

Summary - the bigger picture

- we have introduced several ensembles of random matrices
eg. the Gaussian Orthogonal Ensemble (GOE)

= probability space on the set of all real $N \times N$ matrices $H \in \mathbb{R}^{N \times N}$
that satisfy $H = H^T$ and $\forall i=1, \dots, N$ H_{ii} distrib according to $\mathcal{N}(0, 1)$
 $\forall i, j=1, \dots, N$ H_{ij} " " " " $\mathcal{N}(0, \frac{1}{2})$ } all indep

$$\Leftrightarrow \left[P(H) = \frac{1}{C_N} e^{-\frac{1}{2} \text{Tr} H^2} \right] \text{ or } dP(H) = \frac{1}{N!} \prod_{i=1}^N \frac{e^{-\frac{1}{2} H_{ii}^2}}{\sqrt{2\pi}} \prod_{i < j} \frac{e^{-\frac{1}{2} H_{ij}^2}}{\sqrt{2\pi}} dH_{ij}$$

$C_N = \frac{2^{\frac{N(N+1)}{2}} \pi^{\frac{N(N-1)}{2}}}{2^{\frac{N(N+1)}{2}}}$

- as the matrices from such an ensemble can be diagonalised, $H = O D O^T$
with $O \in O(N)$, $D = \text{diag}(x_1, \dots, x_N)$ the eigenvalues of H (and $\text{Tr} H^2 = \text{Tr} D^2$),

we obtain the jpdf $\boxed{S(x_1, \dots, x_N) = \frac{1}{Z_{N,\beta}} \prod_{i < j} |x_i - x_j|^\beta e^{-\frac{1}{2} \sum_{i=1}^N x_i^2}}$

with GOE $\beta=1$, GUE ($H=H^*$, $H \in \mathbb{C}^{N \times N}$, indep (complex) normal) \nearrow we still need to prove this, only $N=2$

- Coulomb gas picture (rescale $x_i \Rightarrow \sqrt{N} x_i$)

$$Z_{N,\beta} = C_{N,\beta} \int_{\mathbb{R}^N} \underbrace{dx_1 \dots dx_N}_{[dx]} \exp \left[-\beta N \sum_{i=1}^N x_i^2 + \frac{\beta}{2} \sum_{i < j} \ln |x_j - x_i| \right]$$

$\equiv -\beta N^2 V[x]$ Gibbs-Boltzmann weight

- Counting function = empirical spectral measure

$$n_N(x) = \frac{1}{N} \sum_{i=1}^N \delta(x - x_i) \quad \text{claim: for } \lim_{N \rightarrow \infty} n_N(x) = n_*(x) \text{ minimises}$$

$$\mathcal{F}_0[n(x)] = \frac{1}{2} \int dx x^2 n(x) - \frac{1}{2} \iint dx dx' n(x) n(x') \ln |x - x'| \quad \text{energy functional}$$

with $Z_{N,\beta} \approx \exp(-\beta N^2 \mathcal{F}_0[n_*(x)])$

where we also computed lower orders $+ \left(\frac{\beta}{2} - 1\right) N \int dx n(x) \ln(x)$ - c.c.

- check: $Z_{N,\beta=2}$ known exactly (to be derived) agrees asymptotically with result from inserting $n_*(x) = \frac{1}{\pi} \sqrt{2-x^2} = S_{sc}(x)$ the Semi-Circle

To do: derive $n_x(x)$ from saddle point analysis of

$$\left(\int_0^1 dx [u(x)] - N \left(\int dx u(x) - 1 \right) \right)$$
 energy + Lagrange multiplier today (10/1)

reminder in 1D: $I = \int dx e^{-N f(x)}$ gives leading contribution at minimum x_0
 for $N \gg 1$ $\approx \sqrt{\frac{2\pi}{N f''(x_0)}} e^{-N f(x_0)}$ $f'(x_0) = 0 \Rightarrow$ Taylor $f(x) = f(x_0) + \frac{1}{2} (x-x_0)^2 f''(x_0)$

here: $\frac{\delta \int_0^1 dx [u(x)]}{\delta u(x)} \stackrel{!}{=} 0$, $\frac{\delta \int_0^1 dx [u(x)]}{\delta \lambda} = 0$

• the solution to be found must have compact support!

Literature: to Coulomb gas picture:

- S. Saffary 1709.04089
- Peter J. Forrester: Log-Gases and RM (e.g. chapter 14, 1) to general RM 2010
- M.L. Mehta Random Matrices, 3rd Ed. 2004
- G. Anderson, A. Guionnet, O. Zeitouni in Intro to RM, 2010
- Y Fedotkin Intro to RM: GUE and beyond multiple 04/2017

Notice: $S_n(x) = \langle u(x) \rangle \rightarrow S(x)$ too, to be shown

$$= \int dx_2 \dots dx_n S_n(x, x_2, \dots, x_n)$$

5. Saddle-point analysis of the Coulomb gas partition function

- also called steepest descent method or stationary phase approximation ⁽ⁱ⁾ of an integral with a large parameter N

$I = \int dx e^{-N f(x)}$ expand around ^(minimum) extremum x_0 with $f'(x_0) = 0$

$$\Rightarrow I \approx e^{-N f(x_0)} \int_{-\infty}^{\infty} dx e^{-\frac{N f''(x_0)}{2} (x-x_0)^2} = \sqrt{\frac{2\pi}{N f''(x_0)}} e^{-N f(x_0)}$$

"fluctuations"

see e.g. Fyodorov math-ph/0412017 ch 5.2.

- for several minima we have to sum over all of them
- for $\int_{\mathcal{J}}$ with $\mathcal{J} \subseteq \mathbb{R}$ and $\min x_0 \notin \mathcal{J}$ we have to expand to linear order at $\partial \mathcal{J}$ with x_0

5.1 Lagrange multiplier and saddle-point equation

absorb in normalisation

$$Z_{N,\beta} = C_{N,\beta} \int \mathcal{D}[u, \mu] \exp[-\beta N^2 \mathcal{F}_0[u, \mu] + (\frac{\beta}{2} - 1) N \mathcal{F}_1[u, \mu] + \beta N (\ln N - \ln c) + o(N)]$$

we have to insure that the smooth random measure $u(x)$ is normalised:

$$\text{inset } \delta \left[\int_{\mathbb{R}} dx u(x) - 1 \right] = \int_{\mathbb{R}} \frac{dk}{2\pi} e^{ik \left(\int_{\mathbb{R}} dx u(x) - 1 \right)}$$

using standard δ -exp.

- rescaling $ik \rightarrow \beta N^2 k$ and keeping only the leading order, $O(N^2)$ \mathcal{F}_0 -contrib.

$$\Rightarrow Z_{N,\beta} \approx \tilde{C}_{N,\beta} \int \mathcal{D}[u, \mu] \int_{\mathbb{R}} dx \exp \left\{ -\beta N^2 \mathcal{F}_0[u, \mu] - \mu \left(\int_{\mathbb{R}} dx u(x) - 1 \right) \right\} + O(N)$$

$\mathcal{S}[u, \mu]$ ← Lagrange multiplier

$$\text{saddle-point} \approx \text{const } e^{-\beta N^2 \mathcal{S}[u_x, \mu_x]} \quad (\text{we don't consider the fluctuations})$$

with $u_x(x), \mu_x$ solving the saddle point equations (SP eq.)

$$\left\{ \begin{aligned} 0 &= \frac{\delta}{\delta \mu(x)} \mathcal{S}[u_x, \mu_x] \Big|_{\substack{\mu = \mu_x \\ u = u_x}} = \frac{x^2}{2} - \int_{\mathbb{R}} dx' u_x(x') \ln|x-x'| - \mu_x \\ 0 &= \frac{\partial}{\partial x} \mathcal{S}[u_x, \mu_x] \Big|_{\substack{\mu = \mu_x \\ u = u_x}} = \int_{\mathbb{R}} dx' u_x(x') - 1 \end{aligned} \right. \quad (5.9)$$

used $\frac{\delta \ln|y|}{\delta y} = \frac{1}{y}$

* the normalized solution $u_x(x)$ of the SP eq. (equilibrium measure) cannot be supported on all of \mathbb{R} :

$$\text{for } x \rightarrow 1 \text{ we have } 0 = \frac{x^2}{2} - \int_{\mathbb{R}} dx' u_x(x') \ln|x-x'| - x x$$

$$\approx \frac{x^2}{2} - \ln|x| \int_{\mathbb{R}} dx' u_x(x') = \frac{x^2}{2} - \ln|x| \text{ which is inconsistent}$$

\Rightarrow the solution $u_x(x) = u_x(x; a, b)$ has compact support on $[a, b] \subseteq \mathbb{R}$
 - the parameters a, b have to be determined s.t. $u_x(x; a, b)$ mda F !

5.2 Derivative of the SP eq and Tricomi's solution

$\ln|x-x'|$ is not differentiable at $x=x'$

Def weak derivative: For $u(x), \varphi(x) \in \mathcal{F}^1([a, b])$ φ is the weak derivative of u if $\forall \psi(x) \in \mathcal{C}^\infty, \psi(a) = \psi(b) = 0$: $\int_a^b dx u(x) \psi'(x) = - \int_a^b dx \varphi(x) \psi(x)$

Here set $u(x) = \int dx' u_x(x') \ln|x-x'|$, we will regularize the singularity at $x=x'$ by ε

$$\int dx \varphi'(x) \int dx' u_x(x') \frac{\ln|x-x'|}{\varepsilon} = \frac{1}{\varepsilon} \lim_{\varepsilon \rightarrow 0} \int dx \varphi'(x) \int dx' u_x(x') \ln[(x-x')^2 + \varepsilon^2]$$

$$= -\frac{1}{2} \int dx \varphi(x) \int dx' u_x(x') \frac{2(x-x')}{(x-x')^2 + \varepsilon^2} = - \int dx \varphi(x) \text{Pr} \int dx' \frac{u_x(x')}{x-x'}$$

or $\int dx'$ or P.V. $\int dx'$ is Cauchy's principal value

i.e. for x singular point of $F(x)$

$$\underline{D_c} \left(\text{Pr} \int dx' F(x') \right) = \lim_{\varepsilon \rightarrow 0} \int_{x-\varepsilon}^{x+\varepsilon} dx' F(x') + \int dx' F(x')$$

$$\Rightarrow \text{weak derivative of SP eq: } \boxed{0 = x - \text{Pr} \int dx' \frac{u_x(x')}{x-x'}} \quad (5.15)$$

Assume that $[a, b]$ is the single interval support of $f(x)$ (why: $\underbrace{\text{true for } W}$, not true for W)

Thm [Tricomi 85]

$$g(x) = \text{Pr} \int_a^b dx' \frac{f(x')}{x-x'} \Rightarrow f(x) = \frac{1}{\pi \sqrt{(x-a)(b-x)}} \left(c - \text{Pr} \int_a^b \frac{dt \sqrt{(t-a)(b-t)}}{x-t} g(t) \right)$$

↑
arb. (integration) constant

hence $g(x) = t \Rightarrow$ (Exercise 4.1), using $\int_a^b dx \, n_x(x) = 1$ \odot

$$n_x(x) = \frac{1}{\pi \sqrt{(x-a)(b-x)}} \left[1 - x^2 + (a+b)x + \frac{1}{8}(b-a)^2 \right] \quad (5.17)$$

* this is a solution of (5.15) for any choice of a and b !

\Rightarrow we still have to determine these s.t. F is minimised.

5.3 Evaluation of $F_0[n_x(x)]$ and determination of a, b

\bullet $\int_a^b dx \, n_x(x) \cdot (5.9) \stackrel{\odot}{\Rightarrow} \frac{1}{2} \int_a^b dx \, n_x(x) x^2 - n_x = \int_a^b dx \, dx' \, n_x(x) n_x(x') \ln|x-x'|$
sp eqs. Catalan²

\odot set $x=a$ in (5.9) $\Rightarrow 0 = \frac{a^2}{2} - \int_a^b dx' \, n_x(x') \ln|a-x'| - n_x \Rightarrow n_x = \dots$

$$\Rightarrow F_0[n_x(x)] = \frac{1}{2} \int_a^b dx \, n_x(x) x^2 - \frac{1}{2} \int_a^b dx \, dx' \, n_x(x) n_x(x') \ln|x-x'|$$

$$= \frac{1}{4} \int_a^b dx \, n_x(x) x^2 + \frac{a^2}{4} - \frac{1}{2} \int_a^b dx \, n_x(x) \ln(x-a)$$

\bullet inserting $n_x(x)$ this can be solved as a function of (a, b) and minimised w.r.t. a and b . Result $a = -b = \sqrt{2}$

Shortcut: we expect $n_x(x)$ to be symmetric $\Rightarrow b = -a$

$$\Rightarrow n_x(x) = \frac{1}{\pi \sqrt{a^2 - x^2}} \left[1 - x^2 + \frac{1}{2} a^2 \right]; \text{ for } a = \sqrt{2}$$

we get $= \frac{1}{\pi} \sqrt{2-x^2}$, so we only need to show that

$$\text{Pr } \int_{-\sqrt{2}}^{\sqrt{2}} dx \frac{\sqrt{2-x^2}}{\pi(x-x')} = \lim_{\epsilon \rightarrow 0} \left(\int_{-\sqrt{2}}^{-\epsilon} + \int_{\epsilon}^{\sqrt{2}} \right) dx \frac{\sqrt{2-x^2}}{\pi(x-x')} = -x$$

\rightarrow Exercise 4.1. verify this with the anti-derivative of the integrand given in eqs (5.26), (5.27) in the book LNV.

\bullet there are other proofs for the semi-circle, e.g. using combinatorics: $\begin{pmatrix} n \\ k \end{pmatrix} \begin{pmatrix} n-k \\ l \end{pmatrix}$
 Ex 4.3 shows that g_{sc} is the generating function for Catalan numbers. 21

* Q: Is the semi-circle density universal?

- for the deformation 1. = Wigner ensemble (yes), as long as the H_{ij} are iid, that is they have the same variance
- ρ_{sc} also appears in other applications e.g. as the limiting density of the adjacency matrix of certain random graphs (Erdős-Rényi)
- for the deformation 2. = invariant ensembles (No):

(or $P(H) \sim \exp(-N\beta \text{Tr} \phi(H))$ e.g. $\phi(H) = \frac{1}{2} H^2 + \frac{1}{3} H^4 \dots$)

We have $\rho(x_1, \dots, x_n) = \frac{1}{Z} \prod_{i,j} |x_i - x_j|^\beta e^{-N\beta \sum_{i=1}^n \phi(x_i)}$ polynomial

\Rightarrow in $\mathcal{F}_0(n, \omega)$ replace $\int dx \frac{1}{2} x^2 n(x) \rightarrow \int dx \phi(x) n(x)$

$\Rightarrow \dots \Rightarrow$ SP equation $\left[\phi'(x) = P_n \int dy \frac{n(y)}{x-y} \right]$

has a diff. eq. $n_x(x) = \frac{1}{x} \left(\frac{1}{2} + g a^2 + 2g x \right) \sqrt{a^2 - x^2}$, with $a = \frac{a^2}{2} + \frac{3}{2} g a^4$ eq. for endpoints.

\rightarrow Exercise 4.2 for $\phi(x) = x - \alpha \ln x$ gives $S_{MP}(x)$ Maschke-Pastur (also called quarter-circle)

Nevertheless also for deformation 2. the local statistics e.g. the spacing distribution remain universal!

6. Derivation of the SPDF (part I) and the volume of $OU(N)$

6.1 As a reminder: change of variables in 2 dimensions:

Consider $I_1 = \int_{\mathbb{R}^2} dx dy \rho_1(x, y)$, $I_2 = \int_{\mathbb{R}^2} dx dy \rho_2(x, y)$

with $\rho_1(x, y) = f(x^2 + y^2)$, $\rho_2(x, y) = x f(x^2 + y^2)$ for some $f(z)$ s.t.

both integrals converge, e.g. $f(z) = e^{-z}$.

Standard strategy: change of variables from Cartesian $(x, y) \in \mathbb{R}^2$

to polar coordinates $(r, \theta) \in \mathbb{R}_+ \times [0, 2\pi]$: $x = r \cos \theta$, $y = r \sin \theta$

$$\text{Jacobian } J(r, \theta) = \left| \frac{\partial(x, y)}{\partial(r, \theta)} \right| = \begin{vmatrix} \frac{\partial x}{\partial r} & \frac{\partial x}{\partial \theta} \\ \frac{\partial y}{\partial r} & \frac{\partial y}{\partial \theta} \end{vmatrix} = \begin{vmatrix} \cos \theta & -r \sin \theta \\ \sin \theta & r \cos \theta \end{vmatrix} = r(\cos^2 \theta + \sin^2 \theta) = r$$

$$\Leftrightarrow \int_{\mathbb{R}^2} dx dy = \int_0^\infty dr \int_0^{2\pi} d\theta |J(r, \theta)| = \int_0^\infty dr r \int_0^{2\pi} d\theta$$

• due to $x^2 + y^2 = r^2(\cos^2 \theta + \sin^2 \theta) = r^2$ the integral I_1 is simple, as its integrand only depends on r :

$$I_1 = \int_0^\infty dr \int_0^{2\pi} d\theta f(r^2) = 2\pi \int_0^\infty dr f(r^2), \text{ while } g_2(r, \theta) = r \cos \theta f(r^2)$$

marginal $g_1(r) = r f(r^2) 2\pi$

still contains 2 integrals
(here: decouple, goes 0)

6.2. The jpdf of the GOE (part I)

$$\text{GOE: we have } P(H) = C_N e^{-\frac{1}{2} \text{Tr} H^2} = \frac{1}{2^{\frac{N}{2}} \pi^{\frac{N(N-1)}{4}}} \frac{1}{N!} e^{-\frac{1}{2} \sum_{i=1}^N \lambda_i^2} \frac{1}{\pi^{\frac{N}{2}}} e^{-\frac{1}{2} \sum_{i=1}^N \lambda_i^2}$$

$$\text{s.t. } \int_{\mathbb{R}^{N \times N}} P(H) = 1$$

• all matrices $H = H^T$, $H \in \mathbb{R}^{N \times N}$ can be diagonalised by an orthogonal basis $H = O X O^T$, $O \in O(N)$, $X = \text{diag}(\lambda_1, \dots, \lambda_N)$

$$\text{and } \text{Tr} H^2 = \text{Tr}(O \underbrace{X O^T O X^T}_{X^2}) = \text{Tr} X^2 = \sum_{i=1}^N \lambda_i^2$$

* We are in the same situation as for integral I_1 , we only need the

$$\text{Jacobian } J(X, O) = \det \left(\frac{\partial H_{ij}}{\partial (x_i, o_{ij})} \right)$$

$$dP(H) = P(\lambda_1, \lambda_2, \dots, \lambda_N) \prod_{i=1}^N d\lambda_i = C_N e^{-\sum_{i=1}^N \lambda_i^2} |J(X, O)| dO \prod_{i=1}^N dx_i$$

product of Lebesgue on \mathbb{R}

flat Lebesgue on \mathbb{R}^N