

# Non-relativistic leptogenesis

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in collaboration with Dietrich Bödeker

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- Theory to explain the baryon asymmetry of the universe
- Proposed by Yanagida and Fukugita in 1986
- Standard Model extended by heavy right-handed Majorana-neutrinos
- Their decays  $N \rightleftharpoons \varphi + \ell$  and  $N \rightleftharpoons \bar{\varphi} + \bar{\ell}$  can fulfill the Sakharov conditions
  - $L$  violation, cf.  $B - L$  violation
  - $CP$  violation
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  - $L$  violation, cf.  $B - L$  violation
  - $CP$  violation
  - Departure from thermal equilibrium
- Sphaleron processes convert  $B - L$  into  $B$  - asymmetry

- 1 Leptogenesis in the non-relativistic limit
- 2 Relativistic corrections
- 3 Radiative corrections
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- Measure for the washout strength:  $K = \Gamma_0/H(T = M_N)$ 
  - $K \gg 1$ : “Strong washout”
  - $K \ll 1$ : “Weak washout”
- In the strong washout regime, asymmetry that was created at  $T > M_N$  does not play a role
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- $\Delta m_{atm}^2$  and  $\Delta m_{sol}^2$  imply:  $7 \lesssim K \lesssim 46$

**$\Rightarrow$  Final asymmetrie was created at  $T < M_N$ ,  
i.e. in the non-relativistic regime**

$$\left(\frac{d}{dt} + 3H\right)n_N = -\Gamma_N(n_N - n_N^{\text{eq}})$$

$$\left(\frac{d}{dt} + 3H\right)n_{B-L} = \Gamma_{B-L,N}(n_N - n_N^{\text{eq}}) - \Gamma_{B-L}n_{B-L}$$

*These equations are valid to all orders in the SM couplings!*



# Determining $\Gamma_N$ at leading order

- Boltzmann equation

$$(\partial_t - H p \partial_p) f_N = \frac{M_N \Gamma_0}{E_N} (e^{-E_N/T} - f_N)$$

- Yukawa interaction

$$\mathcal{L}_{NYuk} = h_{ij} \overline{N_{Ri}} \tilde{\varphi}^+ \ell_{Lj} + \text{h.c.} \Rightarrow \Gamma_0 = \frac{|h_{11}|^2 M_N}{8\pi}$$

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# Rate equations in the non-relativistic regime

$$\left(\frac{d}{dt} + 3H\right)n_N = -\Gamma_N (n_N - n_N^{\text{eq}})$$

$$\left(\frac{d}{dt} + 3H\right)n_{B-L} = \Gamma_{B-L,N} (n_N - n_N^{\text{eq}}) - \Gamma_{B-L} n_{B-L}$$

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LO coefficients:

$$\Gamma_N = \Gamma_0$$

$$\Gamma_{B-L,N} = \epsilon \Gamma_0$$

$$\Gamma_{B-L} = \frac{3}{\pi^2} \left(c_\ell + \frac{c_\varphi}{2}\right) z^2 K_1(z) \Gamma_0$$

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- Boltzmann equation

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- integrate Boltzmann equation over  $\vec{p}$

- therefor expand  $1/E_N$   $\frac{1}{E_N} \approx \frac{1}{M_N} + \frac{\vec{p}^2}{2M_N^3}$

$$\left(\frac{d}{dt} + 5H\right) u = -\Gamma_u (u - u^{\text{eq}})$$

$$u \equiv \frac{1}{M_N} \cdot 2 \int \frac{d^3 p}{(2\pi)^3} \frac{\vec{p}^2}{2M_N} f_N$$

$$\left(\frac{d}{dt} + 3H\right) n_N = -\Gamma_N (n_N - n_N^{\text{eq}}) + \Gamma_{N,u} (u - u^{\text{eq}})$$

$$\left(\frac{d}{dt} + 3H\right) n_{B-L} = \Gamma_{B-L,N} (n_N - n_N^{\text{eq}}) + \Gamma_{B-L,u} (u - u^{\text{eq}}) - \Gamma_{B-L} n_{B-L}$$

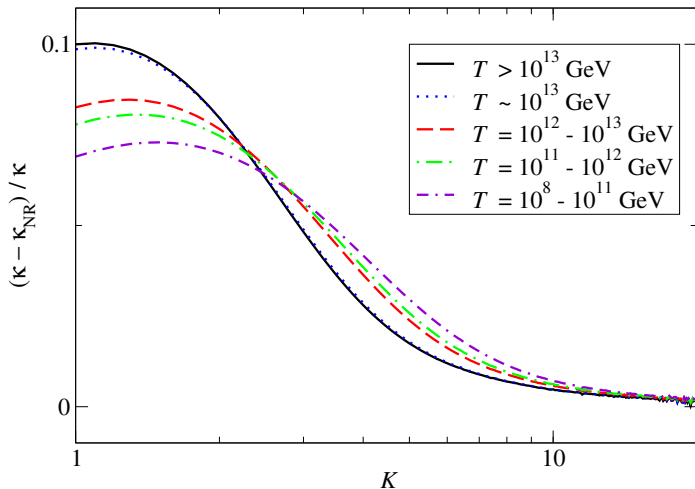
$$\Gamma_u = \Gamma_0$$

$$\Gamma_{N,u} = \Gamma_0$$

$$\Gamma_{B-L,u} = \epsilon \Gamma_0$$



# Relativistic corrections - Numerical results



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# Radiative corrections

- Before: LO
  - Decays ( $1 \rightarrow 2$ )
- Now: NLO
  - $2 \rightarrow 2$
  - $1 \rightarrow 3$
  - $1 \rightarrow 2$  virtual corrections

$$\left. \frac{\partial f_N}{\partial t} \right|_{n_N=0} = f_F(E_N) \Gamma_0 \frac{M_N}{E_N} \left\{ a + \frac{\vec{p}^2}{M_N^2} b + O\left(g^4, \frac{g^3 T^2}{M_N^2}, \frac{g^2 T^6}{M_N^6}\right) \right\}$$

- $a = 1 - \frac{\lambda T^2}{M_N^2} - |h_t|^2 \left[ \frac{21}{2(4\pi)^2} + \frac{7\pi^2}{60} \frac{T^4}{M_N^4} \right] + (g_1^2 + 3g_2^2) \left[ \frac{29}{8(4\pi)^2} - \frac{\pi^2}{80} \frac{T^4}{M_N^4} \right]$
- $b = - \left[ |h_t|^2 \frac{7\pi^2}{45} + (g_1^2 + 3g_2^2) \frac{\pi^2}{60} \right] \frac{T^4}{M_N^4}$

A. Salvio, P. Lodone and A. Strumia, *Towards leptogenesis at NLO: the right-handed neutrino interaction rate*, JHEP 1108 (2011) 116 [arXiv:1106.2814 [hep-ph]]

M. Laine and Y. Schröder, *Thermal right-handed neutrino production rate in the non-relativistic regime*, JHEP 1202 (2012) 068 [arXiv:1112.1205 [hep-ph]]

$$\left(\frac{d}{dt} + 5H\right) u = -\Gamma_u (u - u^{\text{eq}})$$

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NLO coefficients:

$$\Gamma_u = a \Gamma_0$$

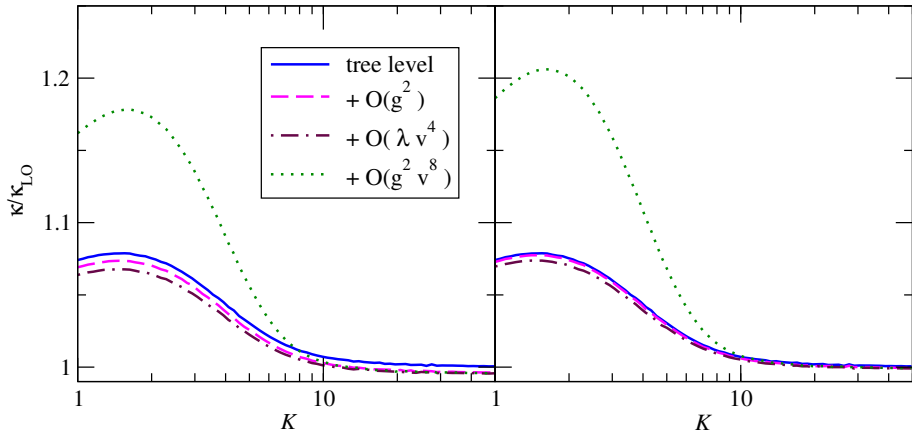
$$\Gamma_N = a \Gamma_0$$

$$\Gamma_{N,u} = (a - 2b) \Gamma_0$$

# Radiative corrections - Numerical results

$$M_N = 10^{10} \text{ GeV}$$

$$M_N = 10^8 \text{ GeV}$$



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## The non-relativistic expansion is a convenient tool for computing the lepton asymmetry!

- Relativistic corrections are small
- Accuracy can be easily controlled
- Rate equations are simple
- Radiative corrections can be included
- Only works in the (favoured) strong washout regime

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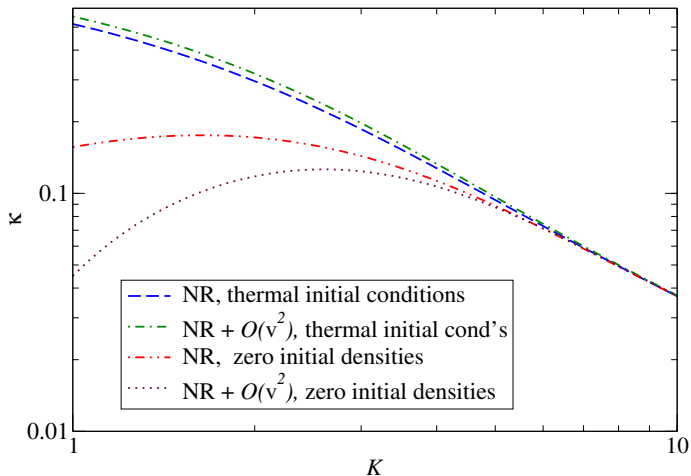
- Relativistic corrections are small
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- Only works in the (favoured) strong washout regime
  
- More work to do
  - Include radiative corrections in the  $B - L$  rate equation
  - Include flavor effects



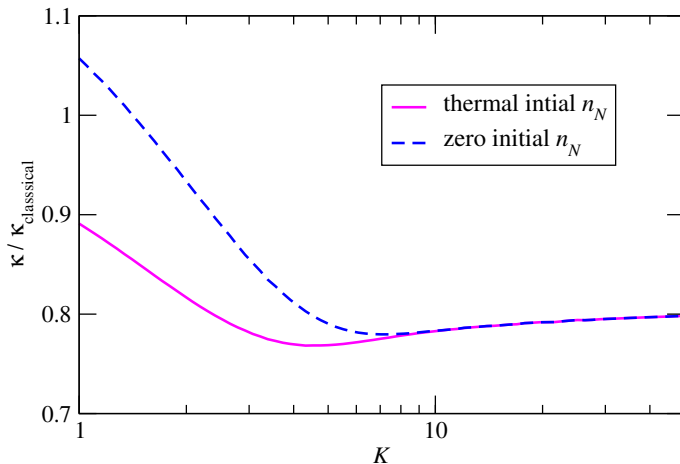
Thank you for your attention!

$$\begin{aligned} K &= \frac{\Gamma_0}{H(T = M_N)} = \frac{\frac{|h_{11}|^2 M_N}{8\pi}}{\sqrt{\frac{8\pi^3 g_*}{90} \frac{M_N^2}{M_{Pl}}}} \\ &= \frac{M_{Pl}}{8\pi \cdot 1.66 \sqrt{g_*} \cdot v^2} \cdot \frac{|h_{11}|^2 v^2}{M_N} \\ &= \text{const.} \cdot \tilde{m}_1 \end{aligned}$$

# Dependence on initial conditions



# Dependence on statistics



# Ratio of reaction rates and the Hubble rate

