Random Matrix Theory

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Lecture 2

- spectral statistics from $\varepsilon \chi PT$ and from RMT
- inclusion of chemical potential why χ PTis sensitive to baryonic μ
- further extensions: Wilson χPT

• Banks Casher Relation $\rho_D(\lambda=0) = \frac{1}{\pi} V\Sigma$ macroscopic density

$$\Sigma(\boldsymbol{m}) \equiv \frac{1}{V} \partial_{\boldsymbol{m}} \log \mathcal{Z}_{\mathbf{QCD}} = \left\langle \sum_{k} \frac{2m}{V(\lambda_{k}^{2} + m^{2})} \right\rangle = \int d\lambda \rho_{D}(\lambda) \frac{2m}{(\lambda^{2} + m^{2})}$$

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- microscopic scaling: $\lim_{V\to\infty} \frac{1}{\Sigma V} \rho_D(\lambda \Sigma V)$ zoom into origin

• Resolvent definition

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 - complex eigenvalues $\lambda \ (\mu \neq 0)$:

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• alternative: Replicas $\lim_{n\to 0} \frac{1}{n} \partial_m \det[D+m]^n$

How to compute densities from $\varepsilon \chi \text{PT-} \mu = 0$

• auxiliary fermion- boson pair
$$G(m) = \partial_m \left\langle \frac{\det[D+m]}{\det[D+m_B]} \Big|_{m=m_B} \right\rangle$$

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- **generating functions** for all *k*-point densities:

 $U(N_F + k|k)$ -supergroup integral of $\varepsilon \chi PT$, take derivatives & discontinuities

vs RMT
$$\rho_k(y_1, \ldots, y_k) \sim \int dy_{k+1} \ldots dy_N \mathcal{P}_{jpdf}$$
 (or + source too)

generating functions agree, so also all individual eigenvalues do! [F. Basile, G.A. 07]

proof: Superbosonisation Theorem also [Littelmann,Sommers,Zirnbauer]

 $\int d\Psi f(\Sigma_k \Psi_k \times \Psi_k) \sim \int_{U(N_F + k|k)} dU_0 \operatorname{sdet}[U_0]^N f(U_0)$

$$\mathcal{Z}_{\varepsilon\chi PT} = \int dU_0 \det[U_0]^{\nu} \exp\left[-\frac{1}{4}\mu^2 F^2 V \mathsf{Tr}[U_0, B][U_0^{\dagger}, B] - \frac{1}{2}\Sigma V \mathsf{Tr}M(U_0 + U_0^{\dagger})\right]$$

• $/D + \gamma_0 \mu$ (non-Hermitian): add μ as **imaginary vector current** $\partial_{\alpha} U \rightarrow \partial_{\alpha} U + \mu [B_{\alpha}, U]$, B_0 charge matrix = diag(1,-1) for isospin

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- **example** $N_f = 2$ & isospin $B = \sigma_3$ (*F* back!) $\overline{Z}_{\nu} = \int_0^1 dtt \exp[-2t^2 \mu^2 F^2 V] I_{\nu} (V \Sigma m_1 t) I_{\nu} (V \Sigma m_2 t)$ check $\mu \to 0$: $[I_{\nu}(\hat{m}_1) \hat{m}_2 I_{\nu+1}(\hat{m}_2) - I_{\nu}(\hat{m}_2) \hat{m}_1 I_{\nu+1}(\hat{m}_1)]/(m_1^2 - m_2^2)$

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- $\mu \rightarrow i \mu$ trivial for \mathcal{Z} , but NOT for spectrum of/ $D + \gamma_0 \mu$!

- density resolvent relation for $\mu \neq 0$: $\rho(\lambda) \sim \partial_{\lambda}^* G(\lambda)$
- other rep. of 2-dim delta:

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• we need an **auxiliary pair of complex conjugate** fermions and bosons to generate the resolvent (and regularize)

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- example quenched density: identify $N\mu^2 = VF^2\mu_{Lat}^2 \equiv \alpha^2$

$$\rho(\xi) = \frac{1}{\alpha^2} |\xi|^2 K_{\nu} \left(\frac{|\xi|^2}{2\alpha^2} \right) \mathbf{e}^{+\frac{1}{4\alpha^2} (\xi^2 + \xi^{*2})} \int_0^1 dt \, t \, \mathbf{e}^{-t^2 \alpha^2} J_{\nu}(\xi) J_{\nu}(t\xi^*)$$

Quenched density in the complex plane

One-point density $\rho_1(\lambda)$, for $\nu = 0$ and $\mu = 0.1$ ($\alpha = 0.591$)

The distribution of the first eigenvalue $p_1(\lambda)$ for $\nu = 0$ and $\mu = 0.1$ ($\alpha = 0.591$)



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• density (left) [Bloch, Wettig 06], "smallest" eigenvalue (right) [+ G.A., Shifrin 07]

$$\mathcal{Z}_{RMT} \equiv \int dW_1 dW_2 \prod_{f=1}^{N_f} \det \begin{bmatrix} m_f & iW_1 + \mu_f W_2 \\ \\ iW_1^{\dagger} + \mu_f W_2^{\dagger} & m_f \end{bmatrix} e^{-N \operatorname{Tr} W_j W_j^{\dagger}}$$
$$\bullet \begin{bmatrix} \mathcal{Z}_{RMT} \sim \int_{\mathbb{C}} \prod_k^N dz_k^2 w(z_k; \mu) \prod_f^{N_f} (z^2 + m_f^2) |\Delta_N(z^2)|^2 \end{bmatrix} [\text{Osborn 04}]$$

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- Vandermonde trick: for arbitrary polynomials

	1	λ_1	•••		$P_0(\lambda_1)$	$P_1(\lambda_1)$	•••
$\Delta_N(\lambda) = \prod_{k>l} \lambda_k - \lambda_j =$	1	λ_2	• • •	=	$P_0(\lambda_2)$	$P_1(\lambda_2)$	•••
	•••				• • •		

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• Orthogonal Polynomials $\int_{\mathbb{C}} d^2 z w(z;\mu;m_f) P_k(z) P_l(z)^* = h_k \delta_{kl}$

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- for $\mu=0$ weight $w(\lambda)=\lambda^{\nu}e^{-\lambda}$ Laguerre polynomials
- for $\mu \neq 0$ Laguerre polynomials on \mathbb{C} (weight exp \cdot K-Bessel) e.g.[Bender, G.A. 10]

• the partition function

- quenched: $\mathcal{Z}_{\nu,RMT} \sim \int_{\mathbb{C}} \prod_{k=1}^{N} dz_{k}^{2} w(z_{k};\mu) |\det[P_{k}]|^{2} \sim N! \prod h_{k}$

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- unquenched:

$$N_f = 1$$
: $\prod_k (\lambda_k + m^2) \Delta_N(\{\lambda\}) = \Delta_{N+1}(\{\lambda\})$ with $\lambda_{N+1} = -m^2$

The orthogonal polynomial method

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$$\mathcal{Z}_{\nu,RMT}^{(N_f)} = \left\langle \prod_f \prod_k (\lambda_k + m_f^2) \right\rangle \sim \det[L_{N+j}^{(\nu)}(-m_f^2)] / \Delta(\{m_f^2\})$$

expectation value of characteristic polynomials

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expectation value of characteristic polynomials

microscopic limit

large-N Laguerre :
$$L_{N+j}^{(\nu)}(-m_f^2) \to I_{\nu}(\hat{m}_f)$$
 gives again $\mathcal{Z}_{\varepsilon\chi}$ PT

- gap probabilities $E(\lambda) = \int_{\lambda}^{\infty} \mathcal{P}_{jpdf}$
 - equals partition function:

$$E(\lambda) = e^{-N^2 \Sigma^2 \lambda} \mathcal{Z}_{\nu=0}^{(N_f+\nu)}(\{m'\})$$

with shifted masses (irrelevant in Vandermonde)

$$m'_f = \sqrt{\lambda + m_f^2}, f = 1, \dots, N_f$$

 $m'_f = \lambda, f = N_f + 1, \dots, N_f + \nu$

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• first eigenvalue distribution

$$p_1(y) = -\partial_y E(y^2)$$
examples $(N_f = 0)$:
$$\nu = 0 \left[p_1(y) = y/2e^{-y^2/4} \right], \nu = 1 \left[p_1(y) = y/2e^{-y^2/4}I_2(y) \right] \text{etc}$$

$$\rho_k(y_1,\ldots,y_k) = \frac{1}{\mathcal{Z}} \frac{N!}{(N-k)!} \int dy_{k+1} \ldots dy_N \mathcal{P}_{jpdf}$$

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- introduce kernel of orthonormal polynomials:

$$K_N(z,v^*) = \sum_{j=0}^{N-1} P_j(z) P_j(v^*) / h_k \rightarrow \text{satisfies}$$

1.
$$\int dz w(z) K_N(z, v^*) = 1$$
, $\int dz w(z) K_N(z, z^*) = N$

2.
$$\int dz w(z) K_N(u, z^*) K_N(z, v^*) = K_N(u, v^*)$$

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$$K_N(z,v^*) = \sum_{j=0}^{N-1} P_j(z) P_j(v^*) / h_k \rightarrow \text{satisfies}$$

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• Integration Theorem Mehta:

 $\int dz_N w(z_N) \det_{k \times k} [K_N(z_i, z_J^*)] = (N - k + 1) \det_{(k-1) \times (k-1)} [K_N(z_i, z_J^*)]$

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- Compute several k-point densities from $\varepsilon \chi PT \rightarrow$ reconstruct 1st eigenvalue
- for **complex eigenvalues** a more direct computation of $p_1(r)$ is difficult also from RMT

orthogonal polynomials (real and complex eigenvalues):

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[Osborn 04, + Splittorff, Verbaarschot, GA; GA 05; Phillips, Sommers, GA.09+10]

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- NEW: corrections from finite lattice spacing

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 $S_{eff} = S_0 + aS_1 + a^2S_2 + \dots$

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$$S = \frac{1}{2} V \Sigma \text{Tr}(U + U^{\dagger}) - a^2 V W_8 \text{Tr}(U^2 + U^{\dagger} 2) -a^2 V W_6 [\text{Tr}(U + U^{\dagger})]^2 - a^2 V W_7 [\text{Tr}(U - U^{\dagger})]^2$$

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- same strategy: ε -regime (with our without RMT approach)
 - compute density and smallest eigenvalue LEC's
 - determine LEC from fit to Lattice data

Wilson $\varepsilon \chi$ PT & Wilson RMT

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- problems at $a \neq 0$:
 - D_W non-Hermitian
 - concept on **topology no longer well defined**, \exists several related densities of D_5 , real modes of D_W , chirality density

- ν zero modes get smeared out and decouple from rest of spectrum

First Wilson results



details:		
[Damgaard,		
Splittorff,	١	Ver-
baarschot	10,	+
G.A. 11]		



• partition function at equal mass (same from $\varepsilon \chi PT$ or RMT):

 $Z_{\nu}^{(N_f)} = \det[Z_{\nu+j-i}^{(N_f)}] \text{ determinant of single flavour as before:}$ $Z_{\nu}^{(N_f=1)} = \int_{-\pi}^{\pi} d\theta e^{i\nu\theta} \exp[mV\Sigma\cos(\theta) - 2a^2VW_8] \rightarrow I_{\nu}(\hat{m})$ for $W_{\nu} = W_{\nu} = 0$ (quiteble op: Gaussian integrals over 7)

- for $W_6 = W_7 = 0$ (switch on: Gaussian integrals over Z)



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• meson correlators in Wilson $\varepsilon \chi PT$: [Shindler, Necco, Bär]

 J.J.M. Verbaarschot, T. Wettig "Random Matrix Theory and Chiral Symmetry in QCD", hep-ph/0003017

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- Introduction to the mathematics of RMT: lectures G.A. \rightarrow webpage
- Yan Fyodorov, "Introduction to Random Matrix Theory", arXiv:math-ph/0412017, lectures Isaak Newton Institute