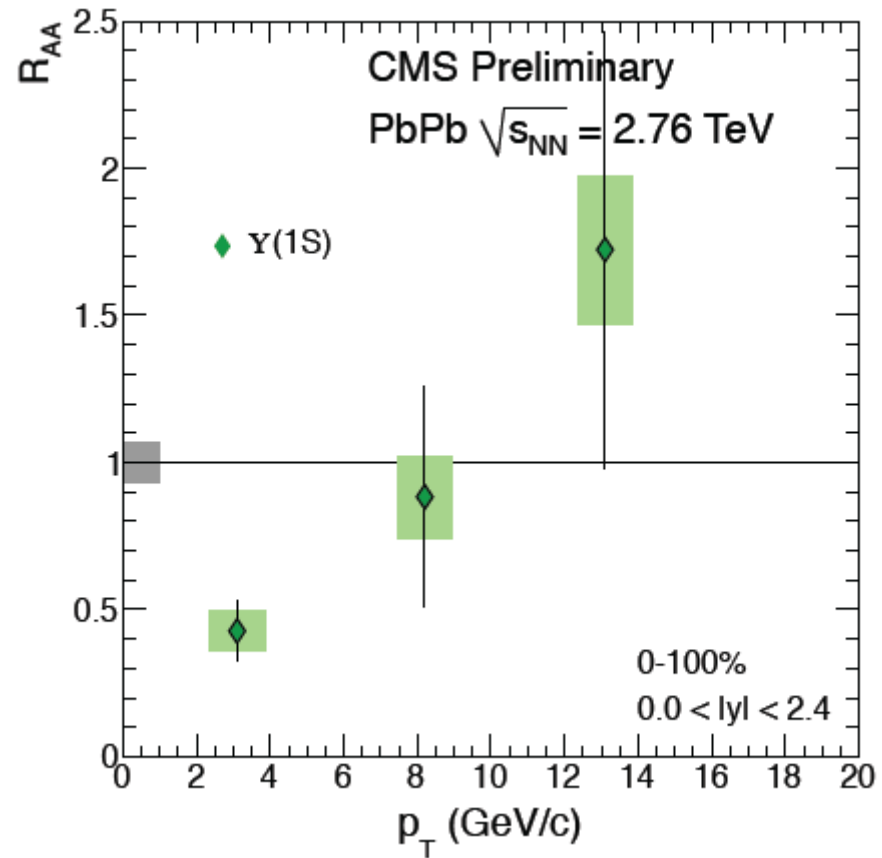


*central Au+Au at $\sqrt{s}=200$
GeV*

Liu, Chen, Xu, Zhuang: arXiv:1009.2585, PLB2011

suppression of excited γ states is sensitive to the fireball temperature !

Y at LHC: $R_{AA}(p_t)$



high p_t is controlled by initial production !