

• In some cases additional info is available:

\* for 1-dim reps. the characters and the  $n \times n$  "matrix" rep are the same

\* for Abelian groups all irreps are 1-dimensional:

proof Schur's first lemma:

take for  $B = D^{(\mu)}(g)$  for any fixed  $g \in G$

$$\Rightarrow \text{as } G \text{ is Abelian } D^{(\mu)}(g) D^{(\nu)}(g') = D^{(\nu)}(g') D^{(\mu)}(g) \quad \forall g, g' \in G$$

as  $D^{(\mu)}(g)$  commutes with all other matrices in that rep

$$\Rightarrow D^{(\mu)}(g) = \lambda_g^{(\mu)} \mathbb{1}, \quad \text{for all } g \in G \text{ as } g \text{ was arbitrary}$$

$\Rightarrow D^{(\mu)}$  is reducible (to  $1 \times 1$  blocks) (i.e.  $\frac{1}{|G|} \sum_g \chi(g) = 1$ )

\* this can be extended to any compact Abelian group. As

We have seen for  $G$  Abelian its conjugacy classes are all

singletons, so  $k = |G|$  and  $K_i = \{g_i\} \quad i=1, \dots, |G|$  are the conjugacy classes

$$\Rightarrow |G| = \sum_{\mu=1}^{n=k=|G|} n_{\mu}^2 \quad \Rightarrow n_{\mu}^2 = 1 \quad \forall \mu \quad \text{so all irreps are 1-dim}$$

example  $C_3$ : is Abelian, so its conj. classes are  $\{e\}, \{c\}, \{c^2\}$

and  $k=3$

character table

$C_3$	$e$	$c$	$c^2$
$D^{(1)}$	1	1	1
$D^{(2)}$	1	$\omega$	$\omega^2$
$D^{(3)}$	1	$\omega^2$	$\omega$

not faithful: kernel =  $C_3$

faithful kernel =  $\{e\}$

" " " "

↑  
irreps, 1 dim

• the characters (traces) and  $1 \times 1$  dim matrix reps are the same

$\Rightarrow$  the characters themselves have to obey the group multiplication:

$$\chi(c^2) = (\chi(c))^2 \quad \text{and} \quad (\chi(c))^3 = \chi(c^3) = \chi(e) = 1$$

$\Rightarrow \chi(c)$  must be a root of unity, here  $1, \omega = e^{\frac{2\pi i}{3}}, e^{\frac{4\pi i}{3}} = \omega^2$

•  $D^{(1)}$  is the trivial rep

•  $D^{(2)}$  and  $D^{(3)}$  are complex conjugate as  $\omega = e^{\frac{-2\pi i}{3}} = e^{\frac{-2\pi i}{3} + \frac{6\pi i}{3}} = \omega^2$

for all 3 reps we have a homomorphism  $G \rightarrow GL(1, \mathbb{C})$

• as  $C_3$  is abelian it has no proper subgroups, only  $\{e\}$  and  $C_3$  are subgroups

• we recall the special relation of normal subgroups and conjugacy classes (labeling via irreps):

$H \subseteq G$  subgroup  
(left = right cosets always)

$\Rightarrow$  a normal subgroup is made of complete conjugacy classes (p. 28)

- because the kernel is a normal subgroup itself we can

distinguish irreps by their kernel: ( $= \{e\}$  or  $C_3$ )

$D^{(1)}$ : kernel  $K = C_3$ , not faithful

$D^{(2)}$  &  $D^{(3)}$  kernel  $K = \{e\}$  so both are faithful

\*  $D : G \rightarrow GL(1, \mathbb{C}) (= \mathbb{C}^\times)$  is a map from group  $G$  to the group  $GL(1, \mathbb{C})$  (with multiplication on  $\mathbb{C}$ , neutral element 1)

## Proof of orthogonality:

$$\frac{1}{[g]} \sum_{i=1}^3 \sum_{j=1}^3 \chi_i^{(\mu)} \chi_j^{(\nu)*} = \delta_{\mu\nu}$$

$\mu = \nu$ :  $\delta^{\mu\nu} = \delta^{11} = 1$ ,  $\delta^{22} = \delta^{33} = 1$  as  $(\omega\omega^*)^{1,2} = 1$

$\mu \neq \nu$ : we need to show  $\omega \neq 1$   $\boxed{1 + \omega + \omega^2 = 0}$  for  $\omega \neq 1$   
 (v=1,  $\mu=2,3$ )  
 which is true for roots of unity  $\omega^3 = 1$   
 $\Rightarrow (\omega - 1)(1 + \omega + \omega^2) = 0$

b)  $\boxed{1 + \omega\omega^{*2} + \omega^2\omega^* = 1 + \omega^* + \omega = 0}$   
 (v=2,  $\mu=3$ )  
 $e^{-\frac{2\pi i}{3}} = e^{\frac{6\pi i}{3} - \frac{2\pi i}{3}} = \omega^2$

## Direct Products of Representations and their Decomposition

- So far we have decomposed reducible reps  $D$  into their irreps,  
 e.g.  $D(g) = A(g) \oplus B(g)$  or more general  $D = \sum_{\nu} a_{\nu} D^{(\nu)}$

• Q: Can we also study products of reps?

Reminder - so far we have studied  $G$  itself as a direct product of subgroups  $A, B$  (p. 314)

$G = A \times B$  if  $\forall a \in A \forall b \in B$ ,  $ab = ba$  and  $\forall g \in G \exists! a \in A \exists! b \in B$   
 (or products of different groups: Cartesian product) s.t.  $g = ab$

Now we'd like to study also products of reps. of the same group  $G$  (not only subgroups)

Example: Consider a qm. 2-particle system, e.g. 2 electrons of atoms on a crystal (lattice) or of free atoms. These e.c. are described by spatial wave functions  $\psi_a(x), \psi_c(x)$ .

Suppose we have a symmetry group  $G$  acting on the system, and (lattice sym. or cont. rotation sym.)  $\psi_a(x)$  and  $\psi_c(x)$  transform w.r.t the two irrep  $D^{(\mu)}, D^{(\nu)}$ :  $g \in G$

$$\psi'_a(x) = D_{ab}^{(\mu)}(g) \psi_b(x) \quad \psi_{ac}(x) \equiv \psi_a(x) \psi_c(x) \quad \text{product wave fun.}$$

$$\psi'_c(x) = D_{cd}^{(\nu)}(g) \psi_d(x) \quad \text{, or}$$

$$\text{with } \psi'_{ac}(x) = \underbrace{D_{ab}^{(\mu)}(g) D_{cd}^{(\nu)}(g)}_{\equiv D_{AB}^{(\mu, \nu)}(g)} \psi_{bd}(x)$$

We can regard the product of the 2 matrices as a single matrix with indices  $A = (ac)$  (ordered pair) and  $B = (bd)$ :  $\psi'_A(x) = D_{AB}^{(\mu, \nu)}(g) \psi_B(x)$

The rep  $D_{AB}^{(\mu, \nu)}(g)$  is then called a direct product rep

of  $\boxed{D^{(\mu)} \otimes D^{(\nu)}}$ . We have to check that it is a rep itself:

$$D^{(\mu, \nu)}(g_1) D^{(\mu, \nu)}(g_2) = D_{ab}^{(\mu)}(g_1) D_{cd}^{(\nu)}(g_1) D_{ab}^{(\mu)}(g_2) D_{cd}^{(\nu)}(g_2) = D_{ab}^{(\mu)}(g_1 g_2) D_{cd}^{(\nu)}(g_1 g_2) = D^{(\mu, \nu)}(g_1 g_2)$$

Q: is  $D^{(\mu, \nu)}$  also an irrep (as were  $D^{(\mu)}$  and  $D^{(\nu)}$ )?

1) the electrons don't interact:

$\Rightarrow$  acting on  $\psi$  with  $g_1$  and on  $\psi$  with  $g_2$  is a symmetry,

provided by  $D_{ac, bd}^{(\mu, \nu)}(g_1, g_2) \equiv D_{ab}^{(\mu)}(g_1) D_{cd}^{(\nu)}(g_2)$

This is an irrep of the product group  $G \times G$

(assume that not:  $\exists$  inv. proper subspace  $\forall g_1, g_2 \in G$ , in particular) for  $g_1 = e \Rightarrow \exists$  ... of  $D^{(\mu, \nu)}$

2) the electrons do interact

⇒ we only have a residual invariance taking  $g_1 = g_2 = g$

In general then  $D^{(\mu)} \otimes D^{(\nu)}$  is reducible:

$$D^{(\mu)} \otimes D^{(\nu)} = \sum_{\oplus} a_{\sigma} D^{(\sigma)} \quad \text{called Clebsch-Gordan series}$$

• the  $a_{\sigma}$  can be determined using the orthogonal simple characters:

we have  $\chi^{(\mu, \nu)}(g) = D_{AA}^{(\mu, \nu)}(g) = D_{AA}^{(\mu)}(g) D_{CC}^{(\nu)}(g) = \chi^{(\mu)}(g) \chi^{(\nu)}(g)$

p. 54  
 $\Rightarrow a_{\sigma} = \frac{1}{|G|} \sum_{g \in G} \chi^{(\mu)}(g) \chi^{(\nu)}(g) \chi^{(\sigma)}(g^{-1}) = \langle \chi^{(\mu)} \chi^{(\nu)}, \chi^{(\sigma)} \rangle$

## Continuous Groups

Apart from some examples we have so far only dealt with discrete groups, i.e. rotations by multiples of a fixed angle  $\frac{2\pi}{n}$ , with  $n \in \mathbb{N}$ .

For obvious reasons in physics rotations with a continuous angle play an important rôle. In general we would like to treat the situation where group elements are characterised by  $n$  real variables, that is some  $x \in \mathbb{R}^n$ . For group operations we then have for  $g(x) \in G$

product  $\forall x, y \in \mathbb{R}^n \exists z \in \mathbb{R}^n$  s.t.  $\boxed{g(x) \cdot g(y) = g(z)}$  with  $z = f_1(x, y)$

inverse  $\forall x \in \mathbb{R}^n \exists w \in \mathbb{R}^n$  s.t.  $g(x) = g(w)^{-1}$

with  $f_2(x) = w$ ,  $f_2: \mathbb{R}^n \rightarrow \mathbb{R}^n$

Def: A group  $G$  is called continuous if the above applies and  $f_1$  and  $f_2$  are continuous

Def If in addition  $f_1$  and  $f_2$  are analytic ( $\in C^\infty$ , expandible in a power-Taylor series) then  $G$  is called a Lie group.

Example:

The matrix rep on p. 32  $\begin{pmatrix} \cos \theta & -\sin \theta & 0 \\ \sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{pmatrix}$  of rotations in  $\mathbb{R}^3$  around the z-axis parametrized

by an angle  $\theta \in [0, 2\pi)$  is a Lie group. It is also called a compact group since  $\theta \in [0, 2\pi) \subseteq \mathbb{R}$  which is a set that is : bounded & single component (not closed)

Examples:

• we have already seen that linear maps / matrix reps that preserve the scalar product are unitary. How many real parameters do they contain?

complex scalar product on  $\mathbb{C}^n$ :

$$\boxed{U(n)} \subseteq GL(n, \mathbb{C}) \quad \text{mit } U_{ij}^{-1} = U_{ji}^* \text{ or } UU^T = \mathbb{1}_{n \times n}$$

- has  $n^2$  complex, so  $2n^2$  real entries, plus constraints to satisfy

$$U_{ij} U_{kj}^* = U_{ie} U_{je}^* = \delta_{ij} \quad \begin{aligned} &\bullet \quad i=j: n \text{ real eqs.} \\ &\bullet \quad i < j: \frac{n(n-1)}{2} \text{ complex eqs} \\ &\quad (\neq j \text{ gives the same eqs.}) \end{aligned}$$

$\Rightarrow$  we have

$$2n^2 - n - 2 \cdot \frac{n(n-1)}{2} = \underline{n^2 \text{ real parameters}}$$

$\boxed{SU(n)}$  has in addition  $\det U = +1$ :

$n$  real eq. ( $\det U^* = \det U^T = \det U^{-1} = \frac{1}{\det U} = 1$  is the same)

$$\Rightarrow \underline{n^2 - 1 \text{ real parameters}}$$