

Summary

Lecture 9

* We have seen an example for the product of 2 coupled random matrices J_1, J_2 , $Y_2 = J_1 J_2$

$$P(J_1, J_2) \sim \exp\left[-\frac{1}{1-c} \text{Tr} (J_1 J_1^T + J_2 J_2^T - c(J_1 J_2 + J_2^T J_1^T))\right]$$

that can be solved using Laguerre polynomials on \mathbb{C}

$$\text{with kernel } K_N(z, \bar{w}) \sim \sum_{\ell=0}^{N-1} \frac{z^{2\ell} \bar{w}^\ell}{(\ell!)^2} L_c^{(\nu)}\left(\frac{z^2}{c}\right) L_c^{(\nu)}\left(\frac{\bar{w}^2}{c}\right)$$

$$\text{and weight } W(z^2) = |z|^{2\nu+2} e^{-\frac{c}{1-c}(z^2 + \bar{z}^2)} K_\nu\left(\frac{2|z|^2}{1-c}\right) \\ \uparrow \\ \sim G_{02}^{20}\left(\frac{1}{c} \mid \frac{2z^2}{1-c}\right)$$

* we have started to investigate the large- N limit distinguishing global/local statistics as for Hermitian RMT

• The large- N limit(s) of the Ginibre ensemble ($\nu=0$ for simplicity)

k -point correlation functions: p.6

$$R_k(z_1, \dots, z_k) = \frac{1}{k!} \prod_{j=1}^k \omega(z_j) \det [K_N(z_i, \bar{z}_m)]_{1 \leq i, m \leq k}$$

$$= \det \left[\underbrace{\omega(z_i)^{\frac{1}{2}} \omega(\bar{z}_m)^{\frac{1}{2}} K_N(z_i, \bar{z}_m)}_{K_N(z_i, \bar{z}_m)} \right] \quad \omega(z) = e^{-|z|^2}$$

\Rightarrow we need to study the limit of weights, kernel

* "naive limit" for Ginibre, p.7

$$\lim_{N \rightarrow \infty} e^{-\frac{1}{2}|z|^2} e^{-\frac{1}{2}|\bar{z}|^2} \frac{N^{-1}}{2} \frac{(z\bar{z})^k}{k!} e^{\frac{1}{N} z\bar{z}} = \begin{pmatrix} -\frac{1}{2}|z|^2 & -\frac{1}{2}|\bar{z}|^2 & z\bar{z} \\ e & \frac{1}{N} e & \end{pmatrix} = \frac{K_{\text{Gin}}(z, \bar{z})}{N}$$

$|z|, |\bar{z}| = o(N), N \text{ indep!}$

$$\Rightarrow \lim_{N \rightarrow \infty} R_k(z_1, \dots, z_k) = \frac{1}{k!} e^{-|z_i|^2} \det [e^{z_i \bar{z}_m}]_{1 \leq i, m \leq k}$$

Distributing
Ginibre-kernel

e.g. $k=1$ spectral density $R_1(z) = \frac{1}{N} K_{\text{Gin}}(z, \bar{z}) = e^{-|z|^2} \frac{1}{N} e^{z\bar{z}} = \frac{1}{N}$

constant density

$k=2$ two-point function: $R_2(z_1, \bar{z}_2) = e^{-|z_1|^2} e^{-|z_2|^2} \frac{1}{N^2} \det \begin{pmatrix} e^{z_1 \bar{z}_1} & e^{z_1 \bar{z}_2} \\ e^{z_2 \bar{z}_1} & e^{z_2 \bar{z}_2} \end{pmatrix}$

only depends on the
distance $z_1 - z_2$, both $R_1, R_2 = o(N)$

$$= \frac{e^{-|z_1|^2 - |z_2|^2}}{N^2} \left(e^{|z_1|^2 + |z_2|^2} - e^{z_1 \bar{z}_2 + z_2 \bar{z}_1} \right)$$

$$= \frac{1}{N^2} \left(1 - e^{-\frac{(z_1 - z_2)(\bar{z}_1 - \bar{z}_2)}{|z_1 - z_2|^2}} \right)$$

Are these local or global correlation functions?

(factored as $R_1(z_1)R_1(z_2)$ - part)

→ For this we need to find out the support of the density close to origin on a global scale:

recall p. 2: $K_n(z, \bar{u}) = \sum_{\ell=0}^{n-1} \frac{(z\bar{u})^\ell}{\ell!} = \frac{e^{z\bar{u}}}{\bar{u}} \frac{\Gamma'(N, z\bar{u})}{\Gamma(N)}$

where $\Gamma'(N, x) = \int_x^\infty dt t^{N-1} e^{-t}$ is the incomplete gamma function

with $\Gamma'(N, 0) = \Gamma(N) = (N-1)!$ asymptotic from Stirling formula

⇒ $R_n(z) = e^{-|z|^2} e^{|z|^2} \frac{1}{\pi} \frac{\Gamma'(N, |z|^2)}{\Gamma(N)} = \frac{1}{\pi} \frac{\Gamma'(N, |z|^2)}{\Gamma(N)}$ and it is well known

that $\lim_{N \rightarrow \infty} R_n(z) = \lim_{N \rightarrow \infty} \frac{1}{\pi} \frac{\Gamma'(N, |z|^2)}{\Gamma(N)} \approx \frac{1}{\pi} \Theta(\sqrt{N} - |z|)$ with $\int_{\mathbb{C}} R_n^2(z) = 1$
Heaviside function.

Exercise show this using

→ a saddle point approx. to the integral for $\Gamma'(N, x)$, around $x = \sqrt{N}$

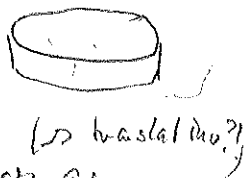
* This is the circular law, that is for $|z| \leq \sqrt{N}$ the density is constant, so to get a compact support we need to rescale matrix elements or complex ev as in the GUE

$z_j \rightarrow \sqrt{N} \hat{z}_j \quad \forall j=1, \dots, N$

⇒ $P(\hat{z}) \sim e^{-N \sqrt{N} \hat{z}^2}$ has compact support, with

$\underline{Q_c(z)} = \lim_{N \rightarrow \infty} R_n(\hat{z}) = \lim_{N \rightarrow \infty} \frac{1}{\pi} \frac{\Gamma'(N, N \sqrt{N} \hat{z}^2)}{\Gamma(N)} = \frac{1}{\pi} \Theta(\sqrt{N} - \sqrt{N} \hat{z}) = \frac{1}{\pi} \Theta(1 - \hat{z})$

which is now normalised to unity $\int_{\mathbb{C}} Q_c^2(z) = 1$



⇒ the previous "naive limit" is a microscopic limit close to the origin as

$\lim_{N \rightarrow \infty} \frac{1}{N} \sum_{\ell=0}^{N-1} \frac{(z\bar{u})^\ell}{\ell!} = e^{-\frac{1}{N} |z|^2 |u|^2} \sum_{\ell=0}^{N-1} \frac{N (z\bar{u})^\ell}{\ell!}$ need a rescaling $\hat{z} \rightarrow \frac{1}{\sqrt{N}} \hat{z}$

* because the density is flat on the unit disc, we may expect that the limiting

$$\text{Ginibre - Kernel } K_{\text{Gin}}(z, \bar{w}) = \frac{1}{\pi} e^{-\frac{|z|^2 - |w|^2}{2} + z\bar{w}}$$

holds everywhere on $|z| < 1$, due to translational invariance.

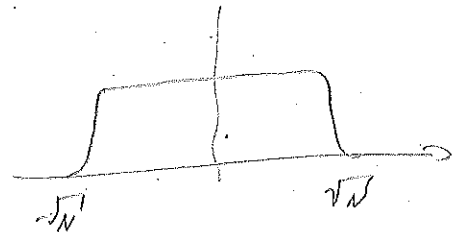
o Limiting kernel at the edge: complementary error function etc

- in the original scaling $P(y) \sim e^{-y^2/4}$ we have seen

$$\text{that } R_N(z) = \frac{1}{\pi} \frac{\Gamma(N, |z|^2)}{\Gamma(N)} \sim \frac{1}{\pi} \Theta(\sqrt{N} - |z|)$$

so let's zoom into the

$$\text{vicinity of the edge: } z = (\sqrt{N} + x) e^{i\varphi}$$



$$\Rightarrow \lim_{N \rightarrow \infty} R_N(z = (\sqrt{N} + x)e^{i\varphi}) \approx \frac{1}{2\pi} \text{erfc}(\sqrt{2}x)$$

Exercise 26

$$\text{where } \text{erfc}(x) = 1 - \text{erf}(x) = \frac{2}{\sqrt{\pi}} \int_x^{\infty} dt e^{-t^2}$$

with asymptote	$\lim_{x \rightarrow -\infty} \text{erfc}(x) = 2$	$\lim_{x \rightarrow +\infty} \text{erfc}(x) = \frac{e^{-x^2}}{\sqrt{\pi} x}$
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for the limiting kernel see e.g. 1312.0068, Prop. 8.11

* note that it can be shown, that the eigenvalues, largest in modulus is Gumbel distributed

$$|P(x)| = e^{-e^{-x}} = \sum_{n=0}^{\infty} \frac{(-1)^n}{n!} (e^{-x})^n = \sum_{n=0}^{\infty} \frac{(-1)^n}{n!} \int_x^{\infty} \int_x^{\infty} e^{-x_1 - x_2} dx_1 dx_2$$

like indep. Poisson distrib. random variables

n degenerate det kernel

see 0808.2608

so far we have: global - circular law

local - Ginibre kernel bulk (origin),
sfc - kernel edge

- Other limiting kernels? Yes, for products and elliptic Gin
- Are these limiting kernels universal, that is do they hold for distributions of matrix element J_{ij} different from $\text{Gin} \propto e^{-\text{Tr} J^2}$?

(~~Ex~~ in Lec: it is difficult to write down convergent $P(z)$ such that eV, z_1, \dots, z_N decouple) Yes, e.g. for Wigner matrices:

J_{ij} indep random variables, s.t. certain (all) moments exist

T. Tao, V. Vu 1206.1893

not much more rigorous results for non-Hermitian RMT (e.g. 1610.06519)

→ the weak non-Hermiticity limit [Fyodorov, Khavichenko, Sommers
chao-dyn/9802025]

it interpolates between some-kernel GUE
and Ginibre kernel, as we will show

→ *Consider the elliptic Ginibre ensemble (p. 7) $z = x + iy$

wigner $\omega(z) = \det \left(-\frac{1}{1-i^2} \left(\frac{1-z^2}{x^2+y^2} - \frac{z}{2(x^2+y^2)} (z^2 + \bar{z}^2) \right) \right) = \exp \left(-\frac{x^2}{1+i} - \frac{y^2}{1-i} \right)$

kernel $K_n(z, \bar{w}) = \sum_{\ell=0}^{n-1} \frac{1}{\ell!} \left(\frac{z}{2} \right)^\ell H_\ell \left(\frac{z}{\sqrt{2i}} \right) H_\ell \left(\frac{\bar{w}}{\sqrt{2i}} \right)$

global limit: one can show (see above reference)

$$\lim_{n \rightarrow \infty} \frac{1}{n} R_n(\sqrt{n}z) = \begin{cases} \frac{1}{n(1-z^2)} & \text{if } \frac{x^2}{(1-i)^2} + \frac{y^2}{(1+i)^2} \leq 1, z = x+iy \\ 0 & \text{else} \end{cases}$$

elliptic law

Universality of the elliptic Ginibre ensemble:

"narrow limit" $\lim_{N \rightarrow \infty} K_N(z, \bar{u}) = \sum_{\ell=0}^{\infty} \frac{1}{\ell!} \left(\frac{\bar{c}}{2}\right)^\ell He_\ell\left(\frac{z}{\sqrt{2}}\right) He_\ell\left(\frac{\bar{u}}{\sqrt{2}}\right)$

Mehler formula $= \frac{1}{1-\bar{c}^2} \exp\left[\frac{\bar{c}}{2} (z\bar{u} - \frac{z^2 + \bar{u}^2}{2})\right]$

= k-point correl. function

Exercise 23, check in the limit $\bar{c} \rightarrow 0$

$$R_N(z_1, \dots, z_k) \sim \frac{1}{N^k} e^{-\frac{1}{2} (|z_j|^2 - \frac{\bar{c}}{2} (z_j^2 + \bar{z}_j^2))} \det_{1 \leq j, m \leq k} \left[\frac{1}{1-\bar{c}^2} e^{\left[\frac{z_j \bar{z}_m}{1-\bar{c}^2} - \frac{\bar{c}}{2} (z_j^2 + \bar{z}_m^2) \right]} \right]$$

take out and cancel

$$= \frac{1}{(1-\bar{c}^2)^{\frac{k}{2}}} \frac{1}{N^k} e^{-\frac{1}{2} |z_j|^2} \det_{1 \leq j, m \leq k} \left[e^{\frac{z_j \bar{z}_m}{1-\bar{c}^2}} \right]$$

same as for Ginibre kernel, after rescaling $z_j \rightarrow \frac{z_j}{\sqrt{1-\bar{c}^2}}$

Weak non-Hermiticity

= universal!

* We consider simultaneously (= double scaling limit)

$N \rightarrow \infty$ and $\tau \rightarrow 1$, the Hermitian limit: $\lim_{\tau \rightarrow 1} \frac{1}{N^k (1-\tau)^k} e^{-\frac{y^2}{2(1-\tau)}} = \delta(y)$

such that $\lim_{\substack{N \rightarrow \infty \\ \tau \rightarrow 1}} N(1-\tau^2) = \kappa^2$ is kept constant

\Rightarrow the global density (Ellipse \Rightarrow) collapses to the GUE semi circle

But, the local correlations still extend into the complex plane

* scaling of z_j 's:

- for compact support we need $y \rightarrow \tau z_j$, weight $e^{-\frac{N}{2} (|z_j|^2 - \frac{\bar{c}}{2} (z_j^2 + \bar{z}_j^2))}$

then $z_j = X + \frac{1}{N} y_1 + i \frac{y_2}{N} = X + \frac{z_{1,2}}{N}$ ($X \in \mathbb{R}$, $\frac{1}{N}$ deviation from X)

wigner function: $w(z_1) = \exp\left[-\frac{N}{1+\tau} \left(x + \frac{x_1}{N}\right)^2 - \frac{N}{1-\tau} \left(\frac{y_1}{N}\right)^2\right]$

$$\xrightarrow[\tau \rightarrow 1]{N \rightarrow \infty} \exp\left[-\frac{N}{2} x^2 - x x_1 - \frac{2 y_1^2}{\alpha^2}\right]$$

↑ GUE weight

i.e. I remaining non-trivial comb

* for the kernel the limit is more tricky, involving \int -rep's of the Hermite poly. and a careful saddle point analysis (see Lot). We cannot use the Mehler formula here, (non uniform)!.

result: $\lim_{\substack{N \rightarrow \infty \\ \tau \rightarrow 1}} w(z_1) \frac{1}{\alpha} K_N(z_1, \bar{z}_2) \sim \frac{e^{-\frac{y_1^2}{\alpha^2} - \frac{y_2^2}{\alpha^2} - \frac{1}{\alpha} \int dy e^{-\frac{\alpha^2}{2} u^2} \cos(u(\frac{z_1}{\alpha} - \bar{z}_2))}}{\alpha}$

(here for $x=0$ at the origin, else $\int \rightarrow \int_0^1$)

↑ this is the kernel

$$K(z_1, \bar{z}_2, \alpha)$$

lim form of hep-th/0206080

Interplay between Sine-kernel & Ginibre

GUE limit $\alpha \rightarrow 0$ $\lim_{\alpha \rightarrow 0} \frac{1}{\sqrt{\alpha}} e^{-\frac{y^2}{\alpha}} = \delta(y)$ collapse to IR

and $\int_0^1 du \cos(u(x_1 - x_2)) = \frac{1}{x_1 - x_2} \left[\sin(u(x_1 - x_2)) \right]_0^1 = \frac{\sin(x_1 - x_2)}{x_1 - x_2}$

Ginibre limit $\alpha \rightarrow \infty$: here the scaling of z_i is different (\sqrt{N} vs N in wigner)

so we have to rescale z_i 's with $\alpha \sim \sqrt{N}$

and one can show that $\lim_{\alpha \rightarrow \infty} K(\frac{z_1}{\alpha}, \bar{z}_2, \alpha) \Big|_{\frac{z_{i,2}}{\alpha} \text{ fixed}} \sim K_{\text{Gin}}$