

Some reminders & formal definition of a Lie Algebra

- G Lie group, parametrised by $\dim[G]$ real parameters

$$\Rightarrow D(\theta) \left\{ \begin{array}{l} g \rightarrow \exp \left[i \sum_{a=1}^{\dim[G]} \theta^a T^a \right] \\ G \rightarrow G \subset (n, \mathbb{C}) \end{array} \right. \begin{array}{l} \text{is an } n \times n \text{ dimensional} \\ \text{unitary matrix rep of } G \end{array}$$

which lies in the component \mathcal{K} of G containing the identity

The generators T^a $a=1, \dots, \dim[G]$ satisfy $T^a = T^{a\dagger}$,

and if furthermore $\det D(\theta) = 1 \quad \forall \theta \Rightarrow \text{Tr } T^a = 0 \quad \forall a$

Furthermore $[T^a, T^b] = i f_{ab}^c T^c$ holds, with $f_{ab}^c \in \mathbb{R}$.

* on p. 38 we defined a vector space over a field.

- The $n \times n$ dim generators span a $\dim[G]$ dimensional vector space over \mathbb{C} , with elements

$$\theta = \theta^a T^a, \quad \text{addition } \theta + \varphi = \theta^a T^a + \varphi^a T^a = (\theta^a + \varphi^a) T^a$$

and scalar multiplication $\alpha \theta = (\alpha \theta^a) T^a$, satisfying all rules A0-A4 and B0-B4 on p 38

- Def A vector space V over field K is called an Algebra over K if there exists a binary operation (multiplication) $V \times V \rightarrow V$ that satisfies

$$\forall x, y, z \in V \quad \begin{array}{l} (x+y) \cdot z = x \cdot z + y \cdot z \\ x \cdot (y+z) = x \cdot y + x \cdot z \end{array} \quad \begin{array}{l} \text{left \& right} \\ \text{distributivity} \end{array}$$

$$\forall \alpha \in K \quad \alpha(x \cdot y) = (\alpha x) \cdot y = x \cdot (\alpha y)$$

Def a vector space V over K is called a Lie Algebra

if the binary operation $x \cdot y \equiv [x, y]_{\mathcal{L}}$

satisfies in addition $V \times V \longrightarrow V$

$[x, x]_{\mathcal{L}} = 0$ and the Jacobi identity holds,

$$[x, [y, z]]_{\mathcal{L}} + [z, [x, y]]_{\mathcal{L}} + [y, [z, x]]_{\mathcal{L}} = 0$$

Such a binary op. is then called a Lie bracket.

Examples for Lie bracket ^{for matrices} are the commutator or the Poisson-bracket

\Rightarrow equipped with the binary operation given by the commutator of the $n \times n$ matrices

$$T^a \cdot T^b \equiv [T^a, T^b] = i f^{ab}_c T^c$$

$$\mathfrak{g} \times \mathfrak{g} \longrightarrow \mathfrak{g}$$

the vector space \mathfrak{g} of our generators T^a , $a = 1, \dots, \dim(\mathfrak{g})$

is a Lie-algebra in the above sense.

• one can also define a scalar product on the vector space of \mathfrak{g} by $(A, B) = \text{Tr}(A B)$. The basis T_a $a=1, \dots, \dim[\mathfrak{g}]$ is called

orthogonal if $\text{Tr}(T_a T_b) = \delta_{ab}$ and orthonormal if $T = \frac{1}{\sqrt{2}}$.

* the requirement that the basis T^a is complete leads to a completeness relation $\sum_{a=1}^{\dim} T^a_i T^a_{jk} = \dots$ (\rightarrow exercise for $SU(n)$)
idea see next page

Representations and Lie-Algebras:

• we can also consider more general reps $D: \mathfrak{G} \rightarrow V$ for V some vector space of $\dim d$, that are not necessarily unitary. What we've said remains true: $D = \exp\left[i \sum_{a=1}^{\dim} \theta^a T^a\right]$, with the same # of T^a 's but these are no longer hermitian in general, and are of $\dim d$.

* we distinguish 3 important examples:

the conjugate rep: $D_* = \exp\left[-i \sum_{a=1}^{\dim} \theta^a T_a^*\right]$

\Rightarrow the generators $-T_a^*$ satisfy the same Lie algebra (which is rep. indep.)

$$[T_a, T_b] = i f_{ab}^c T_c \Leftrightarrow [-T_a^*, -T_b^*] = -i f_{ab}^c T_c^*$$

the adjoint rep (A)

the generators $(T_a)_b^c = -i f_{ab}^c$ satisfy the same Lie algebra (eq 8.2)

Because $b, c = 1, \dots, \dim[\mathfrak{g}]$ this is the dimension $d = \dim[\mathfrak{g}]$ of the rep.
e.g. $SU(n) \Rightarrow d = n^2 - 1$

the fundamental rep (F) as before which is $d = n$ dimensional

for $(SU(n), SO(n))$. It is often denoted by \mathbf{n} .

$$M = \Theta_a \Gamma^a$$

choose $\Gamma_a \Gamma_b = \Gamma \delta_{ab}$ ON
basis

$$\Rightarrow \Gamma_c (M \Gamma^b) = \Gamma_c (\Theta_a \Gamma^a \Gamma^b) = \Theta_a \delta^{ab} \Gamma = \Theta^b \Gamma$$

$$\Rightarrow M = \frac{1}{\Gamma} (\Gamma_c (M \Gamma^c)) \Gamma^0$$

$$M_{ij} = \frac{1}{\Gamma} M_{kc} \Gamma_{ek}^a \Gamma_{ij}^a$$

\Rightarrow

- these reps differ - in the normalisation of the scalar product T
- in the dimension d
- in the quadratic Casimir operator C :

$$\boxed{\frac{\dim}{2} \sum_{a=1} (T_a T_a)_{ij} = C \delta_{ij}}$$

eg. $SU(2)$: - find \bar{F} has $T_F = \frac{1}{2}$ & orthogonal $d_F = n$, $C_F = \frac{n^2 - 1}{n}$
 $\text{Tr}(T_a^F T_b^F) = \delta_{ab} \frac{1}{2}$ $T_a = -\frac{1}{2} \sigma_a$ (from completeness relation)

$SU(n)$ - adj. it has $T_A = n$, $d_A = n^2 - 1$ ($= \dim[G]$), $C_A = n$
 & orthogonal
 $\text{Tr}(T_a^A T_b^A) = n \delta_{ab}$

* from the adjoint rep we can learn more about the f_{abc} :

The corresp. scalar prod. is called $(A, B) \in G$

Killing form $(A, B) \equiv \text{Tr}(\underset{\text{adj rep}}{D_A(A)} \underset{\text{adj rep}}{D_A(B)})$

applied to the generators it is called

Cartan metric $g_{ab} \equiv \text{Tr}_A(T_a^A T_b^A) = f_{ac}^d f_{bd}^c = g_{ba}$

let us define $f_{abc} \equiv f_{ab}^d g_{dc}$

\Rightarrow f_{abc} is totally antisym. in all indices (we were sloppy in notation before, but due to $g_{ab} = \delta_{ab}$ for $SU(2)$ this didn't matter there)

proof: $\text{Tr}(A[B]C) = \text{Tr}(ABC - BAC)$

$= \text{Tr}(C[BA]) = \text{Tr}(C[A]B)$ due to cyclicity

now $\text{Tr}_A([T_a^A, T_b^A] T_c^A) = i f_{ab}^d \text{Tr}_A(T_d^A T_c^A) = i f_{ab}^d g_{dc} = i f_{cab} = -i f_{acb}$ etc