

# Lecture 4

## Summary:

\* Complex eigenvalue correlation functions

$$R_N(z_1, \dots, z_N) = \frac{1}{h} \omega(z) \det \left[ K_N(z_i, \bar{z}_j) \right]_{i,j=1}^N \quad \text{given as } N \times N \text{ det}$$

in terms of the kernel of OP  $K_N(z, \bar{w}) = \sum_{e=0}^{N-1} \frac{P_e(z) \overline{P_e(w)}}{h_e}$

(e.g. density  $R_1(z_1) = \omega(z_1) K_N(z_1, \bar{z}_1)$ )

\* joint density of radii for  $\omega = \omega(|z|)$ ,  $r_j = |z_j|$   $j=1, \dots, N$

$$\frac{P}{N}(r_1, \dots, r_N) = \frac{1}{h} \omega(r) \text{Per} \left[ r_i^{2j-1} \right]_{i,j=1}^N$$

radii become indep random var. with different distrib for  $r_k$

Permanent  $\text{Per}[A] = \sum_{\sigma \in S_N} \prod_{e=1}^N A_{e\sigma(e)}$

as det, without signs  $\Rightarrow$  is also linear, we can multiply in factors into rows (columns) (e.g.  $\frac{1}{h} \omega(r) \text{Per} [r_i^{2j-1}] = \text{Per} [r_i^{2j-1}]$ )

\* distribution of  $k$ th individual ev (for Hermitian real  $\Rightarrow$  real ev  $\rho$ )

$$P_k(s) = h \binom{N}{k} \int_0^s \int_0^s \dots \int_0^s \int_s^\infty \dots \int_s^\infty \frac{1}{C_{N-k}} \rho(\lambda_1, \dots, \lambda_{N-k}, s, \dots, s) d\lambda_1 \dots d\lambda_{N-k}$$

\*  $k$ -th gap probability

$$E_k(s) = \frac{N!}{(N-k)!} \int_0^s \int_0^s \dots \int_0^s \int_s^\infty \dots \int_s^\infty \frac{1}{C_N} \rho(\lambda_1, \dots, \lambda_N) d\lambda_1 \dots d\lambda_N$$

$\rightarrow$  examples

\* relation  $E_k(s)$  to  $p_k(s)$ :

$$\mathcal{D}_s E_0(s) = -p_0(s)$$

smallest ev distribution

$$\mathcal{D}_s E_k(s) = k! (p_k(s) - p_{k+1}(s)) \quad \text{for } k=1, \dots, N-1$$

$$\text{(also } \mathcal{D}_s E_N(s) = N! p_N(s) \text{)}$$

largest ev distrib.

$\Rightarrow$  telescopic sum 
$$p_k(s) = - \sum_{e=0}^{k-1} \frac{1}{e!} \mathcal{D}_s E_e(s) \quad \text{for } k=1, \dots, N$$

Exercise 11: proof

\* not all  $E_k(s)$  are indep. :  $1 - E_0(s) \sim$  probab that at least 1 cv  $< s$   
 can be expressed through cell  $E_1(s), \dots, E_N(s)$

• Relation between  $k$ -point correlation functions and  $k$ -th gap probab:

idea:  $(a-b)^j = \sum_{e=0}^j \binom{j}{e} a^{j-e} (-b)^e$ , apply to  $\int_0^s = \int_0^\infty - \int_0^s$

using the symmetry of  $P_N(\lambda_1, \dots, \lambda_N)$  under relabelling

$$\Rightarrow E_k(s) = \frac{N!}{(N-k)!} \int_0^s d\lambda_1 \dots \int_0^s d\lambda_k \int_0^\infty d\lambda_{k+1} \dots \int_0^\infty d\lambda_N \frac{1}{C_N} P_N(\lambda_1, \dots, \lambda_N)$$

$$= \sum_{e=0}^{N-k} (-1)^e \binom{N-k}{e} \left( \int_0^s \right)^{N-k-e} \left( \int_0^\infty \right)^e d\lambda_{k+1} \dots d\lambda_N$$

using the definition of the

$(k+e)$ -point correl funct.  $R_{k+e}$  given as  $N-(k+e)$  integral  $\int_0^\infty$  over  $\frac{1}{C_N} P_N(\lambda_1, \dots, \lambda_N)$

$$E_k(s) = \sum_{e=0}^{N-k} \frac{(-1)^e}{e!} \int_0^s d\lambda_1 \dots \int_0^s d\lambda_k R_{k+e}(\lambda_1, \dots, \lambda_{k+e})$$

$= \det_k \left( \lambda_i^{-j} \frac{1}{h} \omega(\lambda) \right)_{(k+e) \times (k+e)}$

\* for  $k=0$  e.g. 
$$E_0(s) = \sum_{e=0}^N \frac{(-1)^e}{e!} \int_0^s d\lambda_1 \dots \int_0^s d\lambda_e \det [k_{ij}(\lambda_i, \lambda_j)]_{i,j=1}^e \frac{1}{h} \omega(\lambda_m)$$

where the  $e=0$ -term is 1

This is the definition of the Fredholm determinant for the kernel  $k_N$

as an integral operator on  $(L^2(\tilde{\omega}))$ :

$$(k_N f)(x) = \int_0^\infty dy \tilde{\omega}(y) k_N(x, y) f(y), \quad \text{with } \tilde{\omega}(y) = \omega(y) \mathcal{K}_{(0, s)}^{(g)}$$

here

$$E_0(s) = \det(1 - k_N | L^2(\tilde{\omega}))$$

proof: see chapter 4 in Oxford handbook on RMT

\* in the example with  $k=0$  <sup>(comp B)</sup> we have obtained the 2nd order term as unity:

$$E_0(s) = 1 - \int_0^s dt_1 R_1(t_1) + \frac{1}{2} \int_0^s dt_1 \int_0^s dt_2 R_2(t_1, t_2) - \frac{1}{3!} \int_0^s dt_1 \int_0^s dt_2 \int_0^s dt_3 R_3(t_1, t_2, t_3) + \dots$$

\* for some considerations, e.g. the smallest ev distrib, this sum sometimes "converges" very rapidly, see hep-th/0311171

• generating function:

$$\text{define } E(s; \frac{1}{z}) = 1 - \sum_{e=1}^N (-z)^e \frac{1}{e!} \int_0^s dt_1 \dots \int_0^s dt_e R_e(t_1, \dots, t_e)$$

it holds that  $E_k(s) = (-1)^k \frac{\partial^k}{\partial z^k} E(s; \frac{1}{z}) \Big|_{z=1}$  for  $k=0, 1, \dots, N$

(where  $k=0$  has no derivative).

Exercise A2 proof

Q: How can we generalise this concept to ev in the complex plane, where there is no unique ordering procedure?

→ radial ordering of ev  $z_i$  w.r.t their modulus  $r_i = |z_i|$

But: 2 ev can have  $r = |z_i| = |z_{i+1}|$  without penalty,



→ level repulsion?

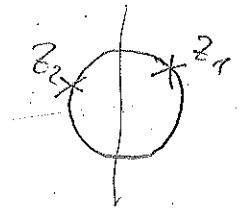
• Distribution of individual ev and gap probabilities in  $\mathbb{C}$

[See e.g. 0901.0897]  
other con-  
variant  $\frac{1}{k}$

On  $\mathbb{C}$  there is no unique ordering for the eigenvalues  
 $\rightarrow$  we will proceed with radial ordering

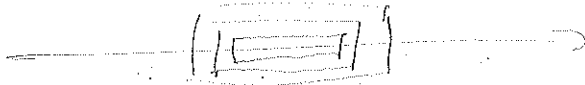
$$\text{s.t. } |z_1| \leq |z_2| \leq \dots \leq |z_N|$$

\* clearly, unlike for real ev the Coulomb repulsion is not necessarily felt for 2 ev with the same radii  $|z_1| = |z_2|$ , unless



$|z_1 - z_2|$  is small! (we have seen that considering only radii these become indep.)

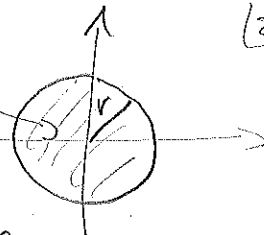
\* other orderings are possible (and more useful) e.g. for the elliptic Ginibre close to the GUE w/ some family of contours, e.g.



Radial gap probability

$$E_k(r) = \frac{N!}{(N-k)!} \int_0^r dr_1 r_1 \dots \int_0^r dr_k r_k \int_{r_k}^{\infty} dr_{k+1} r_{k+1} \dots \int_r^{\infty} dr_N r_N \prod_{\substack{1 \leq i < j \leq N \\ i \neq k}} \frac{1}{|z_i - z_j|} \prod_{i=1}^N \rho(z_i)$$

$\approx$  probability to have  $k$  ev inside and  $N-k$  outside of the disk around  $0$  of radius  $r$   $D_r$



$\int_{\mathbb{C}}$  all ev integrated out

e.g.  $E_0(r)$  prob that this circle is empty

$$\Rightarrow \left| \frac{\partial}{\partial r} E_k(r) = r (P_k(r) - P_{k+1}(r)) \right| \quad (\text{sub } P_0(r) \equiv 0)$$

where  $P_k(r)$  is the radial distribution of the  $k$ -th ev in radial ordering. Ex  $\rightarrow$  write down the def of  $P_k(r)$

$$\Rightarrow P_k(v) = -\frac{\Delta}{v} \frac{\partial}{\partial v} \sum_{\ell=0}^{k-1} E_\ell(v)$$

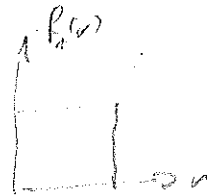
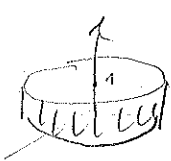
\* relations to the density  $R_n(z_n)$  ( $N = \int_{\mathbb{C}^2} R_n(z) = \int_0^\infty \int_0^{2\pi} R_n(r) r dr d\theta$ )

$$\int_0^{2\pi} \int_0^\infty R_n(r) r dr d\theta = \sum_{k=1}^N P_k(v), \quad \text{with} \quad \int_0^\infty P_k(v) dv = 1 \quad \forall k=1, \dots, N$$

Note that we could also call  $\tilde{P}_k(v) = v P_k(v)$  the radial distrib. w.r.t flat measure, and

$$\tilde{R}_n(v) = v \int_0^{2\pi} R_n(r) r dr d\theta \quad \text{the radial density}$$

e.g. for Ginibre  $\Rightarrow R_1(z) = \int_0^{2\pi} R_1(r) r dr d\theta = 2\pi R_1(v)$



But  $\tilde{R}_1(v)$  "triangular law"

### Fredholm determinant representation:

as for  $v \in \mathbb{R}$  we can show that the det of  $E_v$  is equivalent to

$$E_k(v) = \sum_{\ell=0}^{N-k} \frac{(-1)^\ell v^{k+\ell}}{\ell!} \int_0^{2\pi} \int_0^\infty \prod_{j=1}^k R_{\ell+1}(z_j) R_{\ell+1}(z_{k+\ell+1}) dz_j$$

1st term is  $\sum_{j=1}^k$   
↓  
 $k=0, 1, \dots, N$

and a gen. funct

$$E_k(v; z) = \left( \prod_{j=1}^k z_j \right) \Rightarrow E_k(v) = (v)^k \frac{\partial^k}{\partial z^k} E(v; z) \Big|_{z=1}$$

alternatively we can write with  $\Psi_j(z) = \omega(z)^{\frac{1}{2}} \frac{P_j(z)}{h_j^{\frac{1}{2}}}$  ON wave function from OP

$$E_0(v) = \frac{1}{C_N} \frac{1}{h} \int_{\mathbb{C}^2} \omega(z) |\Delta_N(z)|^2 = \frac{1}{N!} \frac{1}{h} \int_{\mathbb{C}^2} \left| \det[\Psi_{j-1}(z_i)]_{j,i=1}^N \right|^2$$